

# CHAPTER 15 VECTOR CALCULUS

## 15.1 Vector Fields (page 554)

An ordinary function assigns a value  $f(x)$  to each point  $x$ . A vector field assigns a **vector**  $\mathbf{F}(x, y)$  to each point  $(x, y)$ . Think of the vector as going out from the point (not out from the origin). The vector field is like a head of hair! We are placing a straight hair at every point. Depending on how the hair is cut and how it is combed, the vectors have different lengths and different directions.

The vector at each point  $(x, y)$  has two components. Its horizontal component is  $M(x, y)$ , its vertical component is  $N(x, y)$ . *The vector field is*  $\mathbf{F}(x, y) = M(x, y)\mathbf{i} + N(x, y)\mathbf{j}$ . Remember: A vector from every point.

- Suppose all of the vectors  $\mathbf{F}(x, y)$  have length 1, and their directions are *outward* (or *radial*). Find their components  $M(x, y)$  and  $N(x, y)$ .
  - At a point like  $(3, 0)$  on the  $x$  axis, the outward direction is the  $x$  direction. The vector of length 1 from that point is  $\mathbf{F}(3, 0) = \mathbf{i}$ . This vector goes outward *from the point*. At  $(0, 2)$  the outward vector is  $\mathbf{F}(0, 2) = \mathbf{j}$ . At the point  $(-2, 0)$  it is  $\mathbf{F}(-2, 0) = -\mathbf{i}$ . (The *minus*  $x$  direction is outward.) At every point  $(x, y)$ , the outward direction is parallel to  $x\mathbf{i} + y\mathbf{j}$ . This is the “position vector”  $\mathbf{R}(x, y)$ .

We want an outward spreading field of **unit** vectors. So divide the position vector  $\mathbf{R}$  by its length:

$$\mathbf{F}(x, y) = \frac{\mathbf{R}(x, y)}{|\mathbf{R}(x, y)|} = \frac{x\mathbf{i} + y\mathbf{j}}{\sqrt{x^2 + y^2}}. \quad \text{This special vector field is called } \mathbf{u}_r.$$

The letter  $\mathbf{u}$  is for “unit,” the subscript  $r$  is for “radial.” No vector is assigned to the origin, because the outward direction there can’t be decided. Thus  $\mathbf{F}(0, 0)$  is not defined (for this particular field). Then we don’t have to divide by  $r = \sqrt{x^2 + y^2} = 0$  at the origin.

To repeat: The field of outward unit vectors is  $\mathbf{u}_r = \frac{\mathbf{R}}{r}$ . Another way to write it is  $\mathbf{u}_r = \cos \theta \mathbf{i} + \sin \theta \mathbf{j}$ . The components are  $M(x, y) = \frac{x}{\sqrt{x^2 + y^2}} = \cos \theta$  and  $N(x, y) = \frac{y}{\sqrt{x^2 + y^2}} = \sin \theta$ .

- Suppose again that all the vectors  $\mathbf{F}(x, y)$  are unit vectors. But change their directions to be *perpendicular* to  $\mathbf{u}_r$ . The vector at  $(3, 0)$  is  $\mathbf{j}$  instead of  $\mathbf{i}$ . Find a formula for this “**unit spin field**.”
  - We want to take the vector  $\mathbf{u}_r$ , at each point  $(x, y)$  except the origin, and **turn that vector by**  $90^\circ$ . The turn is counterclockwise and the new vector is called  $\mathbf{u}_\theta$ . It is still a unit vector, and its dot product with  $\mathbf{u}_r$  is zero. Here it is, written in two or three different ways:

$$\begin{aligned} \mathbf{u}_\theta &= -\sin \theta \mathbf{i} + \cos \theta \mathbf{j} \text{ is perpendicular to } \mathbf{u}_r = \cos \theta \mathbf{i} + \sin \theta \mathbf{j} \\ \mathbf{u}_\theta &= \frac{-y}{\sqrt{x^2 + y^2}} \mathbf{i} + \frac{x}{\sqrt{x^2 + y^2}} \mathbf{j} = \frac{1}{r}(-y \mathbf{i} + x \mathbf{j}) = \mathbf{S}/r. \end{aligned}$$

Where  $\mathbf{R} = x\mathbf{i} + y\mathbf{j}$  was the position field (outward),  $\mathbf{S} = -y\mathbf{i} + x\mathbf{j}$  is the **spin field** (around the origin). The lengths of  $\mathbf{R}$  and  $\mathbf{S}$  are both  $r = \sqrt{x^2 + y^2}$ , increasing as we move outward. For unit vectors  $\mathbf{u}_r$  and  $\mathbf{u}_\theta$  we divide by this length.

Two other important fields are  $\mathbf{R}/r^2$  and  $\mathbf{S}/r^2$ . At each point one is still outward (parallel to  $\mathbf{R}$ ) and the other is “turning” (parallel to  $\mathbf{S}$ ). But now the lengths decrease as we go outward. The length of  $\frac{\mathbf{R}}{r^2}$  is  $\frac{r}{r^2}$  or  $\frac{1}{r}$ . This is closer to a typical men’s haircut.

So far we have six vector fields: three radial fields  $\mathbf{R}$  and  $\mathbf{u}_r = \mathbf{R}/r$  and  $\mathbf{R}/r^2$  and three spin fields  $\mathbf{S}$  and  $\mathbf{u}_\theta = \mathbf{S}/r$  and  $\mathbf{S}/r^2$ . The radial fields point along **rays**, out from the origin. The spin fields are tangent to **circles**, going around the origin. These rays and circles are the **field lines** or **streamlines** for these particular vector fields.

The field lines give the direction of the vector  $\mathbf{F}(x, y)$  at each point. The length is not involved (that is why  $\mathbf{S}$  and  $\mathbf{S}/r$  and  $\mathbf{S}/r^2$  all have the same field lines). The direction is tangent to the field line so the slope of that line is  $dy/dx = N(x, y)/M(x, y)$ .

3. Find the field lines (streamlines) for the vector field  $\mathbf{F}(x, y) = x \mathbf{i} - y \mathbf{j}$ .

- Solve  $dy/dx = N/M = -y/x$  by separating variables. We have  $\frac{dy}{y} = -\frac{dx}{x}$ . Integration gives

$$\ln y = -\ln x + C. \text{ Therefore } \ln x + \ln y = C \text{ or } \ln xy = C \text{ or } xy = c.$$

The field lines  $xy = c$  are hyperbolas. The vectors  $\mathbf{F} = x\mathbf{i} - y\mathbf{j}$  are tangent to those hyperbolas.

4. **A function  $f(x, y)$  produces a gradient field  $\mathbf{F}$ .** Its components are  $M = \frac{\partial f}{\partial x}$  and  $N = \frac{\partial f}{\partial y}$ . This field has the special symbol  $\mathbf{F} = \nabla f$ . Describe this gradient field for the particular function  $f(x, y) = xy$ .

- The partial derivatives of  $f(x, y) = xy$  are  $\partial f/\partial x = y$  and  $\partial f/\partial y = x$ . Therefore the gradient field is  $\nabla f = y\mathbf{i} + x\mathbf{j}$ . Not a radial field and not a spin field.

Remember that the gradient vector gives the direction in which  $f(x, y)$  changes fastest. This is the “steepest direction.” Tangent to the curve  $f(x, y) = c$  there is no change in  $f$ . **Perpendicular to the curve there is maximum change.** This is the gradient direction. So the gradient field  $y\mathbf{i} + x\mathbf{j}$  of Problem 4 is perpendicular to the field  $x\mathbf{i} - y\mathbf{j}$  of Problem 3.

In Problem 3, the field is tangent to the hyperbolas  $xy = c$ . In Problem 4, the field is perpendicular to those hyperbolas. The hyperbolas are called **equipotential lines** because the “potential”  $xy$  is “equal” (or constant) along those curves  $xy = c$ .

#### **Read-throughs and selected even-numbered solutions :**

A vector field assigns a vector to each point  $(x, y)$  or  $(x, y, z)$ . In two dimensions  $\mathbf{F}(x, y) = M(x, y)\mathbf{i} + N(x, y)\mathbf{j}$ . An example is the position field  $\mathbf{R} = x\mathbf{i} + y\mathbf{j} (+z\mathbf{k})$ . Its magnitude is  $|\mathbf{R}| = r$  and its direction is **out from the origin**. It is the gradient field for  $f = \frac{1}{2}(x^2 + y^2)$ . The level curves are **circles**, and they are **perpendicular** to the vectors  $\mathbf{R}$ .

Reversing this picture, the spin field is  $\mathbf{S} = -y\mathbf{i} + x\mathbf{j}$ . Its magnitude is  $|\mathbf{S}| = r$  and its direction is **around the origin**. It is not a gradient field, because no function has  $\partial f/\partial x = -y$  and  $\partial f/\partial y = x$ .  $\mathbf{S}$  is the velocity field for flow going **around the origin**. The streamlines or **field lines** or **integral curves** are **circles**. The flow field  $\rho\mathbf{V}$  gives the rate at which **mass** is moved by the flow.

A gravity field from the origin is proportional to  $\mathbf{F} = \mathbf{R}/r^3$  which has  $|\mathbf{F}| = 1/r^2$ . This is Newton’s inverse square law. It is a gradient field, with potential  $f = 1/r$ . The equipotential curves  $f(x, y) = c$  are **circles**. They are **perpendicular** to the field lines which are **rays**. This illustrates that the **gradient** of a function  $f(x, y)$  is **perpendicular** to its level curves.

The velocity field  $y\mathbf{i} + x\mathbf{j}$  is the gradient of  $f = xy$ . Its streamlines are **hyperbolas**. The slope  $dy/dx$  of a streamline equals the ratio  $N/M$  of velocity components. The field is **tangent** to the streamlines. Drop a leaf onto the flow, and it goes along a **streamline**.

- 2**  $x\mathbf{i} + \mathbf{j}$  is the gradient of  $f(x, y) = \frac{1}{2}x^2 + y$ , which has **parabolas**  $\frac{1}{2}x^2 + y = c$  as equipotentials (they open down). The streamlines solve  $dy/dx = 1/x$  (this is  $N/M$ ). So  $y = \ln x + C$  gives the streamlines.
- 6**  $x^2\mathbf{i} + y^2\mathbf{j}$  is the gradient of  $f(x, y) = \frac{1}{3}(x^3 + y^3)$ , which has closed curves  $x^3 + y^3 = \text{constant}$  as equipotentials. The streamlines solve  $dy/dx = y^2/x^2$  or  $dy/y^2 = dx/x^2$  or  $y^{-1} = x^{-1} + \text{constant}$ .
- 14**  $\frac{\partial f}{\partial x} = 2x - 2$  and  $\frac{\partial f}{\partial y} = 2y$ ;  $\mathbf{F} = (2x - 2)\mathbf{i} + 2y\mathbf{j}$  leads to circles  $(x - 1)^2 + y^2 = c$  around the center  $(1, 0)$ .
- 26**  $f(x, y) = \frac{1}{2} \ln(x^2 + y^2) = \ln \sqrt{x^2 + y^2}$ . This comes from  $\frac{\partial f}{\partial x} = \frac{x}{x^2 + y^2}$  or  $f = \int \frac{x dx}{x^2 + y^2}$ .
- 32** From the gradient of  $y - x^2$ ,  $\mathbf{F}$  must be  $-2x\mathbf{i} + \mathbf{j}$  (or this is  $-\mathbf{F}$ ).
- 36**  $\mathbf{F}$  is the gradient of  $f = \frac{1}{2}ax^2 + bxy + \frac{1}{2}cy^2$ . The equipotentials are **ellipses** if  $ac > b^2$  and **hyperbolas** if  $ac < b^2$ . (If  $ac = b^2$  we get straight lines.)

## 15.2 Line Integrals (page 562)

The most common line integral is along the  $x$  axis. We have a function  $y(x)$  and we integrate to find  $\int y(x)dx$ . Normally this is just called an integral, without the word “line.” But now we have functions defined at every point in the  $xy$  plane, so we can integrate along curves. A better word for what is coming would be “curve” integral.

Think of a curved wire. The density of the wire is  $\rho(x, y)$ , possibly varying along the wire. Then the total mass of the wire is  $\int \rho(x, y)ds$ . This is a line integral or curve integral (or wire integral). Notice  $ds = \sqrt{(dx)^2 + (dy)^2}$ . We use  $dx$  for integrals along the  $x$  axis and  $dy$  up the  $y$  axis and  $ds$  for integrals along other lines and curves.

1. A circular wire of radius  $R$  has density  $\rho(x, y) = x^2y^2$ . How can you compute its mass  $M = \int \rho ds$ ?
  - Describe the circle by  $x = R \cos t$  and  $y = R \sin t$ . You are free to use  $\theta$  instead of  $t$ . The point is that we need a parameter to describe the path and to compute  $ds$ :

$$ds = \sqrt{\left(\frac{dx}{dt}\right)^2 + \left(\frac{dy}{dt}\right)^2} dt = \sqrt{R^2 \sin^2 t + R^2 \cos^2 t} dt = R dt.$$

Then the mass integral is  $M = \int x^2y^2 ds = \int_{t=0}^{2\pi} (R^2 \cos^2 t)(R^2 \sin^2 t)R dt$ . I won't integrate.

This chapter is about vector fields. But we integrate scalar functions (like the density  $\rho$ ). So if we are given a vector  $\mathbf{F}(x, y)$  at each point, we take its dot product with another vector – to get an ordinary scalar function to be integrated. Two dot products are by far the most important:

$$\begin{aligned} \mathbf{F} \cdot \mathbf{T} ds &= \text{component of } \mathbf{F} \text{ tangent to the curve} = M dx + N dy \\ \mathbf{F} \cdot \mathbf{n} ds &= \text{component of } \mathbf{F} \text{ normal to the curve} = M dy - N dx \end{aligned}$$

The unit tangent vector is  $\mathbf{T} = \frac{d\mathbf{R}}{ds}$  in Chapter 12. Then  $\mathbf{F} \cdot \mathbf{T} ds$  is  $\mathbf{F} \cdot d\mathbf{R} = (M \mathbf{i} + N \mathbf{j}) \cdot (dx \mathbf{i} + dy \mathbf{j})$ . This dot product is  $M dx + N dy$ , which we integrate. Its integral is the **work** done by  $\mathbf{F}$  along the curve. Work is force times distance, but the distance is measured parallel to the force. This is why the tangent component  $\mathbf{F} \cdot \mathbf{T}$  goes into the work integral.

2. Compute the work by the force  $\mathbf{F} = x\mathbf{i}$  around the unit circle  $x = \cos t$ ,  $y = \sin t$ .

- Work =  $\int M dx + N dy = \int x dx + 0 dy = \int_{t=0}^{2\pi} (\cos t)(-\sin t dt)$ . This integral is zero!

3. Compute the work by  $\mathbf{F} = x\mathbf{i}$  around a square: along  $y = 0$ , up  $x = 1$ , back along  $y = 1$ , back down  $x = 0$ .

- Along the  $x$  axis, the direction is  $\mathbf{T} = \mathbf{i}$  and  $\mathbf{F} \cdot \mathbf{T} = x\mathbf{i} \cdot \mathbf{i} = x$ . Work =  $\int_0^1 x dx = \frac{1}{2}$ .

Up the line  $x = 1$ , the direction is  $\mathbf{T} = \mathbf{j}$ . Then  $\mathbf{F} \cdot \mathbf{T} = x\mathbf{i} \cdot \mathbf{j} = 0$ . No work.

Back along  $y = 1$  the direction is  $\mathbf{T} = -\mathbf{i}$ . Then  $\mathbf{F} \cdot \mathbf{T} = -x$ . The work is  $\int \mathbf{F} \cdot \mathbf{T} ds = \int -x dx = -\frac{1}{2}$ .

Note! You might think the integral should be  $\int_1^0 (-x) dx = +\frac{1}{2}$ . Wrong. Going left,  $ds$  is  $-dx$ .

The work down the  $y$  axis is again zero.  $\mathbf{F} = x\mathbf{i}$  is perpendicular to the movement  $d\mathbf{R} = \mathbf{j} dy$ . So  $\mathbf{F} \cdot \mathbf{T} = 0$ .

Total work around square =  $\frac{1}{2} + 0 - \frac{1}{2} + 0 = \text{zero}$ .

4. Does the field  $\mathbf{F} = x\mathbf{i}$  do zero work around every closed path? If so, why?

- Yes, the line integral  $\int \mathbf{F} \cdot \mathbf{T} ds = \int x dx + 0 dy$  is always zero around closed paths. The antiderivative is  $f = \frac{1}{2}x^2$ . **When the start and end are the same point  $P$  the definite integral is  $f(P) - f(P) = 0$ .**

We used the word “antiderivative.” From now on we will say “**potential function**.” This is a function  $f(x, y)$  - if it exists - such that  $df = M dx + N dy$ :

The potential function has  $\frac{\partial f}{\partial x} = M$  and  $\frac{\partial f}{\partial y} = N$ . Then  $\int_P^Q M dx + N dy = \int_P^Q df = f(Q) - f(P)$ .

The field  $\mathbf{F}(x, y)$  is the **gradient** of the potential function  $f(x, y)$ . Our example has  $f = \frac{1}{2}x^2$  and  $\mathbf{F} = \nabla f = x\mathbf{i}$ . Conclusion: Gradient fields are conservative. The work around a *closed* path is zero.

5. Does the field  $\mathbf{F} = y\mathbf{i}$  do zero work around every closed path? If not, why not?

- This is not a gradient field. There is no potential function that  $\frac{\partial f}{\partial x} = M$  and  $\frac{\partial f}{\partial y} = N$ . We are asking for  $\frac{\partial f}{\partial x} = y$  and  $\frac{\partial f}{\partial y} = 0$  which is impossible. The work around the unit circle  $x = \cos t$ ,  $y = \sin t$  is

$$\int M dx + N dy = \int y dx = \int_{t=0}^{2\pi} (\sin t)(-\sin t dt) = -\pi. \text{ Not zero!}$$

*Important* A gradient field has  $M = \frac{\partial f}{\partial x}$  and  $N = \frac{\partial f}{\partial y}$ . Every function has equal mixed derivatives  $\frac{\partial^2 f}{\partial x \partial y} = \frac{\partial^2 f}{\partial y \partial x}$ . Therefore the gradient field has  $\partial M / \partial y = \partial N / \partial x$ . This is the quick test “**D**” for a gradient field.  $\mathbf{F} = y\mathbf{i}$  fails this test as we expected, because  $\partial M / \partial y = 1$  and  $\partial N / \partial x = 0$ .

### Read-throughs and selected even-numbered solutions :

Work is the **integral** of  $\mathbf{F} \cdot d\mathbf{R}$ . Here  $\mathbf{F}$  is the **force** and  $\mathbf{R}$  is the **position**. The **dot** product finds the component of  $\mathbf{F}$  in the direction of movement  $d\mathbf{R} = dx \mathbf{i} + dy \mathbf{j}$ . The straight path  $(x, y) = (t, 2t)$  goes from  $(0,0)$  at  $t = 0$  to  $(1,2)$  at  $t = 1$  with  $d\mathbf{R} = dt \mathbf{i} + 2dt \mathbf{j}$ .

Another form of  $d\mathbf{R}$  is  $\mathbf{T}ds$ , where  $\mathbf{T}$  is the **unit tangent** vector to the path and the arc length has  $ds = \sqrt{(dx/dt)^2 + (dy/dt)^2}$ . For the path  $(t, 2t)$ , the unit vector  $\mathbf{T}$  is  $(\mathbf{i} + 2\mathbf{j})/\sqrt{5}$  and  $ds = \sqrt{5}dt$ . For  $\mathbf{F} = 3\mathbf{i} + \mathbf{j}$ ,  $\mathbf{F} \cdot \mathbf{T} ds$  is still  $5dt$ . This  $\mathbf{F}$  is the gradient of  $f = 3x + y$ . The change in  $f = 3x + y$  from  $(0,0)$  to  $(1,2)$  is  $5$ .

When  $\mathbf{F} = \text{grad } f$ , the dot product  $\mathbf{F} \cdot d\mathbf{R}$  is  $(\partial f / \partial x)dx + (\partial f / \partial y)dy = df$ . The work integral from  $P$  to  $Q$  is  $\int df = f(\mathbf{Q}) - f(\mathbf{P})$ . In this case the work depends on the **endpoints** but not on the **path**. Around a closed

path the work is **zero**. The field is called **conservative**.  $\mathbf{F} = (1 + y)\mathbf{i} + x\mathbf{j}$  is the gradient of  $f = \mathbf{x} + \mathbf{xy}$ . The work from  $(0,0)$  to  $(1,2)$  is **3**, the change in potential.

For the spin field  $\mathbf{S} = -y\mathbf{i} + x\mathbf{j}$ , the work does depend on the path. The path  $(x, y) = (3 \cos t, 3 \sin t)$  is a circle with  $\mathbf{S} \cdot d\mathbf{R} = -y dx + x dy = 9 dt$ . The work is  $18\pi$  around the complete circle. Formally  $\int g(x, y) ds$  is the limit of the sum  $\sum g(x_i, y_i) \Delta s_i$ .

The four equivalent properties of a conservative field  $\mathbf{F} = M\mathbf{i} + N\mathbf{j}$  are **A: zero work around closed paths**, **B: work depends only on endpoints**, **C: gradient field**, **D:  $\partial M/\partial y = \partial N/\partial x$** . Test **D** is passed by  $\mathbf{F} = (y + 1)\mathbf{i} + x\mathbf{j}$ . The work  $\int \mathbf{F} \cdot d\mathbf{R}$  around the circle  $(\cos t, \sin t)$  is **zero**. The work on the upper semicircle equals the work on the lower semicircle (clockwise). This field is the gradient of  $f = \mathbf{x} + \mathbf{xy}$ , so the work to  $(-1, 0)$  is **-1** starting from  $(0,0)$ .

**4** Around the square  $0 \leq x, y \leq 3$ ,  $\int_3^0 y dx = -9$  along the top (backwards) and  $\int_0^3 -x dy = -9$  up the right side. All other integrals are zero: answer **-18**. By Section 15.3 this integral is always  $-2 \times \text{area}$ .

**8 Yes** The field  $x\mathbf{i}$  is the gradient of  $f = \frac{1}{2}x^2$ . Here  $M = x$  and  $N = 0$  so  $\int_P^Q M dx + N dy = f(Q) - f(P)$ .

More directly: up and down movement has no effect on  $\int x dx$ .

**10 Not much**. Certainly the limit of  $\Sigma(\Delta s)^2$  is zero.

**14**  $\frac{\partial N}{\partial x} = \frac{\partial M}{\partial y}$  and  $\mathbf{F}$  is the gradient of  $f = xe^y$ . Then  $\int \mathbf{F} \cdot d\mathbf{R} = f(Q) - f(P) = -1$ .

**18**  $\frac{\mathbf{R}}{r^n}$  has  $M = \frac{x}{(x^2+y^2)^{n/2}}$  and  $\frac{\partial M}{\partial y} = -xny(x^2+y^2)^{-(n/2)-1}$ . This agrees with  $\frac{\partial N}{\partial x}$  so  $\frac{\mathbf{R}}{r^n}$  is a

gradient field for all  $n$ . The potential is  $f = \frac{r^{2-n}}{2-n}$  or  $f = \ln r$  when  $n = 2$ .

**32**  $\frac{\partial N}{\partial x} \neq \frac{\partial M}{\partial y}$  (not conservative):  $\int x^2y dx + xy^2 dy = \int_0^1 2t^3 dt = \frac{1}{2}$  but  $\int_0^1 t^2(t^2) dt + t(t^2)^2(2t dt) = \frac{17}{35}$ .

**34** The potential is  $f = \frac{1}{2} \ln(x^2 + y^2 + 1)$ . Then  $f(1, 1) - f(0, 0) = \frac{1}{2} \ln 3$ .

## 15.3 Green's Theorem (page 571)

The last section studied line integrals of  $\mathbf{F} \cdot \mathbf{T} ds$ . This section connects them to *double integrals*. The work can be found by integrating around the curve (with  $ds$ ) or *inside the curve* (with  $dA = dx dy$ ). The connection is by Green's Theorem. The theorem is for integrals around **closed curves**:

$$\int_C M dx + N dy = \iint_R \left( \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dx dy.$$

We see again that this is zero for gradient fields. Their test is  $\frac{\partial N}{\partial x} = \frac{\partial M}{\partial y}$  so the double integral is immediately zero.

1. Compute the work integral  $\int M dx + N dy = \int y dx$  for the force  $\mathbf{F} = y\mathbf{i}$  around the unit circle.

- Use Green's Theorem with  $M = y$  and  $N = 0$ . The line integral equals the double integral of  $\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} = -1$ . Integrate  $-1$  over a circle of area  $\pi$  to find the answer  $-\pi$ . This agrees with Question 5 in the previous section of the Guide. It also means that the true-false Problem 15.2.44c has answer "False."

**Special case** If  $M = -\frac{1}{2}y$  and  $N = \frac{1}{2}x$  then  $\partial N/\partial x = \frac{1}{2}$  and  $-\partial M/\partial y = \frac{1}{2}$ . Therefore Green's Theorem is  $\frac{1}{2} \int_C -y dx + x dy = \iint_R \left( \frac{1}{2} + \frac{1}{2} \right) dx dy = \text{area of } R$ .

2. Use that special case to find the area of a triangle with corners  $(0,0)$ ,  $(x_1, y_1)$ , and  $(x_2, y_2)$ .

- We have to integrate  $\frac{1}{2} \int -y dx + x dy$  around the triangle to get the area. The first side has  $x = x_1 t$  and  $y = y_1 t$ . As  $t$  goes from 0 to 1, the point  $(x, y)$  goes from  $(0, 0)$  to  $(x_1, y_1)$ . The integral is  $\int (-y_1 t)(x_1 dt) + (x_1 t)(y_1 dt) = 0$ . Similarly the line integral between  $(0, 0)$  and  $(x_2, y_2)$  is zero. The third side has  $x = x_1 + t(x_2 - x_1)$  and  $y = y_1 + t(y_2 - y_1)$ . It goes from  $(x_1, y_1)$  to  $(x_2, y_2)$ .

$$\frac{1}{2} \int -y dx + x dy = (\text{substitute } x \text{ and } y \text{ simplify}) = \frac{1}{2} \int_0^1 (x_1 y_2 - x_2 y_1) dt = \frac{1}{2} (x_1 y_2 - x_2 y_1).$$

This is the area of the triangle. It is half the parallelogram area =  $\begin{vmatrix} x_1 & y_1 \\ x_2 & y_2 \end{vmatrix}$  in Chapter 11.

**Green's Theorem also applies to the flux integral  $\int \mathbf{F} \cdot \mathbf{n} ds$  around a closed curve  $C$ .** Now we are integrating  $M dy - N dx$ . By changing letters in the first form (the work form) of Green's Theorem, we get the second form (the flux form):

$$\text{Flow through curve} = \int_C M dy - N dx = \iint_R \left( \frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} \right) dx dy.$$

We are no longer especially interested in gradient fields (which give zero work). Now we are interested in source-free fields (which give zero flux). The new test is  $\frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} = 0$ . This quantity is the **divergence** of the field  $\mathbf{F} = M(x, y)\mathbf{i} + N(x, y)\mathbf{j}$ . A source-free field has *zero divergence*.

3. Is the position field  $\mathbf{F} = x \mathbf{i} + y \mathbf{j}$  source-free? If not, find the flux  $\int \mathbf{F} \cdot \mathbf{n} ds$  going out of a unit square.
  - The divergence of this  $\mathbf{F}$  is  $\frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} = \frac{\partial x}{\partial x} + \frac{\partial y}{\partial y} = 1 + 1 = 2$ . The field is not source-free. The flux is not zero. Green's Theorem gives flux =  $\iint 2 dx dy = 2 \times \text{area of region} = 2$ , for a unit square.
4. Is the field  $\mathbf{F} = x \mathbf{i} - y \mathbf{j}$  source-free? Is it also a gradient field?
  - *Yes and yes.* The field has  $M = x$  and  $N = -y$ . It passes both tests. Test **D** for a gradient field is  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$  which is  $0 = 0$ . Test **H** for a source-free field is  $\frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} = 1 - 1 = 0$ .

This gradient field has the potential function  $f(x, y) = \frac{1}{2}(x^2 - y^2)$ . This source-free field also has a **stream function**  $g(x, y) = xy$ . The stream function satisfies  $\frac{\partial g}{\partial y} = M$  and  $\frac{\partial g}{\partial x} = -N$ . Then  $g(x, y)$  is the antiderivative for the flux integral  $\int M dy - N dx$ . When it goes around a closed curve from  $P$  to  $P$ , the integral is  $g(P) - g(P) = 0$ . This is what we expect for source-free fields, with stream functions.

5. Show how the combination of "conservative" plus "source-free" leads to Laplace's equation for  $f$  (and  $g$ ).
  - $f_{xx} + f_{yy} = M_x + N_y$  because  $f_x = M$  and  $f_y = N$ . But source-free means  $M_x + N_y = 0$ .

### **Read-throughs and selected even-numbered solutions :**

The work integral  $\oint M dx + N dy$  equals the double integral  $\iint (\mathbf{N}_x - \mathbf{M}_y) dx dy$  by Green's Theorem. For  $\mathbf{F} = 3\mathbf{i} + 4\mathbf{j}$  the work is zero. For  $\mathbf{F} = x\mathbf{j}$  and  $-y\mathbf{i}$  the work equals the area of  $R$ . When  $M = \partial f / \partial x$  and  $N = \partial f / \partial y$ , the double integral is zero because  $f_{xy} = f_{yx}$ . The line integral is zero because  $f(\mathbf{Q}) = f(\mathbf{P})$  when  $\mathbf{Q} = \mathbf{P}$  (closed curve). An example is  $\mathbf{F} = y \mathbf{i} + x \mathbf{j}$ . The direction on  $C$  is **counterclockwise** around the outside and **clockwise** around the boundary of a hole. If  $R$  is broken into very simple pieces with crosscuts between them, the integrals of  $\mathbf{M} dx + \mathbf{N} dy$  cancel along the crosscuts.

Test **D** for gradient fields is  $\partial M / \partial y = \partial N / \partial x$ . A field that passes this test has  $\oint \mathbf{F} \cdot d\mathbf{R} = 0$ . There is a solution to  $f_x = M$  and  $f_y = N$ . Then  $df = M dx + N dy$  is an **exact** differential. The spin field  $\mathbf{S}/r^2$  passes test

**D** except at  $r = 0$ . Its potential  $f = \theta$  increases by  $2\pi$  going around the origin. The integral  $\iint (N_x - M_y) dx dy$  is not zero but  $2\pi$ .

The flow form of Green's theorem is  $\oint_C \mathbf{M} dy - N dx = \iint_R (\mathbf{M}_x + \mathbf{N}_y) dx dy$ . The normal vector in  $\mathbf{F} \cdot \mathbf{n} ds$  points out across  $C$  and  $|\mathbf{n}| = 1$  and  $\mathbf{n} ds$  equals  $dy \mathbf{i} - dx \mathbf{j}$ . The divergence of  $M\mathbf{i} + N\mathbf{j}$  is  $\mathbf{M}_x + \mathbf{N}_y$ . For  $\mathbf{F} = x\mathbf{i}$  the double integral is  $\iint 1 dt = \text{area}$ . There is a source. For  $\mathbf{F} = y\mathbf{i}$  the divergence is zero. The divergence of  $\mathbf{R}/r^2$  is zero except at  $r = 0$ . This field has a point source.

A field with no source has properties **E = zero flux through C, F = equal flux across all paths from P to Q, G = existence of stream function, H = zero divergence**. The stream function  $g$  satisfies the equations  $\partial g/\partial y = M$  and  $\partial g/\partial x = -N$ . Then  $\partial M/\partial x + \partial N/\partial y = 0$  because  $\partial^2 g/\partial x \partial y = \partial^2 g/\partial y \partial x$ . The example  $\mathbf{F} = y\mathbf{i}$  has  $g = \frac{1}{2}y^2$ . There is not a potential function. The example  $\mathbf{F} = x\mathbf{i} - y\mathbf{j}$  has  $g = xy$  and also  $f = \frac{1}{2}x^2 - \frac{1}{2}y^2$ . This  $f$  satisfies Laplace's equation  $f_{xx} + f_{yy} = 0$ , because the field  $\mathbf{F}$  is both conservative and source-free. The functions  $f$  and  $g$  are connected by the Cauchy-Riemann equations  $\partial f/\partial x = \partial g/\partial y$  and  $\partial f/\partial y = -\partial g/\partial x$ .

- 4  $\oint y dx = \int_0^1 t(-dt) = -\frac{1}{2}$ ;  $M = y, N = 0, \iint (-1) dx dy = -\text{area} = -\frac{1}{2}$ .
- 12 Let  $R$  be the square with base from  $a$  to  $b$  on the  $x$  axis. Set  $\mathbf{F} = f(x)\mathbf{j}$  so  $M = 0$  and  $N = f(x)$ . The line integral  $\oint M dx + N dy$  is  $(b - a)f(b)$  up the right side minus  $(b - a)f(a)$  down the left side. The double integral is  $\iint \frac{df}{dx} dx dy = (b - a) \int_a^b \frac{df}{dx} dx$ . Green's Theorem gives equality; cancel  $b - a$ .
- 16  $\oint \mathbf{F} \cdot \mathbf{n} ds = \int xy dy = \frac{1}{2}$  up the right side of the square where  $\mathbf{n} = \mathbf{i}$  (other sides give zero).  
Also  $\int_0^1 \int_0^1 (y + 0) dx dy = \frac{1}{2}$ .
- 22  $\oint \mathbf{F} \cdot \mathbf{n} ds$  is the same through a square and a circle because the difference is  $\iint (\frac{\partial M}{\partial x} + \frac{\partial N}{\partial y}) dx dy = \iint \text{div } \mathbf{F} dx dy = 0$  over the region in between.
- 30  $\frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} = 3y^2 - 3y^2 = 0$ . Solve  $\frac{\partial g}{\partial y} = 3xy^2$  for  $g = xy^3$  and check  $\frac{\partial g}{\partial x} = y^3$ .
- 38  $g(Q) = \int_P^Q \mathbf{F} \cdot \mathbf{n} ds$  starting from  $g(P) = 0$ . Any two paths give the same integral because forward on one and back on the other gives  $\oint \mathbf{F} \cdot \mathbf{n} ds = 0$ , provided the tests  $E - H$  for a stream function are passed.

## 15.4 Surface Integrals (page 581)

The length of a curve is  $\int ds$ . The area of a surface is  $\iint dS$ . Curves are described by functions  $y = f(x)$  in single-variable calculus. Surfaces are described by functions  $z = f(x, y)$  in multivariable calculus. When you have worked with  $ds$ , you see  $dS$  as the natural next step:

$$ds = \sqrt{1 + \left(\frac{dy}{dx}\right)^2} dx \quad \text{and} \quad dS = \sqrt{1 + \left(\frac{\partial z}{\partial x}\right)^2 + \left(\frac{\partial z}{\partial y}\right)^2} dx dy$$

The basic step  $dx$  is along the  $x$  axis. The extra  $\left(\frac{dy}{dx}\right)^2$  in  $ds$  accounts for the extra length when the curve slopes up or down. Similarly  $dx dy$  is the area  $dA$  down in the base plane. The extra  $\left(\frac{\partial z}{\partial x}\right)^2$  and  $\left(\frac{\partial z}{\partial y}\right)^2$  account

for the extra area when the surface slopes up or down.

1. Find the length of the line  $x + y = 1$  cut off by the axes  $y = 0$  and  $x = 0$ . The line segment goes from  $(1,0)$  to  $(0,1)$ . Find the area of the plane  $x + y + z = 1$  cut off by the planes  $z = 0$  and  $y = 0$  and  $x = 0$ . This is a triangle with corners at  $(1,0,0)$  and  $(0,1,0)$  and  $(0,0,1)$ .

- The line  $y = 1 - x$  has  $\frac{dy}{dx} = -1$ . Therefore  $ds = \sqrt{1 + (-1)^2} dx = \sqrt{2} dx$ . The integral goes from  $x = 0$  to  $x = 1$  along the base. The length is  $\int_0^1 \sqrt{2} dx = \sqrt{2}$ . Check: The line from  $(1,0)$  to  $(0,1)$  certainly has length  $\sqrt{2}$ .
- The plane  $z = 1 - x - y$  has  $\frac{\partial z}{\partial x} = -1$  and  $\frac{\partial z}{\partial y} = -1$ . Therefore  $dS = \sqrt{1 + (-1)^2 + (-1)^2} dx dy = \sqrt{3} dx dy$ . **The integral is down in the  $xy$  plane!** The equilateral triangle in the sloping plane is over a right triangle in the base plane. Look only at the  $xy$  coordinates of the three corners:  $(1,0)$  and  $(0,1)$  and  $(0,0)$ . Those are the corners of the *projection* (the “shadow” down in the base). This shadow triangle has area  $\frac{1}{2}$ . The surface area above is:

$$\text{area of sloping plane} = \iint_{\text{shadow}} dS = \iint_{\text{base area}} \sqrt{3} dx dy = \sqrt{3} \cdot \frac{1}{2}.$$

Check: The sloping triangle has sides of length  $\sqrt{2}$ . That is the distance between its corners  $(1,0,0)$  and  $(0,1,0)$  and  $(0,0,1)$ . An equilateral triangle with sides  $\sqrt{2}$  has area  $\sqrt{3}/2$ .

**All these problems have three steps: Find  $dS$ . Find the shadow. Integrate  $dS$  over the shadow.**

2. Find the area on the plane  $x + 2y + z = 4$  which lies inside the vertical cylinder  $x^2 + y^2 = 1$ .

- The plane  $z = 4 - x - 2y$  has  $\frac{\partial z}{\partial x} = -1$  and  $\frac{\partial z}{\partial y} = -2$ . Therefore  $dS = \sqrt{1 + (-1)^2 + (-2)^2} dx dy = \sqrt{6} dx dy$ . The shadow in the base is the inside of the circle  $x^2 + y^2 = 1$ . This unit circle has area  $\pi$ . So the surface area on the sloping plane above it is  $\iint \sqrt{6} dx dy = \sqrt{6} \times \text{area of shadow} = \sqrt{6} \pi$ .

The region on that sloping plane is an *ellipse*. This is automatic when a plane cuts through a circular cylinder. The area of an ellipse is  $\pi ab$ , where  $a$  and  $b$  are the half-lengths of its axes. The axes of this ellipse are hard to find, so the new method that gave  $\text{area} = \sqrt{6}\pi$  is definitely superior.

3. Find the surface area on the sphere  $x^2 + y^2 + z^2 = 25$  between the horizontal planes  $z = 2$  and  $z = 4$ .

- The lower plane cuts the sphere in the circle  $x^2 + y^2 + 2^2 = 25$ . This is  $r^2 = 21$ . The upper plane cuts the sphere in the circle  $x^2 + y^2 + 4^2 = 25$ . This is  $r^2 = 9$ . **The shadow in the  $xy$  plane is the ring between  $r = 3$  and  $r = \sqrt{21}$ .**

The spherical surface has  $x^2 + y^2 + z^2 = 25$ . Therefore  $2x + 2z\frac{\partial z}{\partial x} = 0$  or  $\frac{\partial z}{\partial x} = -\frac{x}{z}$ . Similarly  $\frac{\partial z}{\partial y} = -\frac{y}{z}$ :

$$dS = \sqrt{1 + \left(-\frac{x}{z}\right)^2 + \left(-\frac{y}{z}\right)^2} dx dy = \frac{\sqrt{x^2 + y^2 + z^2}}{z} dx dy = \frac{5}{z} dx dy.$$

Remember  $z = \sqrt{25 - x^2 - y^2} = \sqrt{25 - r^2}$ . Integrate  $\frac{5}{z}$  over the shadow (the ring) using  $r$  and  $\theta$ :

$$\begin{aligned} \text{Surface area} &= \iint_{\text{ring}} \frac{5}{z} dx dy = \int_0^{2\pi} \int_3^{\sqrt{21}} \frac{5r dr d\theta}{\sqrt{25-r^2}} \\ &= (-5)(2\pi)\sqrt{25-r^2} \Big|_3^{\sqrt{21}} = -10\pi(2-4) = 20\pi. \end{aligned}$$

**Surface equations with parameters** Up to now the surface equation has been  $z = f(x, y)$ . This is restrictive. Each point  $(x, y)$  in the base has only one point above it in the surface. A complete sphere is not allowed. We solved a similar problem for curves, by allowing a parameter:  $x = \cos t$  and  $y = \sin t$  gave a complete circle. For surfaces we need *two parameters*  $u$  and  $v$ . Instead of  $x(t)$  we have  $x(u, v)$ . Similarly  $y = y(u, v)$  and  $z = z(u, v)$ . As  $u$  and  $v$  go over some region  $R$ , the points  $(x, y, z)$  go over the surface  $S$ .

For a circle, the parameter  $t$  is really the angle  $\theta$ . For a sphere, the parameters  $u$  and  $v$  are the angle  $\phi$  down from the North Pole and the angle  $\theta$  around the Equator. These are just spherical coordinates from Section 14.4:  $x = \sin u \cos v$  and  $y = \sin u \sin v$  and  $z = \cos u$ . In this case the region  $R$  is  $0 \leq u \leq \pi$  and  $0 \leq v \leq 2\pi$ . Then the points  $(x, y, z)$  cover the surface  $S$  of the unit sphere  $x^2 + y^2 + z^2 = 1$ .

We still have to find  $dS$ ! The general formula is equation (7) on page 575. For our  $x, y, z$  that equation gives  $dS = \sin u \, du \, dv$ . (In spherical coordinates you remember the volume element  $dV = \rho^2 \sin \phi \, d\rho \, d\phi \, d\theta$ . We are on the surface  $\rho = 1$ . And the letters  $\phi$  and  $\theta$  are changed to  $u$  and  $v$ .) This good formula for  $dS$  is typical of good coordinate systems – equation (7) is not as bad as it looks.

Integrate  $dS$  over the base region  $R$  in  $uv$  space to find the surface area above.

4. Find the surface area (known to be  $4\pi$ ) of the unit sphere  $x^2 + y^2 + z^2 = 1$ .

- Integrate  $dS$  over  $R$  to find  $\int_0^{2\pi} \int_0^\pi \sin u \, du \, dv = (2\pi)(-\cos u)|_0^\pi = 4\pi$ .

5. Recompute Question 3, the surface area on  $x^2 + y^2 + z^2 = 25$  between the planes  $z = 2$  and  $z = 4$ .

- This sphere has radius  $\sqrt{25} = 5$ . Multiply the points  $(x, y, z)$  on the unit sphere by 5:

$$x = 5 \sin u \cos v \text{ and } y = 5 \sin u \sin v \text{ and } z = 5 \cos u \text{ and } dS = 25 \sin u \, du \, dv.$$

Now find the region  $R$ . The angle  $v$  (or  $\theta$ ) goes around from 0 to  $2\pi$ . Since  $z = 5 \cos u$  goes from 2 to 4, the angle  $u$  is between  $\cos^{-1} \frac{2}{5}$  and  $\cos^{-1} \frac{4}{5}$ . Integrate  $dS$  over this region  $R$  and compare with  $20\pi$  above:

$$\text{Surface area} = \iint_R 25 \sin u \, du \, dv = 25(2\pi)(-\cos u) = 50\pi\left(\frac{4}{5} - \frac{2}{5}\right) = 20\pi.$$

6. Find the surface area of the cone  $z = 1 - \sqrt{x^2 + y^2}$  above the base plane  $z = 0$ .

- *Method 1:* Compute  $\frac{\partial z}{\partial x}$  and  $\frac{\partial z}{\partial y}$  and  $dS$ . Integrate over the shadow, a circle in the base plane. (Set  $z = 0$  to find the shadow boundary  $x^2 + y^2 = 1$ .) The integral takes 10 steps in Schaum's Outline. The answer is  $\pi\sqrt{2}$ .
- *Method 2:* Use parameters. Example 2 on page 575 gives  $x = u \cos v$  and  $y = u \sin v$  and  $z = u$  and  $dS = \sqrt{2}u \, du \, dv$ . The cone has  $0 \leq u \leq 1$  (since  $0 \leq z \leq 1$ ). The angle  $v$  (alias  $\theta$ ) goes from 0 to  $2\pi$ . This gives the parameter region  $R$  and we integrate in one step:

$$\text{cone area} = \int_0^{2\pi} \int_0^1 \sqrt{2} u \, du \, dv = (2\pi)(\sqrt{2})\left(\frac{1}{2}\right) = \pi\sqrt{2}.$$

The discussion of surface integrals ends with the calculation of *flow through a surface*. We are given the flow field  $\mathbf{F}(x, y, z)$  – a vector field with three components  $M(x, y, z)$  and  $N(x, y, z)$  and  $P(x, y, z)$ . The flow *through* a surface is  $\iint \mathbf{F} \cdot \mathbf{n} \, dS$ , where  $\mathbf{n}$  is the unit normal vector to the surface.

For the surface  $z = f(x, y)$ , you would expect a big square root for  $dS$ . It is there, but it is cancelled by a square root in  $\mathbf{n}$ . We divide the usual normal vector  $\mathbf{N} = -\frac{\partial f}{\partial x} \mathbf{i} - \frac{\partial f}{\partial y} \mathbf{j} + \mathbf{k}$  by its length to get the unit vector  $\mathbf{n} = \mathbf{N}/|\mathbf{N}|$ . That length is the square root that cancels. This leaves

$$\text{Flow through surface} = \iint_S \mathbf{F} \cdot \mathbf{n} \, dS = \iint_R \mathbf{F} \cdot \mathbf{N} \, dx \, dy = \iint_R \left(-M \frac{\partial f}{\partial x} - N \frac{\partial f}{\partial y} + P\right) \, dx \, dy.$$

The main job is to find the shadow region  $R$  in the  $xy$  plane and integrate. The “shadow” is the range of  $(x, y)$  down in the base, while  $(x, y, z)$  travels over the surface. With parameters  $u$  and  $v$ , the shadow region  $R$  is in the  $uv$  plane. It gives the range of parameters  $(u, v)$  as the point  $(x, y, z)$  travels over the surface  $S$ . This is not the easiest section in the book.

**Read-throughs and selected even-numbered solutions :**

A small piece of the surface  $z = f(x, y)$  is nearly flat. When we go across by  $dx$ , we go up by  $(\partial z/\partial x)dx$ . That movement is  $\mathbf{A}dx$ , where the vector  $\mathbf{A}$  is  $\mathbf{i} + dz/dx \mathbf{k}$ . The other side of the piece is  $\mathbf{B}dy$ , where  $\mathbf{B} = \mathbf{j} + (\partial z/\partial y)\mathbf{k}$ . The cross product  $\mathbf{A} \times \mathbf{B}$  is  $\mathbf{N} = -\partial z/\partial x \mathbf{i} - \partial z/\partial y \mathbf{j} + \mathbf{k}$ . The area of the piece is  $dS = |\mathbf{N}|dx \, dy$ . For the surface  $z = xy$ , the vectors are  $\mathbf{A} = \mathbf{i} + y \mathbf{k}$  and  $\mathbf{B} = \mathbf{j} + x \mathbf{k}$ . The area integral is  $\iint dS = \iint \sqrt{1 + y^2 + x^2} \, dx \, dy$ .

With parameters  $u$  and  $v$ , a typical point on a  $45^\circ$  cone is  $x = u \cos v, y = u \sin v, z = u$ . A change in  $u$  moves that point by  $\mathbf{A} \, du = (\cos v \mathbf{i} + \sin v \mathbf{j} + \mathbf{k})du$ . The change in  $v$  moves the point by  $\mathbf{B} \, dv = (-u \sin v \mathbf{i} + u \cos v \mathbf{j})dv$ . The normal vector is  $\mathbf{N} = \mathbf{A} \times \mathbf{B} = -u \cos v \mathbf{i} - u \sin v \mathbf{j} + u \mathbf{k}$ . The area is  $dS = \sqrt{2} \, u \, du \, dv$ . In this example  $\mathbf{A} \cdot \mathbf{B} = 0$  so the small piece is a rectangle and  $dS = |\mathbf{A}||\mathbf{B}| \, du \, dv$ .

For flux we need  $\mathbf{n}dS$ . The unit normal vector  $\mathbf{n}$  is  $\mathbf{N} = \mathbf{A} \times \mathbf{B}$  divided by  $|\mathbf{N}|$ . For a surface  $z = f(x, y)$ , the product  $\mathbf{n}dS$  is the vector  $\mathbf{N} \, dx \, dy$  (to memorize from table). The particular surface  $z = xy$  has  $\mathbf{n}dS = (-y\mathbf{i} - x\mathbf{j} + \mathbf{k})dx \, dy$ . For  $\mathbf{F} = x\mathbf{i} + y\mathbf{j} + z\mathbf{k}$  the flux through  $z = xy$  is  $\mathbf{F} \cdot \mathbf{n}dS = -xy \, dx \, dy$ .

On a  $30^\circ$  cone the points are  $x = 2u \cos v, y = 2u \sin v, z = u$ . The tangent vectors are  $\mathbf{A} = 2 \cos v \mathbf{i} + 2 \sin v \mathbf{j} + \mathbf{k}$  and  $\mathbf{B} = -2u \sin v \mathbf{i} + 2u \cos v \mathbf{j}$ . This cone has  $\mathbf{n}dS = \mathbf{A} \times \mathbf{B} \, du \, dv = (-2u \cos v \mathbf{i} - 2u \sin v \mathbf{j} + 4u \mathbf{k})du \, dv$ . For  $\mathbf{F} = x \mathbf{i} + y \mathbf{j} + z \mathbf{k}$ , the flux element through the cone is  $\mathbf{F} \cdot \mathbf{n}dS = \text{zero}$ . The reason for this answer is that  $\mathbf{F}$  is along the cone. The reason we don't compute flux through a Möbius strip is that  $\mathbf{N}$  cannot be defined (the strip is not orientable).

2  $\mathbf{N} = -2x \mathbf{i} - 2y \mathbf{j} + \mathbf{k}$  and  $dS = \sqrt{1 + 4x^2 + 4y^2} \, dx \, dy$ . Then  $\iint dS = \int_0^{2\pi} \int_2^{\sqrt{8}} \sqrt{1 + 4r^2} \, r \, dr \, d\theta = \frac{\pi}{6}(33^{3/2} - 17^{3/2})$ .

8  $\mathbf{N} = -\frac{x}{r} \mathbf{i} - \frac{y}{r} \mathbf{j} + \mathbf{k}$  and  $dS = \frac{x^2 + y^2 + r^2}{r^2} \, dx \, dy = \sqrt{2} \, dx \, dy$ . Then area  $= \int_0^{2\pi} \int_a^b \sqrt{2} \, r \, dr \, d\theta = \sqrt{2}\pi(b^2 - a^2)$ .

16 On the sphere  $dS = \sin \phi \, d\phi \, d\theta$  and  $g = x^2 + y^2 = \sin^2 \phi$ . Then  $\int_0^{2\pi} \int_0^{\pi/2} \sin^3 \phi \, d\phi \, d\theta = 2\pi(\frac{2}{3}) = \frac{4\pi}{3}$ .

20  $\mathbf{A} = v\mathbf{i} + \mathbf{j} + \mathbf{k}, \mathbf{B} = u\mathbf{i} + \mathbf{j} - \mathbf{k}, \mathbf{N} = \mathbf{A} \times \mathbf{B} = -2\mathbf{i} + (u+v)\mathbf{j} + (v-u)\mathbf{k}, dS = \sqrt{4 + 2u^2 + 2v^2} \, du \, dv$ .

24  $\iint \mathbf{F} \cdot \mathbf{n}dS = \int_0^{2\pi} \int_2^{\sqrt{8}} -r^3 \, dr \, d\theta = -24\pi$ .

30  $\mathbf{A} = \cos \theta \mathbf{i} + \sin \theta \mathbf{j} - 2r\mathbf{k}, \mathbf{B} = -r \sin \theta \mathbf{i} + r \cos \theta \mathbf{j}, \mathbf{N} = \mathbf{A} \times \mathbf{B} = 2r^2 \cos \theta \mathbf{i} + 2r^2 \sin \theta \mathbf{j} + r\mathbf{k}$ ,

$\iint \mathbf{k} \cdot \mathbf{n} \, dS = \iint \mathbf{k} \cdot \mathbf{N} \, du \, dv = \int_0^{2\pi} \int_0^a r \, dr \, d\theta = \pi a^2$  as in Example 12.

## 15.5 The Divergence Theorem (page 588)

This theorem says that the total source inside a volume  $V$  equals the total flow through its closed surface  $S$ . We need to know how to measure the *source* and how to measure the *flow out*. Both come from the vector field  $\mathbf{F}(x, y, z) = M\mathbf{i} + N\mathbf{j} + P\mathbf{k}$  which assigns a flow vector to every point inside  $V$  and on  $S$ :

$$\text{Source} = \text{divergence of } \mathbf{F} = \nabla \cdot \mathbf{F} = \frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} + \frac{\partial P}{\partial z}$$

$$\text{Flow out} = \text{normal component of } \mathbf{F} = \mathbf{F} \cdot \mathbf{n} = \mathbf{F} \cdot \mathbf{N}/|\mathbf{N}|$$

The balance between source and outward flow is the *Divergence Theorem*. It is Green's Theorem in 3 dimensions:

$$\iint_S \mathbf{F} \cdot \mathbf{n} \, dS = \iiint_V (\text{divergence of } \mathbf{F}) \, dV = \iiint_V \left( \frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} + \frac{\partial P}{\partial z} \right) dx \, dy \, dz.$$

An important special case is a "source-free" field. This means that the divergence is zero. Then the integrals are zero and the total outward flow is zero. There may be flow out through one part of  $S$  and flow in through another part – they must cancel when  $\text{div } \mathbf{F} = 0$  inside  $V$ .

1. For the field  $\mathbf{F} = y\mathbf{i} + x\mathbf{j}$  in the unit ball  $x^2 + y^2 + z^2 \leq 1$ , compute both sides of the Divergence Theorem.

- The divergence is  $\frac{\partial y}{\partial x} + \frac{\partial x}{\partial y} + \frac{\partial 0}{\partial z} = 0$ . This source-free field has  $\iiint \text{div } \mathbf{F} \, dV = 0$ .
- The normal vector to the unit sphere is radially outward:  $\mathbf{n} = x\mathbf{i} + y\mathbf{j} + z\mathbf{k}$ . Its dot product with  $\mathbf{F} = y\mathbf{i} + x\mathbf{j}$  gives  $\mathbf{F} \cdot \mathbf{n} = 2xy$  for flow out through  $S$ . This is not zero, but *its integral is zero*. One proof is by symmetry:  $2xy$  is equally positive and negative on the sphere. The direct proof is by integration (use spherical coordinates):

$$\iint 2xy \, dS = \int_0^{2\pi} \int_0^\pi 2(\sin \phi \cos \theta)(\sin \phi \sin \theta) \sin \phi \, d\phi \, d\theta = 0 \text{ because } \int_0^{2\pi} 2 \cos \theta \sin \theta \, d\theta = 0.$$

*To emphasize:* The flow out of *any volume*  $V$  is zero because this field has divergence = 0. For strange shapes we can't do the surface integral. But the volume integral is still  $\iiint 0 \, dV$ .

2. (This is Problem 15.5.8) Find the divergence of  $\mathbf{F} = x^3\mathbf{i} + y^3\mathbf{j} + z^3\mathbf{k}$  and the flow out of the sphere  $\rho = a$ .

- Divergence =  $\frac{\partial x^3}{\partial x} + \frac{\partial y^3}{\partial y} + \frac{\partial z^3}{\partial z} = 3x^2 + 3y^2 + 3z^2 = 3\rho^2$ . The triple integral for the total source is

$$\int_0^{2\pi} \int_0^\pi \int_0^a 3\rho^2(\rho^2 \sin \phi \, d\rho \, d\phi \, d\theta) = 3\left(\frac{a}{5}\right)^5(2)(2\pi) = 12\pi a^5/5.$$

- Flow out has  $\mathbf{F} \cdot \mathbf{n} = (x^3, y^3, z^3) \cdot \left(\frac{x}{a}, \frac{y}{a}, \frac{z}{a}\right) = \frac{x^4 + y^4 + z^4}{a}$ . The integral of  $\frac{x^4}{a}$  over the sphere  $\rho = a$  is

$$\int_0^{2\pi} \int_0^\pi \frac{1}{a}(a \cos \phi)^4(a^2 \sin \phi \, d\phi \, d\theta) = -2\pi a^5 \frac{\cos^5 \phi}{5} \Big|_0^\pi = \frac{4\pi a^5}{5}.$$

By symmetry  $\frac{x^4}{a}$  and  $\frac{y^4}{a}$  have this same integral. So multiply by 3 to get the same  $12\pi a^5/5$  as above.

The text assigns special importance to the vector field  $\mathbf{F} = \mathbf{R}/\rho^3$ . This is radially outward (remember  $\mathbf{R} = x\mathbf{i} + y\mathbf{j} + z\mathbf{k}$ ). The length of  $\mathbf{F}$  is  $|\mathbf{R}|/\rho^3 = \rho/\rho^3 = 1/\rho^2$ . This is the *inverse-square law* – the force of gravity from a point mass at the origin decreases like  $1/\rho^2$ .

The special feature of this radial field is to have zero divergence – except at one point. This is not typical of radial fields:  $\mathbf{R}$  itself has divergence  $\frac{\partial x}{\partial x} + \frac{\partial y}{\partial y} + \frac{\partial z}{\partial z} = 3$ . But dividing  $\mathbf{R}$  by  $\rho^3$  gives a field with  $\text{div } \mathbf{F} = 0$ . (Physically: No divergence where there is no mass and no source of gravity.) The exceptional point is  $(0,0,0)$ , where there is a mass. That point source is enough to produce  $4\pi$  on both sides of the divergence theorem (provided  $S$  encloses the origin). The point source has strength  $4\pi$ . The divergence of  $\mathbf{F}$  is  $4\pi$  times a “delta function.”

The other topic in this section is the vector form of two familiar rules: the *product rule* for the derivative of  $u(x)v(x)$  and the reverse of the product rule which is *integration by parts*. Now we are in 2 or 3 dimensions and  $v$  is a vector field  $\mathbf{V}(x, y, z)$ . The derivative is replaced by the divergence or the curl. We just use the old product rule on  $uM$  and  $uN$  and  $uP$ . Then collect terms:

$$\text{div}(u\mathbf{V}) = u \text{div} \mathbf{V} + (\text{grad } u) \cdot \mathbf{V} \quad \text{curl}(u\mathbf{V}) = u \text{curl } \mathbf{V} + (\text{grad } u) \times \mathbf{V}.$$

**Read-throughs and selected even-numbered solutions :**

In words, the basic balance law is **flow in = flow out**. The flux of  $\mathbf{F}$  through a surface  $S$  is the double integral  $\iint \mathbf{F} \cdot \mathbf{n} dS$ . The divergence of  $M\mathbf{i} + N\mathbf{j} + P\mathbf{k}$  is  $M_x + N_y + P_z$ . It measures the **source at the point**. The total source is the triple integral  $\iiint \text{div } \mathbf{F} dV$ . That equals the flux by the **Divergence Theorem**.

For  $\mathbf{F} = 5z\mathbf{k}$  the divergence is 5. If  $V$  is a cube of side  $a$  then the triple integral equals  $5a^3$ . The top surface where  $z = a$  has  $\mathbf{n} = \mathbf{k}$  and  $\mathbf{F} \cdot \mathbf{n} = 5a$ . The bottom and sides have  $\mathbf{F} \cdot \mathbf{n} = \text{zero}$ . The integral  $\iint \mathbf{F} \cdot \mathbf{n} dS = 5a^3$ .

The field  $\mathbf{F} = \mathbf{R}/\rho^3$  has  $\text{div } \mathbf{F} = 0$  except at the origin.  $\iint \mathbf{F} \cdot \mathbf{n} dS$  equals  $4\pi$  over any surface around the origin. This illustrates Gauss’s Law: **flux =  $4\pi$  times source strength**. The field  $\mathbf{F} = x\mathbf{i} + y\mathbf{j} - 2z\mathbf{k}$  has  $\text{div } \mathbf{F} = 0$  and  $\iint \mathbf{F} \cdot \mathbf{n} dS = 0$ . For this  $\mathbf{F}$ , the flux out through a pyramid and in through its base are equal.

The symbol  $\nabla$  stands for  $(\partial/\partial x)\mathbf{i} + (\partial/\partial y)\mathbf{j} + (\partial/\partial z)\mathbf{k}$ . In this notation  $\text{div } \mathbf{F}$  is  $\nabla \cdot \mathbf{F}$ . The gradient of  $f$  is  $\nabla f$ . The divergence of  $\text{grad } f$  is  $\nabla \cdot \nabla f$  or  $\nabla^2 f$ . The equation  $\text{div grad } f = 0$  is **Laplace’s equation**.

The divergence of a product is  $\text{div}(u\mathbf{V}) = u \text{div } \mathbf{V} + (\text{grad } u) \cdot \mathbf{V}$ . Integration by parts in 3D is  $\iiint u \text{div } \mathbf{V} dx dy dz = -\iiint \mathbf{V} \cdot \text{grad } u dx dy dz + \iint u \mathbf{V} \cdot \mathbf{n} dS$ . In two dimensions this becomes  $\iint u(\partial M/\partial x + \partial N/\partial y) dx dy = -\int (M \partial u/\partial x + N \partial u/\partial y) dx dy + \int u \mathbf{V} \cdot \mathbf{n} ds$ . In one dimension it becomes **integration by parts**. For steady fluid flow the continuity equation is  $\text{div } \rho \mathbf{V} = -\partial \rho/\partial t$ .

**14**  $\mathbf{R} \cdot \mathbf{n} = (x\mathbf{i} + y\mathbf{j} + z\mathbf{k}) \cdot \mathbf{i} = x = 1$  on one face of the box. On the five other faces  $\mathbf{R} \cdot \mathbf{n} = 2, 3, 0, 0, 0$ .

The integral is  $\int_0^3 \int_0^2 1 dy dz + \int_0^3 \int_0^1 2 dx dz + \int_0^2 \int_0^1 3 dx dy = 18$ . Also  $\text{div } \mathbf{R} = 1 + 1 + 1 = 3$  and  $\int_0^3 \int_0^2 \int_0^1 3 dx dy dz = 18$ .

**18**  $\text{grad } f \cdot \mathbf{n}$  is the **directional derivative** in the normal direction  $\mathbf{n}$  (also written  $\frac{\partial f}{\partial \mathbf{n}}$ ).

The Divergence Theorem gives  $\iiint \text{div}(\text{grad } f) dV = \iint \text{grad } f \cdot \mathbf{n} dS = \iint \frac{\partial f}{\partial \mathbf{n}} dS$ .

But we are given that  $\text{div}(\text{grad } f) = f_{xx} + f_{yy} + f_{zz}$  is zero.

**26** When the density  $\rho$  is constant (incompressible flow), the continuity equation becomes  $\text{div } \mathbf{V} = 0$ . If the flow is irrotational then  $\mathbf{F} = \text{grad } f$  and the continuity equation is  $\text{div}(\rho \text{grad } f) = -d\rho/dt$ .

If also  $\rho = \text{constant}$ , then  $\text{div grad } f = 0$ : Laplace’s equation for the “potential.”

**30** The boundary of a solid ball is a sphere. A sphere has no boundary. Similarly for a cube or a cylinder – the boundary is a closed surface and **that surface's boundary** is empty. This is a crucial fact in topology.

## 15.6 Stokes' Theorem and the Curl of $\mathbf{F}$ (page 595)

The curl of  $\mathbf{F}$  measures the “spin”. A spin field in the plane is  $\mathbf{S} = -y\mathbf{i} + x\mathbf{j}$ , with third component  $P = 0$ :

$$\text{curl } \mathbf{S} = \left(\frac{\partial P}{\partial y} - \frac{\partial N}{\partial z}\right)\mathbf{i} + \left(\frac{\partial M}{\partial z} - \frac{\partial P}{\partial x}\right)\mathbf{j} + \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y}\right)\mathbf{k} = 0\mathbf{i} + 0\mathbf{j} + (1 + 1)\mathbf{k}.$$

The curl is  $2\mathbf{k}$ . It points along the spin axis (the  $z$  axis, perpendicular to the plane of spin). Its magnitude is 2 times the rotation rate. This special spin field  $\mathbf{S}$  gives a rotation counterclockwise in the  $xy$  plane. It is counterclockwise because  $\mathbf{S}$  points that way, and also because of the right hand rule: thumb in the direction of curl  $\mathbf{F}$  and fingers “curled in the direction of spin.” Put your right hand on a table with thumb upward along  $\mathbf{k}$ .

Spin fields can go around any axis vector  $\mathbf{a}$ . The field  $\mathbf{S} = \mathbf{a} \times \mathbf{R}$  does that. Its curl is  $2\mathbf{a}$  (after calculation). Other fields have a curl that changes direction from point to point. Some fields have no spin at all, so their curl is zero. *These are gradient fields!* This is a key fact:

$$\text{curl } \mathbf{F} = \mathbf{0} \text{ whenever } \mathbf{F} = \text{grad } f = \frac{\partial f}{\partial x}\mathbf{i} + \frac{\partial f}{\partial y}\mathbf{j} + \frac{\partial f}{\partial z}\mathbf{k}.$$

You should substitute those three partial derivatives for  $M, N, P$  and see that the curl formula gives zero. **The curl of a gradient is zero.** The quick test **D** for a gradient field is  $\text{curl } \mathbf{F} = \mathbf{0}$ .

The twin formula is that *the divergence of a curl is zero*. The quick test **H** for a source-free “curl field” is  $\text{div } \mathbf{F} = 0$ . A gradient field can be a gravity field or an electric field. A curl field can be a magnetic field.

1. Show that  $\mathbf{F} = yz\mathbf{i} + xy\mathbf{j} + (xy + 2z)\mathbf{k}$  passes the quick test **D** for a gradient field (a conservative field).

The test is  $\text{curl } \mathbf{F} = \mathbf{0}$ . Find the potential function  $f$  that this test guarantees:  $\mathbf{F}$  equals the gradient of  $f$ .

- The curl of this  $\mathbf{F}$  is  $\left(\frac{\partial P}{\partial y} - \frac{\partial N}{\partial z}\right)\mathbf{i} + \left(\frac{\partial M}{\partial z} - \frac{\partial P}{\partial x}\right)\mathbf{j} + \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y}\right)\mathbf{k} = (x - x)\mathbf{i} + (y - y)\mathbf{j} + (z - z)\mathbf{k} = \mathbf{0}$ . The field passes test **D**. There must be a function whose partial derivatives are  $M, N, P$ :  $\frac{\partial f}{\partial x} = yz$  and  $\frac{\partial f}{\partial y} = xz$  and  $\frac{\partial f}{\partial z} = xy + 2z$  lead to the potential function  $f = xyz + z^2$ .

We end with **Stokes' Theorem**. It is like the original Green's Theorem, where work around a plane curve  $C$  was equal to the double integral of  $\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y}$ . Now the curve  $C$  can go out of the plane. The region inside is a curved surface – also not in a plane. The field  $\mathbf{F}$  is three-dimensional – its component  $P\mathbf{k}$  goes out of the plane. Stokes' Theorem has a line integral  $\int_C \mathbf{F} \cdot d\mathbf{R}$  (work around  $C$ ) equal to a surface integral:

$$\int_C M dx + N dy + P dz = \iint_S (\text{curl } \mathbf{F}) \cdot \mathbf{n} dS = \iint_{\text{base}} (\text{curl } \mathbf{F}) \cdot \mathbf{N} dx dy.$$

If the surface is  $z = f(x, y)$  then its normal is  $\mathbf{N} = -\frac{\partial z}{\partial x}\mathbf{i} - \frac{\partial z}{\partial y}\mathbf{j} + \mathbf{k}$  and  $dz$  is  $\frac{\partial z}{\partial x}dx + \frac{\partial z}{\partial y}dy$ . Substituting for  $\mathbf{N}$  and  $dz$  reduces the 3-dimensional theorem of Stokes to the 2-dimensional theorem of Green.

2. (Compare 15.6.12) For  $\mathbf{F} = \mathbf{i} \times \mathbf{R}$  compute both sides in Stokes' Theorem when  $C$  is the unit circle.

- The cross product  $\mathbf{F} = \mathbf{i} \times \mathbf{R}$  is  $\mathbf{i} \times (x\mathbf{i} + y\mathbf{j} + z\mathbf{k}) = y\mathbf{k} - z\mathbf{j}$ . This spin field has  $\text{curl } \mathbf{F} = 2\mathbf{i}$ . either substitute  $M = 0, N = -z, P = y$  in the curl formula or remember that  $\text{curl}(\mathbf{a} \times \mathbf{R}) = 2\mathbf{a}$ . Here  $\mathbf{a} = \mathbf{i}$ .

The line integral of  $\mathbf{F} = y\mathbf{k} - z\mathbf{j}$  around the unit circle is  $\int 0 dx - z dy + y dz$ . All those are zero because  $z = 0$  for the circle  $x^2 + y^2 = 1$  in the  $xy$  plane. **The left side of Stokes' Theorem is zero for this  $\mathbf{F}$ .**

The double integral of  $(\text{curl } \mathbf{F}) \cdot \mathbf{n} = 2\mathbf{i} \cdot \mathbf{n}$  is certainly zero if the surface  $S$  is the flat disk inside the unit circle. The normal vector to that flat surface is  $\mathbf{n} = \mathbf{k}$ . Then  $2\mathbf{i} \cdot \mathbf{n}$  is  $2\mathbf{i} \cdot \mathbf{k} = 0$ .

The double integral of  $2\mathbf{i} \cdot \mathbf{n}$  is zero even if the surface  $S$  is not flat.  $S$  can be a mountain (always with the unit circle as its base boundary). The normal  $\mathbf{n}$  out from the mountain can have an  $\mathbf{i}$  component, so  $2\mathbf{i} \cdot \mathbf{n}$  can be non-zero. But Stokes' Theorem says: The *integral* of  $2\mathbf{i} \cdot \mathbf{n}$  over the mountain is zero. That is the theorem. The mountain integral must be zero because the baseline integral is zero.

**Read-throughs and selected even-numbered solutions :**

The curl of  $M\mathbf{i} + N\mathbf{j} + P\mathbf{k}$  is the vector  $(P_y - N_z)\mathbf{i} + (M_z - P_x)\mathbf{j} + (N_x - M_y)\mathbf{k}$ . It equals the 3 by 3 determinant  $\begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ M & N & P \end{vmatrix}$ . The curl of  $x^2\mathbf{i} + z^2\mathbf{k}$  is zero. For  $\mathbf{S} = y\mathbf{i} - (x+z)\mathbf{j} + y\mathbf{k}$  the curl is  $2\mathbf{i} - 2\mathbf{k}$ . This  $\mathbf{S}$  is a spin field  $\mathbf{a} \times \mathbf{R} = \frac{1}{2}(\text{curl } \mathbf{F}) \times \mathbf{R}$ , with axis vector  $\mathbf{a} = \mathbf{i} - \mathbf{k}$ . For any gradient field  $f_x\mathbf{i} + f_y\mathbf{j} + f_z\mathbf{k}$  the curl is zero. That is the important identity  $\text{curl grad } f = \mathbf{zero}$ . It is based on  $f_{xy} = f_{yx}$  and  $f_{xz} = f_{zx}$  and  $f_{yz} = f_{zy}$ . The twin identity is  $\text{div curl } \mathbf{F} = 0$ .

The curl measures the spin (or turning) of a vector field. A paddlewheel in the field with its axis along  $\mathbf{n}$  has turning speed  $\frac{1}{2}\mathbf{n} \cdot \text{curl } \mathbf{F}$ . The spin is greatest when  $\mathbf{n}$  is in the direction of  $\text{curl } \mathbf{F}$ . Then the angular velocity is  $\frac{1}{2}|\text{curl } \mathbf{F}|$ . Stokes' Theorem is  $\oint_C \mathbf{F} \cdot d\mathbf{R} = \iint_S (\text{curl } \mathbf{F}) \cdot \mathbf{n} dS$ . The curve  $C$  is the boundary of the surface  $S$ . This is Green's Theorem extended to three dimensions. Both sides are zero when  $\mathbf{F}$  is a gradient field because the curl is zero.

The four properties of a conservative field are A :  $\oint \mathbf{F} \cdot d\mathbf{R} = 0$  and B :  $\int_P^Q \mathbf{F} \cdot d\mathbf{R}$  depends only on P and Q and C :  $\mathbf{F}$  is the gradient of a potential function  $f(x, y, z)$  and D :  $\text{curl } \mathbf{F} = 0$ . The field  $y^2z^2\mathbf{i} + 2xy^2z\mathbf{k}$  fails test D. This field is the gradient of no  $f$ . The work  $\int \mathbf{F} \cdot d\mathbf{R}$  from  $(0,0,0)$  to  $(1,1,1)$  is  $\frac{3}{5}$  along the straight path  $x = y = z = t$ . For every field  $\mathbf{F}$ ,  $\iint \text{curl } \mathbf{F} \cdot \mathbf{n} dS$  is the same out through a pyramid and up through its base because they have the same boundary, so  $\oint \mathbf{F} \cdot d\mathbf{R}$  is the same.

**14**  $\mathbf{F} = (x^2 + y^2)\mathbf{k}$  so  $\text{curl } \mathbf{F} = 2(y\mathbf{i} - x\mathbf{j})$ . (Surprise that this  $\mathbf{F} = \mathbf{a} \times \mathbf{R}$  has  $\text{curl } \mathbf{F} = 2\mathbf{a}$  even with nonconstant  $\mathbf{a}$ .) Then  $\oint \mathbf{F} \cdot d\mathbf{R} = \iint \text{curl } \mathbf{F} \cdot \mathbf{n} dS = 0$  since  $\mathbf{n} = \mathbf{k}$  is perpendicular to  $\text{curl } \mathbf{F}$ .

**18** If  $\text{curl } \mathbf{F} = 0$  then  $\mathbf{F}$  is the gradient of a potential:  $\mathbf{F} = \text{grad } f$ . Then  $\text{div } \mathbf{F} = 0$  is  $\text{div grad } f = 0$  which is Laplace's equation.

**24** Start with one field that has the required curl. (Can take  $\mathbf{F} = \frac{1}{2}\mathbf{i} \times \mathbf{R} = -\frac{z}{2}\mathbf{j} + \frac{y}{2}\mathbf{k}$ ). Then add any  $\mathbf{F}$  with curl zero (particular solution plus homogeneous solution as always). The fields with  $\text{curl } \mathbf{F} = 0$  are gradient fields  $\mathbf{F} = \text{grad } f$ , since  $\text{curl grad} = 0$ . Answer:  $\mathbf{F} = \frac{1}{2}\mathbf{i} \times \mathbf{R} + \text{any grad } f$ .

**26**  $\mathbf{F} = y\mathbf{i} - x\mathbf{k}$  has  $\text{curl } \mathbf{F} = \mathbf{j} - \mathbf{k}$ . (a) Angular velocity =  $\frac{1}{2} \text{curl } \mathbf{F} \cdot \mathbf{n} = \frac{1}{2}$  if  $\mathbf{n} = \mathbf{j}$ .

(b) Angular velocity =  $\frac{1}{2}|\text{curl } \mathbf{F}| = \frac{\sqrt{2}}{2}$  (c) Angular velocity = 0.

**36**  $\text{curl } \mathbf{F} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ z & x & xyz \end{vmatrix} = \mathbf{i}(xz) + \mathbf{j}(1 - yz) + \mathbf{k}(1)$  and  $\mathbf{n} = x\mathbf{i} + y\mathbf{j} + z\mathbf{k}$ . So  $\text{curl } \mathbf{F} \cdot \mathbf{n} =$

$x^2z + y - y^2z + z$ . By symmetry  $\iint x^2z dS = \iint y^2z dS$  on the half sphere and  $\iint yz dS = 0$ .

This leaves  $\iint z dS = \int_0^{2\pi} \int_0^{\pi/2} \cos \phi (\sin \phi d\phi d\theta) = \frac{1}{2}(2\pi) = \pi$ .

38 (The expected method is trial and error)  $\mathbf{F} = 5yzi + 2xy\mathbf{k} + \text{any grad } f$ .

## 15 Chapter Review Problems

### Review Problems

- R1** For  $f(x, y) = x^2 + y^2$  what is the gradient field  $\mathbf{F} = \nabla f$ ? What is the unit field  $\mathbf{u} = \mathbf{F}/|\mathbf{F}|$ ?
- R2** Is  $\mathbf{F}(x, y) = \cos x \mathbf{i} + \sin x \mathbf{j}$  a gradient field? Draw the vectors  $\mathbf{F}(0,0)$ ,  $\mathbf{F}(0, \pi)$ ,  $\mathbf{F}(\pi, 0)$ ,  $\mathbf{F}(\pi, \pi)$ .
- R3** Is  $\mathbf{F}(x, y) = \cos y \mathbf{i} + \sin x \mathbf{j}$  a gradient field or not? Draw  $\mathbf{F}(0,0)$ ,  $\mathbf{F}(0, \pi)$ ,  $\mathbf{F}(\pi, 0)$ ,  $\mathbf{F}(\pi, \pi)$ .
- R4** Is  $\mathbf{F}(x, y) = \cos x \mathbf{i} + \sin y \mathbf{j}$  a gradient field? If so find a potential  $f(x, y)$  whose gradient is  $\mathbf{F}$ .
- R5** Integrate  $-y dx + x dy$  around the unit circle  $x = \cos t, y = \sin t$ . Why twice the area?
- R6** Integrate the gradient of  $f(x, y) = x^3 + y^3$  around the unit circle.
- R7** With Green's Theorem find an integral around  $C$  that gives the area inside  $C$ .
- R8** Find the flux of  $\mathbf{F} = x^2\mathbf{i}$  through the unit circle from both sides of Green's Theorem.
- R9** What integral gives the area of the surface  $z = f(x, y)$  above the square  $|x| \leq 1, |y| \leq 1$ ?
- R10** Describe the cylinder given by  $x = 2 \cos v, y = 2 \sin v$ , and  $z = u$ . Is  $(2, 2, 2)$  on the cylinder? What parameters  $u$  and  $v$  produce the point  $(0, 2, 4)$ ?

### Drill Problems

- D1** Find the gradient  $\mathbf{F}$  of  $f(x, y, z) = x^3 + y^3$ . Then find the divergence and curl of  $\mathbf{F}$ .
- D2** What integral gives the surface area of  $z = 1 - x^2 - y^2$  above the  $xy$  plane?
- D3** What is the area on the sloping plane  $z = x + y$  above a base area (or shadow area)  $A$ ?
- D4** Write down Green's Theorem for  $\int Mdx + Ndy = \text{work}$  and  $\int Mdy - Ndx = \text{flux}$ . Write down the flux form with  $Mdy - Ndx$ .
- D5** Write down the Divergence Theorem. Say in words what it balances.
- D6** If  $\mathbf{F} = y^2 \mathbf{i} + (1 + 2x)y \mathbf{j}$  is the gradient of  $f(x, y)$ , find the potential function  $f$ .
- D7** When is the area of a surface equal to the area of its shadow on  $z = 0$ ? (Surface is  $z = f(x, y)$ ).
- D8** For  $\mathbf{F} = 3xi + 4yj + 5k$ , the flux  $\iint \mathbf{F} \cdot \mathbf{n}dS$  equals \_\_\_ times the volume inside  $S$ .
- D9** What vector field is the curl of  $\mathbf{F} = xyz \mathbf{i}$ ? Find the gradient of that curl.
- D10** What vector field is the gradient of  $f = xyz$ ? Find the curl of that gradient.

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