

## 8.575J, 10.44J, 22.52J Statistical Thermodynamics of Complex Liquids

(Spring 2004)

Problem Set 1 (Prof. Chen)

Due March 18.

1. Given the following table for the scattering lengths of common elements:

Isotope	Hydrogen	Deuterium	Carbon	Nitrogen	Oxygen
$b_{\text{coh}} (10^{-12} \text{ cm})$	-0.37423	0.6674	0.66484	0.936	0.5805

estimate the molecular volume (from the molecular weight and density) and then calculate the scattering length densities of the following molecules, in unit of  $10^{10} \text{ cm}^{-2}$ :  $\text{H}_2\text{O}$ ,  $\text{D}_2\text{O}$ , Octane, Deuterated octane, and Pluronic P-84, a tri-block co-polymer,  $[(\text{PEO})_{19}(\text{PPO})_{43}(\text{PEO})_{19}]$ , where  $\text{PEO} = -(\text{CH}_2)_2\text{O}-$ , having a molecular volume  $72.4 \text{ \AA}^3$ , and  $\text{PPO} = -(\text{CH}_2)_3\text{O}-$ , having a molecular volume  $95.4 \text{ \AA}^3$ .

2. Show that the form factor of a spherical particle with an internal core of radius  $R_1$  and a scattering length density (sld)  $\rho_1$ , surrounded by a shell with an outer radius  $R_2$  and sld  $\rho_2$ , immersed in a solvent of sld  $\rho_s$ , is given by:

$$F_{2\text{-shell}}(Q) = \frac{4}{3} \pi R_1^3 (\rho_1 - \rho_2) \left[ \frac{3j_1(QR_1)}{QR_1} \right] + \frac{4}{3} \pi R_2^3 (\rho_2 - \rho_s) \left[ \frac{3j_1(QR_2)}{QR_2} \right]$$

Use this result to calculate and plot the normalized particle structure factor  $\bar{P}(Q)$  of a co-polymer micelle having an inner core radius  $R_1$  and an outer radius  $R_2$ . In a core-shell model of the micelle [Y.C. Liu et al, Phys. Rev. E 54, 1698 (1996)], the inner and outer radii can be determined from the aggregation (N) and hydration (H) numbers of the micelle. Plot the  $\bar{P}(Q)$  for the case of  $N = 63$  and  $H = 290$ .

3. Show that the form factor of a randomly oriented prolate spheroid, with semi-major and minor axes of  $a$  and  $b$ , is given by:

$$F_{\text{ellipsoid}}(Q) = \frac{4\pi}{3} ab^2 (\Delta\rho) \int_0^1 d\mu \left[ \frac{3j_1(u)}{u} \right] \quad (2)$$

$$u = Q \sqrt{a^2 \mu^2 + b^2 (1 - \mu^2)}$$

where  $(\Delta\rho)$  is the contrast between the particle and the solvent,  $\mu$  the cosine of the angle between the major axis of the spheroid and the Q-vector.

4. Derive a normalized particle structure factor  $\bar{P}(Q)$  of a uniform cylindrical particle of radius  $R$  and length  $L$ . Assuming that the particle is randomly oriented with respect to the  $\hat{Q}$  vector.

(A) Show that:

$$\bar{P}(Q) = \left\langle \left| \frac{1}{V_p} \int_{V_p} e^{i\hat{Q} \cdot \vec{r}} d^3r \right|^2 \right\rangle = \frac{1}{2} \int_{-1}^1 d\mu \left[ \frac{\sin QL\mu/2}{QL\mu/2} \right]^2 \left[ \frac{2J_1(QR\sqrt{1-\mu^2})}{QR\sqrt{1-\mu^2}} \right]^2. \quad (3)$$

In Eq.3,  $V_p$  denotes the volume of the particle,  $\vec{r}$  the position vector of an arbitrary point in the interior of the particle, and  $\mu$  the cosine of the angle between the axis of the cylinder and the  $Q$ -vector. The bracket means that we are considering an average over random orientations of the particle.

(B) Show that for a long and thin cylinder, one has asymptotic formulae:

$$\bar{P}(Q) \xrightarrow{QL > 2\pi} \frac{\pi}{QL} \left[ \frac{2J_1(QR)}{QR} \right]^2 \xrightarrow{QR < 1} \frac{\pi}{QL} e^{-\frac{1}{4}Q^2R^2} \quad (4)$$

(C) Show that for a flat particle (a lamellar) of  $QR \gg 1$ ,

$$\bar{P}(Q) \xrightarrow{QR > 1} \frac{2}{Q^2R^2} \left[ \frac{\sin QL/2}{QL/2} \right]^2 \xrightarrow{QL < 1} \frac{2}{Q^2R^2} e^{-\frac{1}{12}Q^2L^2} \quad (5)$$

where  $L$  is the thickness of the flat plate.

(D) From Eq.4 and 5, one can conclude that for a long rod a  $\ln[QI(Q)]$  vs  $Q^2$  plot, and for a flat disk, a  $\ln[Q^2I(Q)]$  vs  $Q^2$  plot, will result in a straight line at large  $Q$  with slopes proportional to  $R^2/4$  and  $L^2/12$  respectively. Explore additional system parameters you can extract from the intercept at  $Q = 0$ .

(E) In polymer literature, another approximate formula is often used. It is the limit when  $R$  goes to zero, the so called "stiff thin rod" limit. Show that:

$$\bar{P}(Q) \underset{R \rightarrow 0}{=} \int_0^1 d\mu \left[ \frac{\sin QL\mu/2}{QL\mu/2} \right]^2 = \int_0^1 d\mu \left[ \frac{\sin x\mu}{x\mu} \right]^2 = \frac{1}{x} \int_0^{2x} \frac{\sin u}{u} du - \left( \frac{\sin x}{x} \right)^2. \quad (6)$$

Explore graphically the difference between approximations of the last equation and Eq. 3

5. (This problem is added for your interest only. Your answer is optional. In case you solve it, you will get a bonus points of 20/100)

Scattering intensity of a Gaussian chain. Consider a flexible polymer chain of a contour length  $L = Na$ , where  $N$  is the no. of segments and  $a$  is the Kuhn length. If the chain makes a random walk in space, then the mean square end-to-end distance is  $R^2 = Na^2$ . For this chain, the distribution of distances between two links  $(i,j)$  is Gaussian, namely,

$$P(R_{ij})dR_{ij} = \left[ \frac{3}{2\pi\langle R_{ij}^2 \rangle} \right]^{3/2} 4\pi R_{ij}^2 \exp\left[ \frac{-3R_{ij}^2}{2\langle R_{ij}^2 \rangle} \right] dR_{ij}, \text{ where } \langle R_{ij}^2 \rangle = |i-j|a^2. \quad (7)$$

In order to calculate the normalized particle structure factor of such chain. (a) Start from the definition:

$$\bar{P}(Q) = \frac{1}{N^2} \left\langle \sum_{i=1}^N \sum_{j=1}^N e^{i\vec{Q} \cdot (\vec{r}_i - \vec{r}_j)} \right\rangle = \frac{1}{N^2} \sum_{i=1}^N \sum_{j=1}^N \left\langle e^{i\vec{Q} \cdot \vec{R}_{ij}} \right\rangle_{\text{gaussian}} \quad (8)$$

so that we can evaluate the Gaussian average by integrating the exponential phase factor using the distribution function given by Eq.7. (b) Show first that :

$$\bar{P}(Q) = \frac{1}{N^2} \sum_{i=1}^N \sum_{j=1}^N \exp\left[-\frac{1}{6}\langle R_{ij}^2 \rangle\right] = \frac{1}{N^2} \sum_{i=1}^N \sum_{j=1}^N \exp\left[-\frac{1}{6}Q^2 a^2 |i-j|\right] \quad (9)$$

(c) Prove a theorem: For an arbitrary function  $f(x)$ ,

$$\sum_{i=1}^N \sum_{j=1}^N f(|i-j|) = Nf(0) + 2 \sum_{n=1}^{N-1} (N-n)f(n) \quad (10)$$

(d) Use the theorem to show that the sum in Eq.9 can be evaluated as:

$$\bar{P}(Q) = \frac{1}{N^2} \left[ N + 2 \frac{(N-1)\alpha - N\alpha^2 + \alpha^{N+1}}{(1-\alpha)^2} \right] \quad (11)$$

where  $\alpha = \exp\left(\frac{1}{6}Q^2 a^2\right) \approx 1 - \frac{1}{6}Q^2 a^2$ , because  $Qa$  is much smaller than unity in practice.

(e) Show that in the limit  $N \rightarrow \infty$ ,  $\bar{P}(Q)$  approach the Debye function

$$\bar{P}(Q) = \frac{2}{x^2} (x - 1 + e^{-x}), \text{ where } x = \frac{1}{6}Q^2 R^2 = \frac{1}{6}Q^2 a^2 N \quad (12)$$

(f) Discuss the small and large  $Q$  behavior of Eq.12. In particular, show that:

$$\lim_{x \rightarrow \infty} \frac{1}{P(Q)} = \frac{1}{2} + \frac{1}{12} Q^2 R^2 \quad (13)$$