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8.512 Theory of Solids II  
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Consider a Fermi gas with dispersion  $\epsilon_k$  and a repulsive interaction  $U\delta(\mathbf{r})$ . Now if  $N(0)U > 1$ , we find in mean field theory the spontaneous appearance of the order parameter:

$$\Delta = U\langle n_\uparrow - n_\downarrow \rangle ,$$

and the splitting of the up- and down-spin bands,

$$\begin{aligned}\tilde{\epsilon}_{k\uparrow} &= \epsilon_k - \Delta/2 \\ \tilde{\epsilon}_{k\downarrow} &= \epsilon_k + \Delta/2 .\end{aligned}$$

1. Show that the **transverse** spin susceptibility of a system described by the mean field Hamiltonian is given by  $\chi_\perp^0 = \mu_B^2 \Gamma_0$  where

$$\Gamma_0(q, \omega) = \sum_k \frac{f(\tilde{\epsilon}_{k+q,\uparrow}) - f(\tilde{\epsilon}_{k,\downarrow})}{\omega - \tilde{\epsilon}_{k+q\uparrow} + \tilde{\epsilon}_{k+q\downarrow} + i\eta} .$$

This is the generalization of the Lindhard function to a spin split band.

2. Now include the interaction term in the response to the transverse field in a self consistent field approximation. Show that

$$\chi_\perp(q, \omega) = \frac{\mu_B^2 \Gamma_0(q, \omega)}{1 - U\Gamma_0(q, \omega)} .$$

3. The poles of the numerator in  $\chi_\perp$  describe the single particle-hole excitations. Sketch the region in  $(\omega, q)$  space where  $Im\chi_\perp \neq 0$  due to these excitations.
4. The other pole in  $\chi_\perp(q, \omega)$  occurs when the denominator vanishes. Calculate the dispersion of this pole which we identify as the spin wave excitation as follows:

- (a) Show that at  $q = \omega = 0$ , the denominator vanishes. [Hint: the condition  $1 - U\Gamma_0 = 0$  is the same as the self-consistency equation for  $\Delta(T)$ .]

- (b) Expand  $\Gamma_0(q, \omega)$  for small  $q, \omega$  and show that the location of the pole of  $\chi_\perp$  is given by  $\omega(q) = Dq^2$ . Note that unlike the Lindhard function for free fermions, the existence of the gap  $\Delta$  makes the expansion well behaved.