22.615, MHD Theory of Fusion Systems Prof. Freidberg

Lecture 15: Variational Techniques

- 1. Variational Procedure
- 2. Variational formulation of MHD
- 3. Energy Principle
- 4. Extended Energy Principle

Variational Procedure

- 1. Alternate representation of a differential equation
- 2. Consider classic eigenvalue problem

$$\frac{d}{dx}\bigg(F\frac{dy}{dx}\bigg) + \left(\lambda - g\right)y = 0 \qquad y\left(0\right) = y\left(1\right) = 0 \qquad F = F\left(x\right), g = g\left(x\right)$$

- 3. Methods of solution
 - a. inspection if F, g simple
 - b. computer
 - c. power series (not useful usually)
 - d. fourier series (not useful usually)
 - e. variational procedure (allows guess and a solution) λ is more accurate than guess for y(x)
- 4. Variational principle: multiply by $\int_0^1 y dx \left[\int \rightarrow \int_0^1 \right]$ below

$$\int \left[y \frac{d}{dx} \left(F \frac{dy}{dx} \right) + (\lambda - g) y^2 \right] dx = 0$$

$$+ \int \left[-Fy^2 + \lambda y^2 - gy^2 \right] dx + Fy y \Big|_0^1 = 0$$

$$\lambda = \frac{\int \left(Fy^2 + gy^2 \right) dx}{\int y^2 dx} \tag{1}$$

5. Why is this variational? Substitute all allowable y(x) into (1). When resulting λ exhibits an extremum (maximum, minimum, saddle pt) then λ and y are actual eigenvalue and eigenfunction.

6. Proof: assume $y_0 = (x)$ is substituted into (1) yielding λ_0 , Modify y a little bit by a small but arbitrary perturbation.

$$y(x) = y_0(x) + \delta y(x)$$
 $\delta y(0) = \delta y(1) = 0$

This produces a change in $\lambda = \lambda_0 + \delta \lambda$ given by

$$\delta\lambda = \frac{\int\!\!\left[F\left(y_0^{} + \delta y\right)^2^{} + g\!\left(y_0^{} + \delta y\right)^2\right]\!dx}{\int\!\left(y_0^{} + \delta y\right)^2^{}dx} - \frac{\int\!\!\left(Fy_0^{'2}^{} + gy_0^2\right)\!dx}{\int\!y_0^2dx}$$

7. For small δy

$$\delta\lambda = \frac{\int \left(Fy_0^2 + gy_0^2\right) dx}{\int y_0^2 dx} - \frac{\int \left(Fy_0^2 + gy_0^2\right) dx}{\int y_0^2 dx} + 2 \frac{\int \left(Fy_0 \delta y + gy_0 \delta y\right) dx}{\int y_0^2 dx}$$

$$- \frac{2\int \left(Fy_0^2 + gy_0^2\right) dx}{\int y_0^2 dx} \frac{\int y_0 \delta y dx}{\int y_0^2 dx}$$

$$= \frac{2\int \left[Fy_0 \delta y + gy_0 \delta y - \lambda_0 y_0 \delta y\right] dx}{\int y_0^2 dx}$$

$$= \frac{2\int \left[Fy_0 \delta y + gy_0 \delta y - \lambda_0 y_0 \delta y\right] dx}{\int y_0^2 dx}$$

$$=\frac{2\int \left[Fy_0 \delta y' + gy_0 \delta y - \lambda_0 y_0 \delta y\right] dx}{\int y_0^2 dx}$$

$$\partial \lambda = \frac{-2 \int dx \, \delta y \left[\left(F y_0 \right)^{\cdot} + \left(\lambda_0 - g \right) y_0 \right]}{\int y_0^2 dx}$$

8. At an extremum $\delta \lambda = 0$ for <u>arbitrary</u> δy , implying

$$\left(Fy_{0}^{'}\right)^{'}+\left(\lambda_{0}-g\right)y_{0}=0$$

This is equivalent to original problem

9. One can multiply original equation by h(x) g(x) when h is arbitrary resulting

$$\lambda = \frac{\int \left\{ hFy^{'2} + \left[\left(hg - \left(Fh' \right)^{1/2} \right) \right] y^2 \right\} dx}{\int hy^2 dx}$$

is correct mathematical expression but <u>is not</u> a variational principle unless h=1. Varying y gives

$$\frac{d}{dx}\left(hF\frac{dy}{dx}\right) + \left[\lambda h - hg + \frac{\left(Fh'\right)^{2}}{2}\right]y = 0$$

or
$$\frac{d}{dx}\left(F\frac{dy}{dx}\right) + F\frac{h}{h}\frac{dy}{dx} + \left[\lambda - g + \frac{\left(Fh'\right)}{2h}\right]y = 0$$

- 10. Since $\delta\lambda=0$ when y coincides with a true eigenfunction, this implies that an estimate for λ using a guess (trial function) for g is more accurate than the trial function itself.
- 11. Proof: Write $y = y_0 + \delta y$ where y_0 is true eigenfunction and δy is the error in the guess, assumed of order ϵ . Substitute into (1) yields

$$\begin{split} \lambda &= \frac{N_0 + N_1 + N_2}{D_0 + D_1 + D_2} = \frac{N_0 + N_1 + N_2}{D_0} \left[1 - \frac{D_1}{D_0} - \frac{D_2}{D_0} + \frac{D_1^2}{D_0^2} + \ldots \right] \\ &= \frac{N_0}{D_0} + \frac{1}{D_0} \left[N_1 - D_1 \frac{N_0}{D_0} \right] + \frac{1}{D_0} \left[N_2 - \frac{N_1 D_1}{D_0} - D_2 \frac{N_0}{D_0} + N_0 \frac{D_1^2}{D_0^2} \right] + \ldots \\ &= \lambda_0 \qquad \left(\delta \lambda = 0 \right) \qquad + \frac{1}{D_0} \left[N_2 - D_2 \lambda_0 \right] \\ &= \lambda_0 + \frac{\int dx \left[F \left(\delta y' \right)^2 + \left(g - \lambda_0 \right) \left(\delta y \right)^2 \right]}{\int dx \ y_0^2} \\ &= \lambda_0 + 0 \left(\epsilon^2 \right) \end{split}$$

Thus, error is λ is $O(\epsilon^2)$ while error in y is $O(\epsilon)$

Generalized Boundary Conditions

1. Problem of practical importance

$$(Fy')' - gy = 0$$
 $y(0) = 1$ $y'(1) = Ay(1)$

2. Simple Variational Principle $> \int y dx$

$$\lambda = \frac{\int \left(Fy^2 - gy^2 \right) dx - fyy' \Big|_{x=1}}{\int y^2 dx}$$
 (2)

3. Let $y = y_0 + \delta y$, $\lambda = \lambda_0 + \delta \lambda$

$$\delta\lambda = -\frac{2\int\delta y \bigg[\Big(F\dot{y_0}\Big)^{'} + \Big(\lambda - g\Big)y_0\bigg] dx + F\Big(\dot{y_0}\delta y - y_0\delta\dot{y'}\Big)\bigg|_1}{\int y_0^2 dx}$$

- 4. $\delta\lambda$ vanishes if
 - a. y₀ satisfies original ODE
 - b. y_0 and δy satisfy $y_0(1) = Ay_0(1)$, $\delta y_0(1) = A\delta y(1)$
- 5. Using trial functions which satisfy g'(1) = Ag(1) is not unexpected but can be cumbersome in practice.
- 6. Alternate, more practical, more elegant variational principle. Replace y'(1) in (2) with Ay(1). Then

$$\lambda = \frac{\int \left(Fy_1^{2} + gy^2\right) dx + Afy^2 \Big|_1}{\int y^2 dx}$$

7. Let $y = y_0 + \delta y$, and $\lambda = \lambda_0 + \delta \lambda$

$$\delta \lambda = \frac{-2 \int \delta y \left[\left(f y_0 \right)^{'} + \left(\lambda_0 - g \right) y_0 \right] dx + 2 F \delta y \left(y_0 - A y_0 \right)}{\int y_0^2 \ dx}$$

- 8. δλ vanishes if
 - a. original ODE satisfied

b.
$$y_0' = Ay_0$$

If we choose trial functions which allow y(1) to float freely, then variational principle forces Ay(1) = y'(1). This is a <u>natural boundary condition</u>.

Practical Applications

1. "Exact" solution and eigenvalue. Let

complete set of orthornormal functions
$$y = \sum a_n \, Y_n \left(x \right)$$

2. Then

$$\lambda = \frac{\sum a_n a_m W_{nm}}{\sum a_n^2}, \quad W_{nm} = \int dx \Big[F y_n y_m + g y_n y_m \Big] dx$$

- 3. Minimize with respect to the a_n . This is equivalent to finding the eigenvalues of $\underline{\underline{W}}$. Simple numerical procedure. In the limit $N \to \infty \left(\sum_{1}^{N}\right)$, we obtain exact eigenvalue and eigenfunction. Since Y_n is a complete set.
- 4. Send "estimate" of eigenvalue. Gives a trial function

$$y = y(x, c_1, c_2, c_3) \quad \left[\text{e.g. } y = x(1-x) \Big[1 + c_1 x + c_2 x^2 + c_3 x^3 \Big] \right]$$
 or
$$y = x(1-x^2) e^{-c_1 x^2} \left[1 + c_2 x^{c_3} \right]$$

Evaluate $\lambda = \lambda(c_1, c_2, c_3)$

Find c_{1},c_{2},c_{3} to extremize λ

$$\frac{\partial \lambda}{\partial c_1} = 0$$

$$\frac{\partial \lambda}{\partial c_2} = 0$$

$$\frac{\partial \lambda}{\partial c_3} = 0$$

Substitute c_{1},c_{2},c_{3} back into λ to get good estimate.

Application of the Variational Principle to MHD

- 1. Generalize previous analysis to include 3-D and sectors
- 2. Conceptually the same

$$-\omega^2 \rho \underline{\xi} = \underline{F}\left(\underline{\xi}\right) > \int \underline{\xi}^* \cdot d\underline{r}$$

$$\omega^2 = \frac{\delta W}{K}$$

$$K = \frac{1}{2} \int \rho \left| \underline{\xi} \right|^2 dr$$

$$\omega^{2} = \frac{\delta W}{K}$$

$$K = \frac{1}{2} \int \rho \left| \underline{\xi} \right|^{2} dr$$

$$\delta W = -\frac{1}{2} \int \underline{\xi}^{*} \cdot \underline{F} \left(\underline{\xi} \right) d\underline{r}$$

3. Proof: Let $\underline{\xi} \rightarrow \underline{\xi} + \delta \underline{\xi}$, $\omega^2 \rightarrow \omega^2 + \delta \omega^2$

$$\omega^2 + \delta\omega^2 = \frac{\delta W\left(\underline{\xi}^{\star}, \underline{\xi}\right) + \delta W\left(\underline{\delta}\underline{\xi}^{\star}, \underline{\xi}\right) + \delta W\left(\underline{\xi}^{\star}, \delta\underline{\xi}\right) + \delta W\left(\delta\underline{\xi}^{\star}, \delta\underline{\xi}\right)}{K\left(\underline{\xi}^{\star}, \underline{\xi}\right) + K\left(\underline{\delta}\underline{\xi}^{\star}, \xi\right) + K\left(\underline{\delta}\underline{\xi}^{\star}, \delta\underline{\xi}\right) + K\left(\underline{\delta}\underline{\xi}^{\star}, \delta\underline{\xi}\right)}$$

4. Assume $\delta \xi$, $\delta \omega^2$ are small

$$\delta\omega^{2} = \frac{\delta W\left(\underline{\xi}^{*}, \delta\,\underline{\xi}\right) + \delta W\left(\delta\underline{\xi}^{*}, \underline{\xi}\right)}{K\left(\underline{\xi}^{*}, \xi\right)} - \frac{\delta W\left(\underline{\xi}^{*}, \xi\right)}{\underbrace{K\left(\underline{\xi}^{*}, \xi\right)}_{\omega^{2}}} \frac{\left[K\left(\delta\underline{\xi}^{*}, \underline{\xi}\right) + K\left(\underline{\xi}^{*}, \delta\,\underline{\xi}\right)\right]}{K\left(\underline{\xi}^{*}, \xi\right)}$$

$$=\frac{\delta W\left(\delta \underline{\xi}^{*},\underline{\xi}\right)-\omega^{2} K\left(\delta \underline{\xi}^{*},\xi\right)+\delta W\left(\underline{\xi}^{*},\delta \underline{\xi}\right)-\omega^{2} K\left(\xi^{*},\delta \underline{\xi}\right)}{K\left(\underline{\xi}^{*},\underline{\xi}\right)}$$

Use self adjoint property: $K\left(\underline{\xi}^*, \partial \underline{\xi}\right) = K\left(\underline{\delta \xi}, \underline{\xi}^*\right)$

$$\delta W\left(\underline{\xi}^{\star},\underline{\delta}\xi\right)=\delta W\left(\underline{\delta}\xi,\underline{\xi}^{\star}\right)$$

Then

$$\delta\omega^{2} = \int d\underline{r} \ \underline{\delta}\underline{\xi}^{\star} \cdot \left[F\left(\xi\right) + \omega^{2}\rho\underline{\xi} \right] + \underline{\delta}\underline{\xi} \left[\underline{F}\left(\underline{\xi}^{\star}\right) + \omega^{2}\rho\underline{\xi}^{\star} \right] \! \middle/ \! K$$

If $\delta \xi$ is arbitrary and $\delta \omega^2$ = 0 (extremum) then

$$\omega^2 \rho \underline{\xi} = -\underline{F}\left(\underline{\xi}\right)$$

Simple Interpretation

