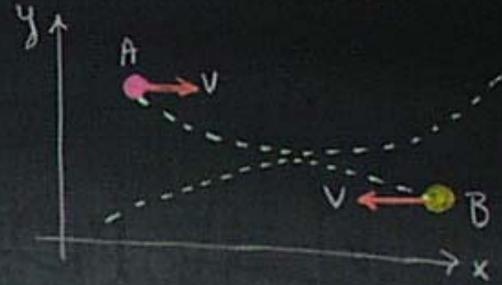


Relativistic Momentum

- What does Special Relativity say about Dynamics?
 - How do we preserve Cons. of Momentum?
- Elastic Collision of two identical particles.
- Frame-A moves along x-axis fixed to A.
- Frame-B moves along x-axis fixed to B.

- Collision completely symmetric
- Each particle has speed u_0 along its y-axis.
- Collision changes y-velocities leaves x-motion constant.
- Relative velocity of frames is V



Law of Trans. of Velocities for opposite particle:

$$\frac{u_0}{\gamma} = u_0 \sqrt{1 - v^2/c^2}$$

• After collision y-velocities reversed.

Assume $\vec{p} = m(w) \vec{w}$

$m(w)$: something that may depend on particle speed \vec{w} .

Frame-A:

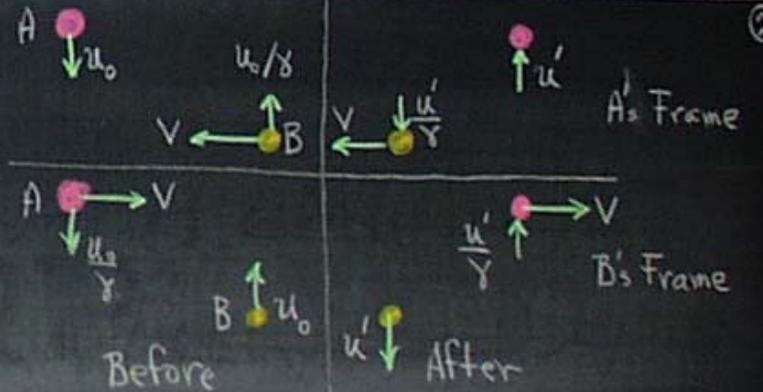
x-momentum only due to particle-B

Before collision B's speed

$$w^2 = V^2 + \frac{u_0^2}{\gamma^2}$$

After collision

$$w^2 = V^2 + \frac{u_1^2}{\gamma^2}$$



Law of Trans of Velocities for opposite particle:
 $\frac{u_0}{\gamma} = u_0 \sqrt{1 - v^2/c^2}$

After collision y-velocities reversed.

Assume $\vec{p} = m(u) \vec{u}$

$m(u)$: something that may depend on particle speed \vec{u} .

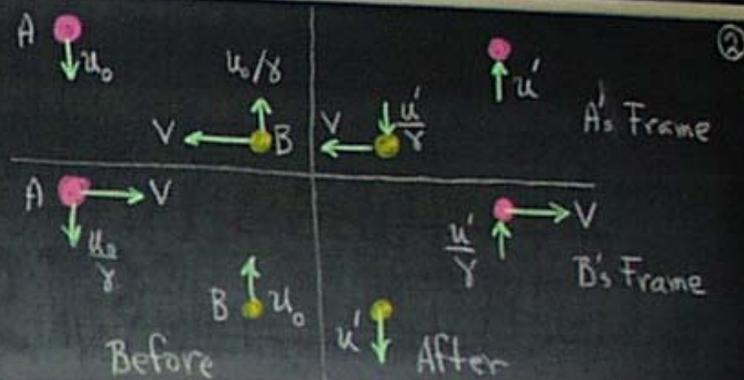
Frame-A:

x-momentum only due to particle-B
 Before collision B's speed

$$\omega^2 = v^2 + \frac{u_0^2}{\gamma^2}$$

After collision

$$\omega'^2 = v^2 + \frac{u'^2}{\gamma^2}$$



Conservation of momentum along-x

$$m(u)V = m(u')V$$

$$\therefore \omega = \omega'$$

$$\therefore u' = u_0$$

Conservation of momentum along-y
 in Frame-A

$$-m(u_0)u_0 + m(u)\frac{u_0}{\gamma} = m(u_0)u_0 - m(u)\frac{u_0}{\gamma}$$

$$\therefore m(u) = \gamma m(u_0)$$

let $m(u_0) = m_0$ as $u_0 \rightarrow 0$

m_0 = rest mass. Mass measured in frame where it is stationary.

In this limit $\omega = v$.

$$\therefore m(v) = \gamma m_0 = \frac{m_0}{\sqrt{1 - v^2/c^2}}$$

Interpret as dependence of mass on speed.

In general:

$$\vec{p} = \frac{m_0 \vec{u}}{\sqrt{1 - u^2/c^2}} = m \vec{u} = \gamma m_0 \vec{u}$$

$$m = \frac{m_0}{\sqrt{1 - u^2/c^2}} = \gamma m_0$$

Relativistic Energy.

classical:

$$k_b - k_a = \int_a^b \vec{F} \cdot d\vec{r} = \int_a^b \frac{d\vec{p}}{dt} \cdot \vec{r} dt$$

$$= \frac{1}{2} m u_b^2 - \frac{1}{2} m u_a^2$$

Relativity:

$$\vec{p} = \gamma m_0 \vec{u}$$

$$k_b - k_a = \int_a^b \frac{d\vec{p}}{dt} \cdot \vec{r} dt$$

$$= \int_a^b \frac{d}{dt} \left[\frac{m_0 \vec{u}}{\sqrt{1-u^2/c^2}} \right] \cdot \vec{u} dt$$

$$= \int_a^b \vec{u} \cdot d \left[\frac{m_0 \vec{u}}{\sqrt{1-u^2/c^2}} \right]$$

Integrand is $\vec{u} \cdot d\vec{p} = d(\vec{u} \cdot \vec{p}) - \vec{p} \cdot d\vec{u}$

$$k_b - k_a = \vec{u} \cdot \vec{p} \Big|_a^b - \int_a^b \vec{p} \cdot d\vec{u}$$

$$= \frac{m_0 u^2}{\sqrt{1-u^2/c^2}} \Big|_a^b - \int_a^b \frac{m_0 u du}{\sqrt{1-u^2/c^2}}$$

$$= \frac{m_0 u^2}{\sqrt{1-u^2/c^2}} \Big|_a^b + m_0 c^2 \sqrt{1-u^2/c^2} \Big|_a^b$$

b is arbitrary point
Assume $u_a = 0$ at a. (4)

$$K = k_b - k_a = \gamma m_0 u^2 + \frac{m_0 c^2}{\gamma} - m_0 c^2$$

$$K = m c^2 - m_0 c^2 \quad \text{where } m = \gamma m_0$$

Suppose $u \ll c$
 $\gamma \sim 1 + \frac{1}{2} \frac{u^2}{c^2} + \dots$

$$K = \frac{m_0 c^2}{\sqrt{1-u^2/c^2}} - m_0 c^2$$

$$\approx m_0 c^2 \left[1 + \frac{1}{2} \frac{u^2}{c^2} - 1 \right]$$

$$K = \frac{1}{2} m_0 u^2$$

Classical Result!!

Rewrite:

$$m c^2 = K + m_0 c^2$$

= Work Done + $m_0 c^2$

Einstein:

$$m c^2 = \text{Total energy } E \text{ of particle.}$$

$$= \text{Ext. Work} + \text{Rest Energy.}$$

$$\boxed{E = m c^2} = \gamma m_0 c^2$$

If energy ΔE is added to any object, it's mass changes by

$$\Delta m = \frac{\Delta E}{c^2}$$

Energy/Momentum

classically: $E = \frac{1}{2} m v^2 = \frac{p^2}{2m}$ ($u=0$) (5)

$$\vec{p} = m \vec{u} = \frac{m_0 \vec{u}}{\sqrt{1-u^2/c^2}} = \gamma m_0 \vec{u}$$

$$E = m c^2 = \gamma m_0 c^2$$

$$p^2 = \gamma^2 m_0^2 u^2 = \frac{1}{1-u^2/c^2} m_0^2 u^2$$

$$K = \frac{m_0 c^2}{\sqrt{1-u^2/c^2}} - m_0 c^2$$

$$\approx m_0 c^2 \left[1 + \frac{1}{2} \frac{u^2}{c^2} - 1 \right]$$

$$K = \frac{1}{2} m_0 u^2$$

Classical Result!!

Rewrite

$$m c^2 = K + m_0 c^2$$

= Work Done + $m_0 c^2$

Einstein:

$$m c^2 = \text{Total energy } E \text{ of particle.}$$

$$= \text{Ext. Work} + \text{Rest Energy.}$$

$$\therefore \boxed{E = m c^2} = \gamma m_0 c^2$$

If energy ΔE is added to any object, it's mass changes by

$$\Delta m = \frac{\Delta E}{c^2}$$

Energy/Momentum

Classically: $E = \frac{1}{2} m v^2 = \frac{p^2}{2m}$ ($u=0$)

$$\vec{p} = m \vec{u} = \frac{m_0 \vec{u}}{\sqrt{1-u^2/c^2}} = \gamma m_0 \vec{u}$$

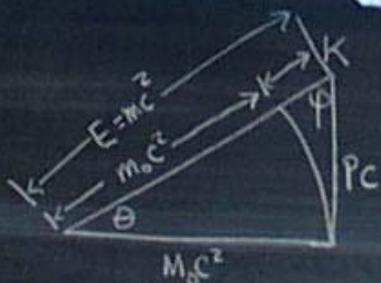
$$E = m c^2 = \gamma m_0 c^2$$

$$p^2 = \gamma^2 m_0^2 u^2 = \frac{1}{1-u^2/c^2} m_0^2 u^2$$

$$\frac{u^2}{c^2} = \frac{p^2}{p^2 + m_0^2 c^2}$$

$$\gamma = \frac{1}{\sqrt{1-u^2/c^2}} = \sqrt{1 + \frac{p^2}{m_0^2 c^2}}$$

$$\therefore E = m_0 c^2 \sqrt{1 + \frac{p^2}{m_0^2 c^2}}$$



$$\sin \theta = \beta = \frac{u}{c}$$

$$\sin \phi = \frac{1}{\gamma}$$

$$\therefore E^2 = (pc)^2 + (m_0 c^2)^2$$

$$E^2 - (pc)^2 = (m_0 c^2)^2$$

4-Vector!
Invariant Quantity

Massless Particles

$$E^2 = (pc)^2 + (m_0 c^2)^2$$

If $m_0 = 0$

$$E = pc$$

$$p = \gamma m_0 u = \frac{m_0}{\sqrt{1-u^2/c^2}} \vec{u}$$

As $m_0 \rightarrow 0$ \vec{p} must remain finite!

Only possible if $u \rightarrow c$ as $m_0 \rightarrow 0$

\therefore massless particles must travel with c .

Photons have $m_0 = 0$

Neutrinos have $m_0 \neq 0$

Force and Relativity.

$$\vec{F} = \frac{d\vec{p}}{dt} = \frac{d}{dt} (m\vec{v}) = m\frac{d\vec{v}}{dt} + \vec{v} \frac{dm}{dt} \quad (1)$$

$$m = E/c^2$$

$$\therefore \frac{dm}{dt} = \frac{1}{c^2} \frac{dE}{dt} = \frac{1}{c^2} \frac{d}{dt} (k + m_0 c^2) = \frac{1}{c^2} \frac{dk}{dt}$$

$$\text{But } \frac{dk}{dt} = \frac{\vec{F} \cdot d\vec{r}}{dt} = \vec{F} \cdot \vec{v}$$

$$\therefore \frac{dm}{dt} = \frac{1}{c^2} \vec{F} \cdot \vec{v}$$

Sub in (1)

$$\vec{F} = m\frac{d\vec{v}}{dt} + \frac{\vec{v}(\vec{F} \cdot \vec{v})}{c^2}$$

$$\vec{a} = \frac{d\vec{v}}{dt} = \frac{1}{m} \left(\vec{F} - \frac{\vec{v}(\vec{F} \cdot \vec{v})}{c^2} \right)$$

\vec{a} is not // to \vec{F} has comp. along \vec{v}

