

# Quantization of the electromagnetic field

# The classical electromagnetic field

# Maxwell Equations

Gauss's law

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

Gauss's law for magnetism

$$\nabla \cdot \mathbf{B} = 0$$

Maxwell-Faraday equation

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

(Faraday's law of induction)

Ampere's circuital law

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

(with Maxwell's correction)

# Maxwell Equations

- In empty space  $(c = 1/\sqrt{\mu_0\epsilon_0})$

Gauss's law

$$\nabla \cdot \mathbf{E} = 0$$

Gauss's law for magnetism

$$\nabla \cdot \mathbf{B} = 0$$

Maxwell-Faraday equation

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

Ampere's circuital law

$$\nabla \times \mathbf{B} = \frac{1}{c} \frac{\partial \mathbf{E}}{\partial t}$$

# Wave Equations

$$\nabla^2 E - \frac{1}{c^2} \frac{\partial^2 E}{\partial t^2} = 0$$

$$\nabla^2 B - \frac{1}{c^2} \frac{\partial^2 B}{\partial t^2} = 0$$

# Derivation of wave equations

- Curl of Maxwell Faraday equation

$$\nabla \times (\nabla \times \mathbf{E}) = -\frac{1}{c} \frac{\partial \nabla \times \mathbf{B}}{\partial t}$$

- Use Ampere's Law

and vector identity  $\nabla \times (\nabla \times \vec{v}) = \nabla(\nabla \cdot \vec{v}) - \nabla^2 \vec{v}$

$$\nabla(\nabla \cdot \mathbf{E}) - \nabla^2 \mathbf{E} = -\frac{1}{c} \frac{\partial}{\partial t} \left( \frac{1}{c} \frac{\partial \mathbf{E}}{\partial t} \right)$$

# Derivation of wave equations

- Use Gauss Law

$$\nabla(\nabla \cdot \mathbf{E}) - \nabla^2 \mathbf{E} = -\frac{1}{c} \frac{\partial}{\partial t} \left( \frac{1}{c} \frac{\partial \mathbf{E}}{\partial t} \right)$$

- Obtain wave equation

$$-\nabla^2 \mathbf{E} = -\frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2}$$

# Wave equation

$$\nabla^2 E - \frac{1}{c^2} \frac{\partial^2 E}{\partial t^2} = 0$$

- Eigenvalue equation from separation of

variables:  $\mathbf{E}(\vec{x}, t) = \sum_m f_m(t) \vec{u}_m(\vec{x})$

$$\nabla^2 u_m = -k_m^2 u_m \quad \frac{d^2 f_m}{dt^2} + c^2 k_m^2 f_m(t) = 0$$

# Normal modes

- $\{u_m\}$  are eigenfunctions of the wave equation
- Boundary conditions (from Maxwell eqs.)

$$\nabla \cdot u_m = 0, \quad \vec{n} \times u_m = 0$$

- Orthonormality condition

$$\int \vec{u}_m(x) \vec{u}_n(x) d^3x = \delta_{n,m}$$

- They form a basis.

# B-field

- Electric field in  $\{u_m\}$  basis:

$$\mathbf{E}(\vec{x}, t) = \sum_m f_m(t) \vec{u}_m(\vec{x})$$

- Magnetic field in  $\{u_m\}$  basis

$$\mathbf{B}(x, t) = \sum_m h_m(t) (\nabla \times u_m(x))$$

# B-field solution

- What are the coefficients  $h_n$ ?
- We still need to satisfy Maxwell equations:

$$\nabla \times E = -\frac{1}{c} \partial_t B \quad \rightarrow$$

$$\sum_n f_n(t) \nabla \times u_n = -\frac{1}{c} \sum_n \partial_t h_n(t) \nabla \times u_n$$

- Solution:  $\frac{d h_n}{d t} = -c f_n$

# Eigenvalues of $h_n$

- Find equation for  $h_n$  only: Ampere's law

$$\nabla \times B = \frac{1}{c} \frac{\partial E}{\partial t}$$

$$\sum_n h_n(t) \nabla \times (\nabla \times u_n) = \frac{1}{c} \sum_n \frac{d f_n}{d t} u_n$$

$$\rightarrow - \sum_n h_n \nabla^2 u_n = \frac{1}{c} \sum_n \frac{d f_n}{d t} u_n$$

# Eigen-equations

- Eigenvalue equation for  $h_n$

$$\frac{d^2}{dt^2} h_n(t) + c^2 k_n^2 h_n(t) = 0$$

- Eigenvalue equation for  $f_n$

$$\frac{d^2}{dt^2} f_n + c^2 k_n^2 f_n(t) = 0$$

# E.M. field Hamiltonian

- Total energy:

$$\mathcal{H} \propto \frac{1}{2} \int (E^2 + B^2) d^3x$$

- Substituting, integrating and using orthonormality conditions:

$$\mathcal{H} = \frac{1}{8\pi} \sum_{n,m} \left( f_n f_m \int u_n(x) u_m(x) d^3x + h_n h_m \int (\nabla \times u_n) \cdot (\nabla \times u_m) d^3x \right)$$

$$\mathcal{H} = \sum_n \frac{1}{8\pi} (f_n^2 + k_n^2 h_n^2)$$

# E.M. field as H.O.

- Hamiltonian looks very similar to a sum of harmonic oscillators:

$$\mathcal{H}_{h.o.} = \frac{1}{2} \sum_n (p_n^2 + \omega_n^2 q_n^2) \Leftrightarrow \mathcal{H}_{e.m.} = \frac{1}{2} \sum_n \frac{1}{4\pi} (f_n^2 + k_n^2 h_n^2)$$

- $h_n$  is derivative of  $f_n$   
 $\Rightarrow$  identify with momentum

# Quantized electromagnetic field

# Operators

- We associate quantum operators to the coefficients  $f_n$ ,  $f_n \rightarrow \hat{f}_n$
- We write this operator in terms of annihilation and creation operators

$$\hat{f}_n = \sqrt{2\pi\omega_n\hbar}(a_n^\dagger + a_n)$$

that create or destroy one mode of the e.m. field

# Operator fields

- Electric field

$$\mathbf{E}(x, t) = \sum_n \sqrt{2\hbar\pi\omega_n} [a_n^\dagger(t) + a_n(t)] \mathbf{u}_n(x)$$

- Magnetic field

$$\mathbf{B}(x, t) = \sum_n i c_n \sqrt{\frac{2\pi\hbar}{\omega_n}} [a_n^\dagger - a_n] \nabla \times \mathbf{u}_n(x)$$

# Hamiltonian

- The Hamiltonian is then simply expressed in terms of the  $a_n$  operators

$$\mathcal{H} = \sum_n \omega_n \left( a_n^\dagger a_n + \frac{1}{2} \right)$$

- The frequencies are

$$\omega_n(k) = c|\vec{k}_n|$$

# Gauges

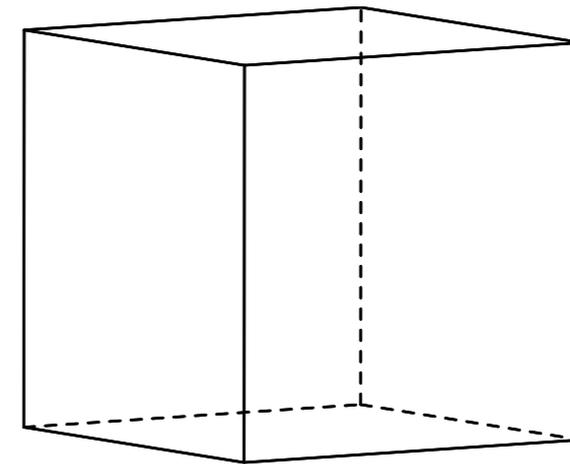
Lorentz (scalar potential  $\varphi = 0$ )  
Coulomb (vector potential  $\nabla \cdot \vec{A} = 0$ )

# Zero-Point Energy

# Field in cavity

- Field in a cavity of volume  $V = L_x L_y L_z$

- Given the boundary conditions, the normal modes are:



$$u_{n,\alpha} = A_\alpha \cos(k_{n,x} r_x) \sin(k_{n,y} r_y) \sin(k_{n,z} r_z)$$

- with  $k_{n,\alpha} = \frac{n_\alpha \pi}{L_\alpha}$ ,  $n_\alpha \in \mathcal{N}$

# Polarization

- Because of the boundary condition,

$$\nabla \cdot \vec{u}_n = 0$$

- the coefficients  $A$  must satisfy:

$$A_x k_{n,x} + A_y k_{n,y} + A_z k_{n,z} = 0$$

- For each set  $\{n_x, n_y, n_z\}$  there are 2 solutions

**Two polarizations per each mode**

# Electric field in cavity

- The electric field has a simple form

$$E(x, t) = \sum_{\alpha=1,2} (\mathcal{E}_\alpha + \mathcal{E}_\alpha^\dagger)$$

- with  $\mathcal{E}_\alpha^\dagger = \hat{e}_\alpha \sum_n \mathcal{E}_n a_n^\dagger e^{i(\vec{k}_n \cdot \vec{r} - \omega t)}$

- and  $\mathcal{E}_n = \sqrt{\frac{\hbar \omega_n}{2\epsilon_0 V}}$  the field of one photon  
of frequency  $\omega_n$

# Energy density

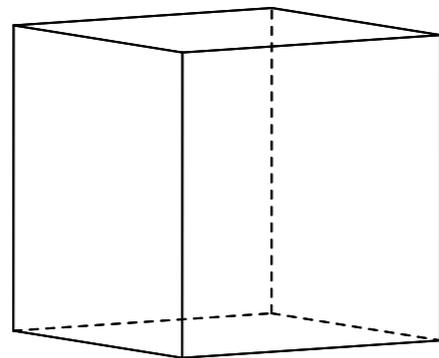
$$E = \langle \mathcal{H} \rangle = 2 \sum_{k=1}^{k_c} \hbar \omega_k \left\langle a_k^\dagger a_k + \frac{1}{2} \right\rangle$$

- The Zero-point energy density is then

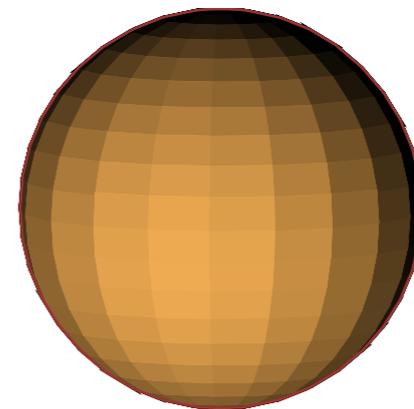
$$E_0 = \frac{2}{V} \sum_{k=1}^{k_c} \frac{1}{2} \hbar \omega_k$$

# Energy density

- If cavity is large, wavevector is almost continuous



$$\sum_{k>0}$$



$$\frac{1}{8} \int d^3 k \rho(k)$$

# Zero-point energy

- Integrating over the positive octant

$$E_0 = \frac{2}{V} \frac{2V}{\pi^3} \frac{4\pi}{8} \int_{k=0}^{k_c} dk \frac{1}{2} \hbar k^3 c$$

- setting a cutoff  $k_c$ , we have

$$E_0 = \frac{c\hbar}{2\pi^2} \int_{k=0}^{k_c} dk k^3 = \frac{\hbar c k_c^4}{8\pi^2}$$

# Zero-point energy

- It's huge!

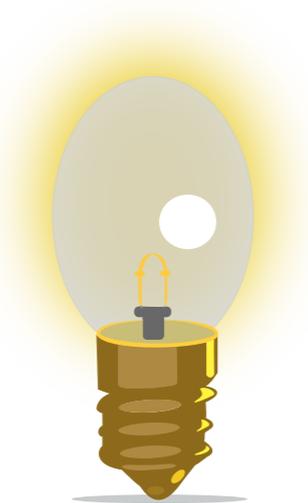


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Cutoff at visible frequency

$$\lambda_c = 2\pi/k = 0.4 \times 10^{-6} m$$

$$2.7 \times 10^{-8} \text{ J/m}^3 \text{ @ } 1\text{m}$$

$$23 \text{ J/m}^3$$

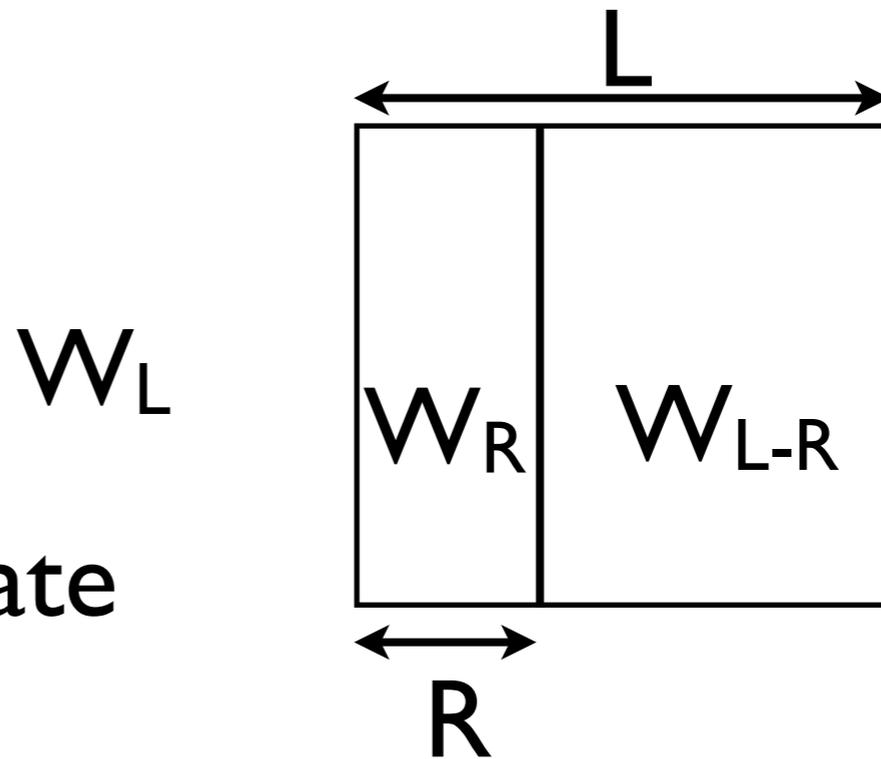
- But is it ever seen?

# Casimir Effect

- Dutch theoretical physicist Hendrik Casimir (1909–2000) first predicted in 1948 that when two mirrors face each other in vacuum, fluctuations in the vacuum exert “radiation pressure” on them

# Casimir Effect

- Cavity bounded by conductive walls
- Add a conductive plate @ distance  $R$
- Change in energy is:



$$\Delta W = (W_R + W_{L-R}) - W_L$$

# Casimir effect

- Each term is calculated from zero-point energy
- Continuous approximation is not valid if  $R$  is small
- Thus the difference  $\Delta W$  is not zero

$$\Delta W = -\hbar c \frac{\pi^2}{720} \frac{L^2}{R^3}$$

# Casimir Force

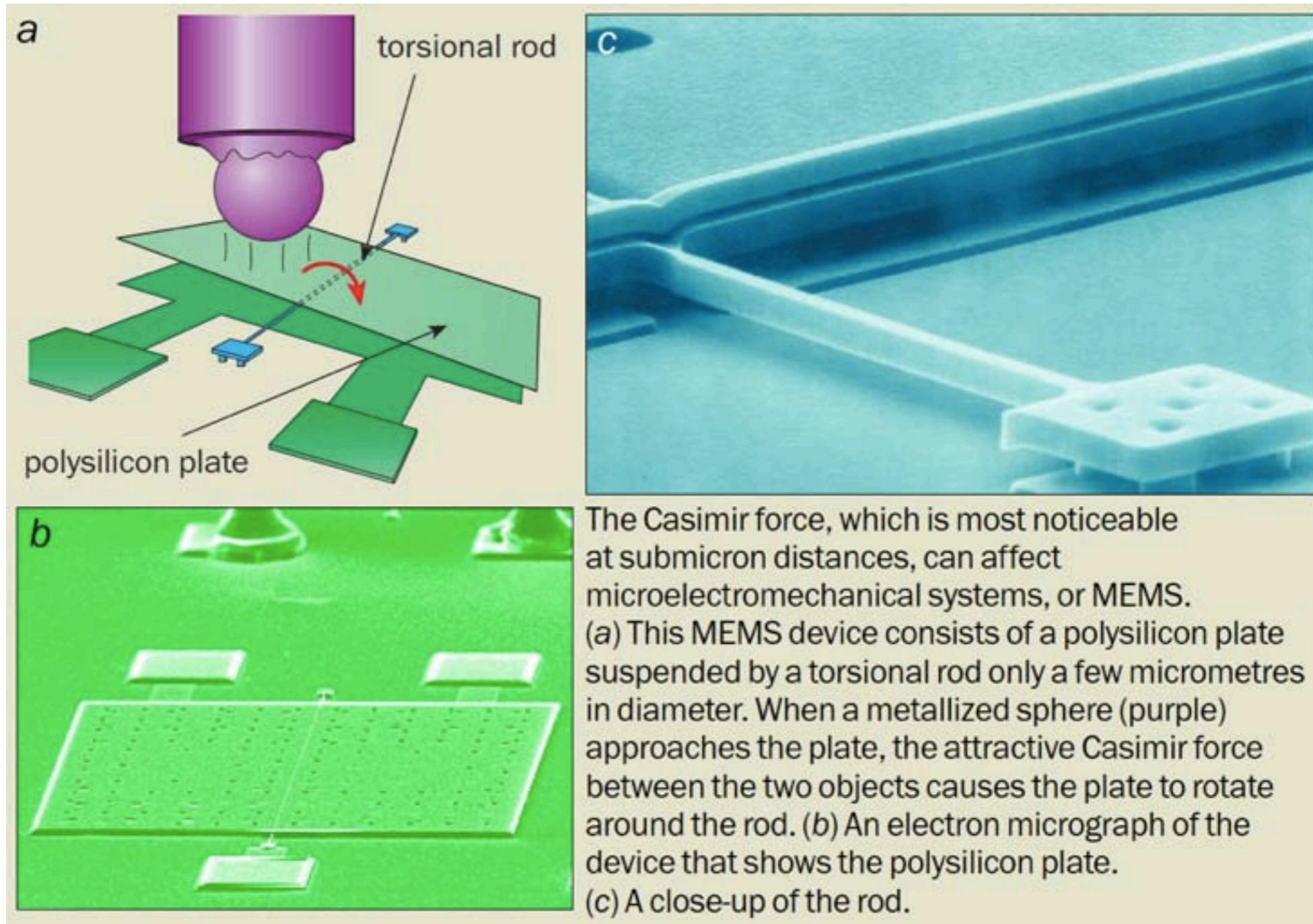
- The difference in energy corresponds to an attractive force

$$F = -\frac{\partial \Delta W}{\partial R} = -\hbar c \frac{\pi^2}{240} \frac{L^2}{R^4}$$

- or a pressure

$$P = -\frac{\pi^2}{240} \frac{\hbar c}{R^4}$$

# Casimir in MEMS



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## Quantum Mechanical Actuation of Microelectromechanical Systems by the Casimir Force

H. B. Chan, V. A. Aksyuk, R. N. Kleiman, D. J. Bishop and Federico Capasso

Science 9 March 2001: Vol. 291 no. 5510 pp. 1941-1944

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22.51 Quantum Theory of Radiation Interactions  
Fall 2012

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