

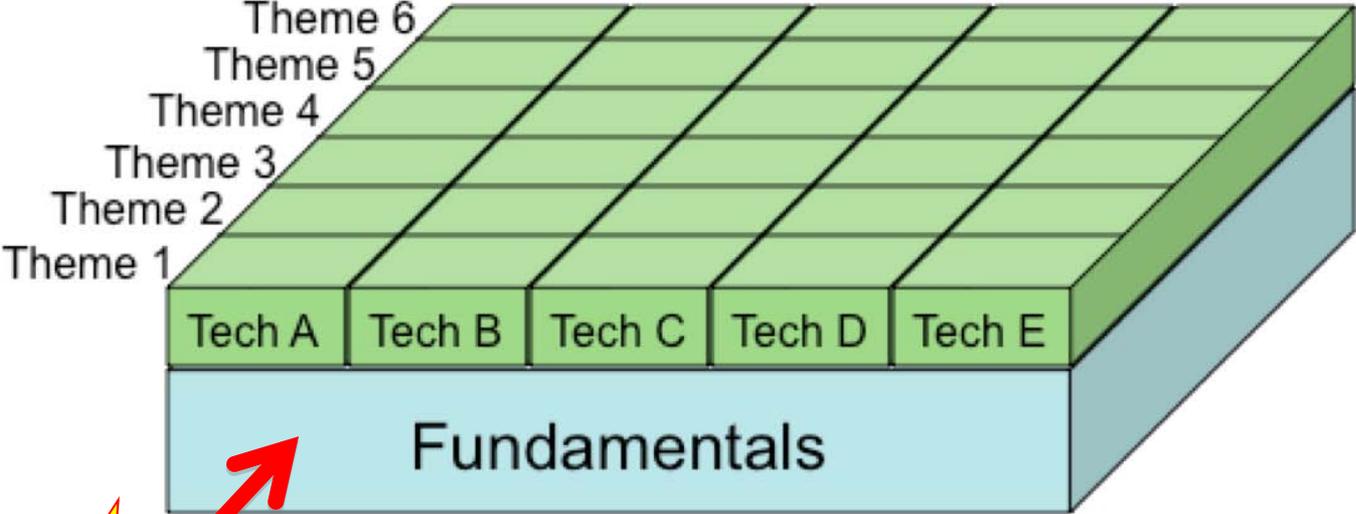
Toward a 1D Device Model Part 2: Material Fundamentals

Lecture 8 – 10/4/2011

MIT Fundamentals of Photovoltaics
2.626/2.627 – Fall 2011

Prof. Tonio Buonassisi

2.626/2.627 Roadmap



You Are Here

2.626/2.627: Fundamentals

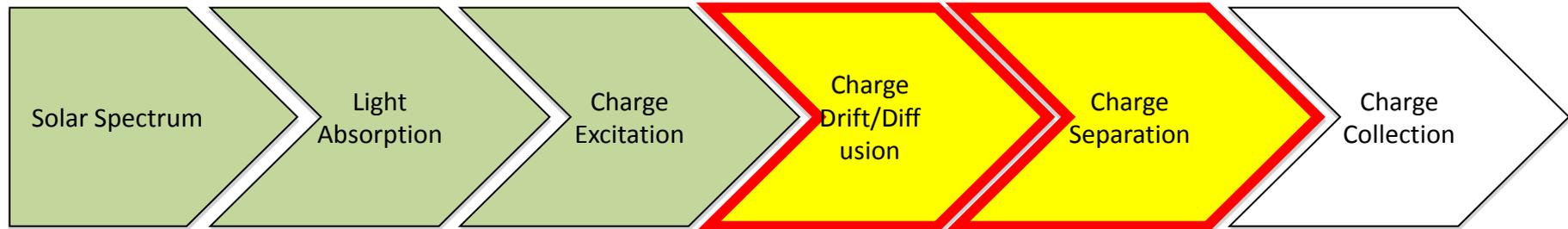
Every photovoltaic device must obey:

$$\text{Conversion Efficiency } (\eta) \equiv \frac{\text{Output Energy}}{\text{Input Energy}}$$

For most solar cells, this breaks down into:

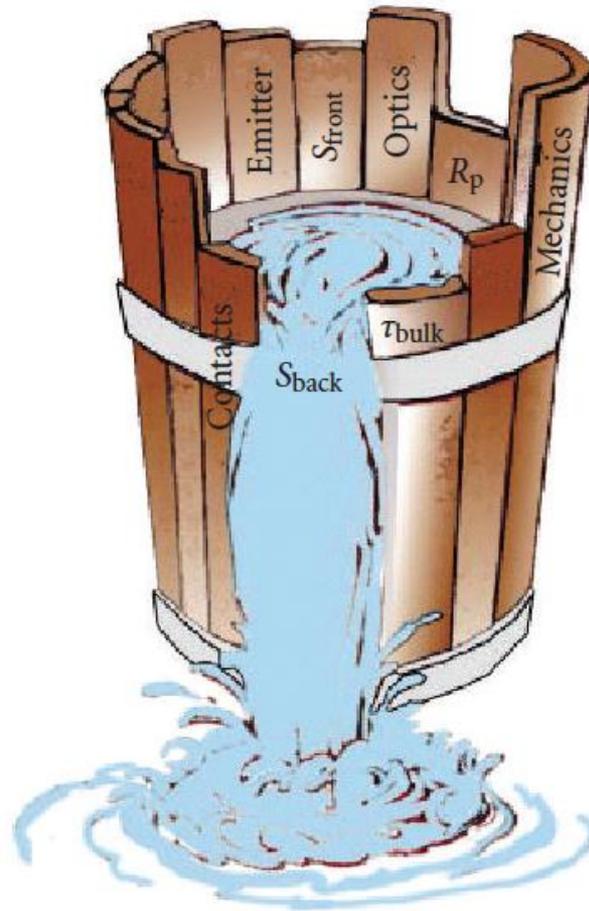
Inputs

Outputs



$$\eta_{\text{total}} = \eta_{\text{absorption}} \times \eta_{\text{excitation}} \times \eta_{\text{drift/diffusion}} \times \eta_{\text{separation}} \times \eta_{\text{collection}}$$

Liebig's Law of the Minimum

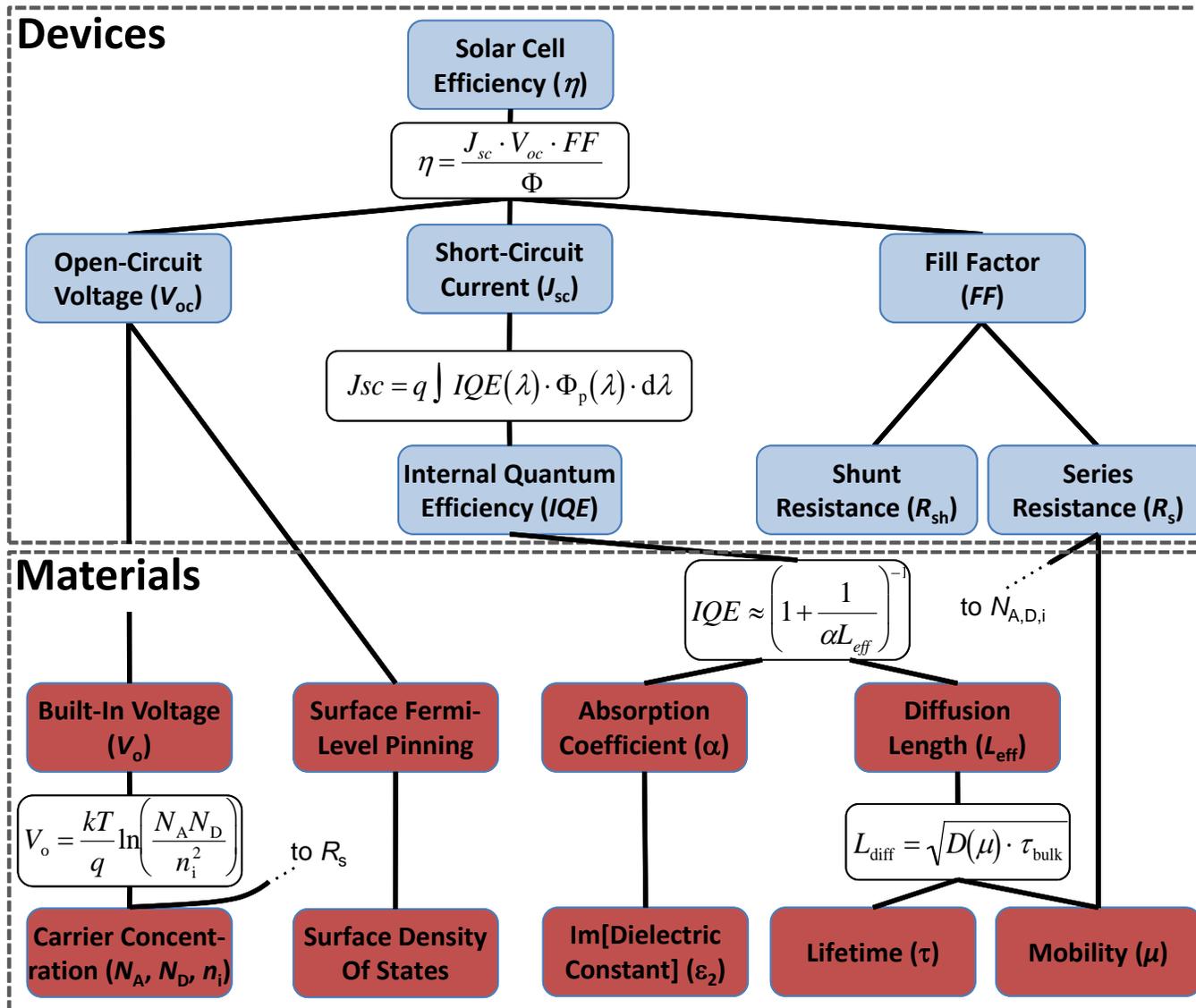


S. Glunz, *Advances in Optoelectronics* 97370 (2007)

Image by S. W. Glunz. License: CC-BY. Source: "[High-Efficiency Crystalline Silicon Solar Cells](#)." *Advances in OptoElectronics* (2007).

$$\eta_{\text{total}} = \eta_{\text{absorption}} \times \eta_{\text{excitation}} \times \eta_{\text{drift/diffusion}} \times \eta_{\text{separation}} \times \eta_{\text{collection}}$$

Rough Depiction of Interrelated Materials & Device Effects



Source: T. Buonassisi, unpublished.

Buonassisi (MIT) 2011

Learning Objectives: Toward a 1D Device Model

1. Describe what minority carrier diffusion length is, and calculate its impact on J_{sc} , V_{oc} . Describe how minority carrier diffusion length is affected by minority carrier lifetime and minority carrier mobility.
2. Describe how minority carrier diffusion length is measured.
3. Lifetime:
 - Describe basic recombination mechanisms in semiconductor materials.
 - Calculate excess carrier concentration as a function of carrier lifetime and generation rate. Compare to background (intrinsic + dopant) carrier concentrations.
4. Mobility:
 - Describe common mobility-limiting mechanisms (dopants, temperature, ionic semiconductors).

Minority Carrier Diffusion Length

Definition: Minority carrier diffusion length is the average distance a minority carrier moves before recombining.

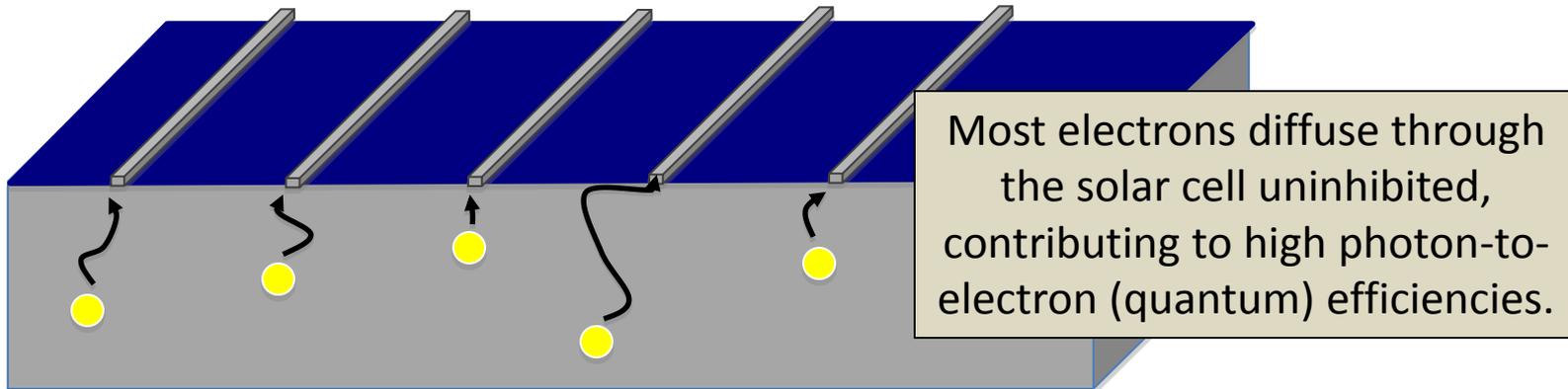
Importance to a Solar Cell: Photoexcited carriers must be able to move from their point of generation to where they can be collected. Longer diffusion lengths generally result in better performance.

Minority Carrier Diffusion Length

Definition: The average distance a minority carrier moves before recombining.

Importance to a Solar Cell: Carriers must be able to move from their point of generation to where they can be collected.

Cross section of solar cell made of high-quality material
Minority carrier diffusion length (L_{diff}) is **LARGE**.
Solar cell current output (J_{sc}) is large.

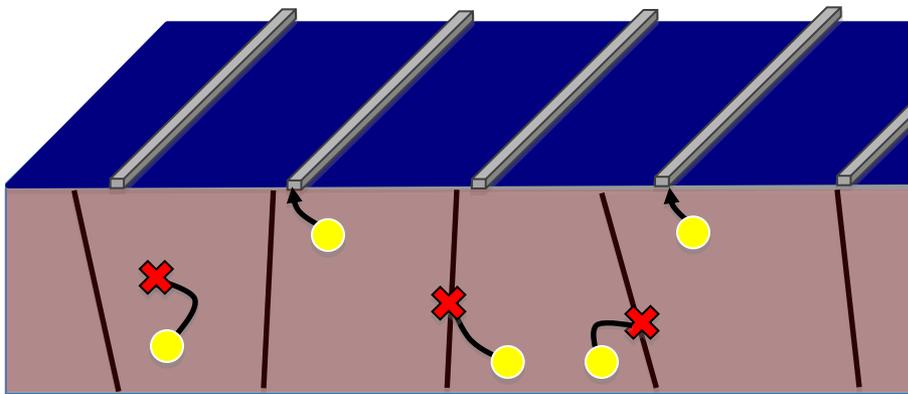


Minority Carrier Diffusion Length

Definition: The average distance a minority carrier moves before recombining.

Importance to a Solar Cell: Carriers must be able to move from their point of generation to where they can be collected.

Cross section of solar cell made of defect-ridden material
Minority carrier diffusion length (L_{diff}) is small.
Solar cell current output (J_{sc}) is small.



Electrons generated closer to the surface make it to the contacts, but those in the bulk are likely to “recombine” (lose their energy, e.g., at bulk defects, and not contribute to the solar cell output current).

Mathematical Formalism

Recall that the current produced in an illuminated pn-junction device, is limited by the minority carrier flux at the edge of the space-charge region.

$$J_{sc} \approx qGL_{diff}$$

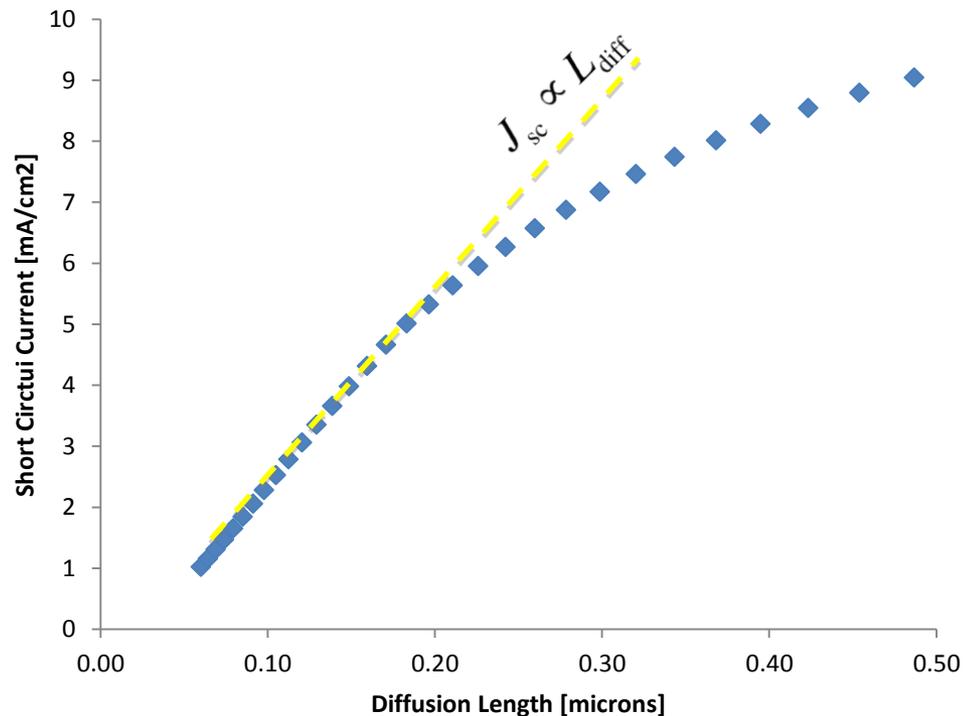
$$J_{sc} \propto L_{diff}$$

From Eq. 4.44 in Green.

J_{sc} = illuminated current = J_L

G = carrier generation rate

PC1D Simulation for 300um thick Si Solar Cell



Mathematical Formalism

Note that the voltage is also affected by L_{diff} .

$$V_{\text{oc}} = \frac{k_{\text{B}}T}{q} \ln\left(\frac{J_{\text{sc}}}{J_0} + 1\right)$$

From Eq. 4.45 in Green.

V_{oc} = open-circuit voltage

J_0 = saturation current density

$$J_0 \approx \frac{qDn_i^2}{L_{\text{diff}}N}$$

From Eq. 4.37 in Green.

D = minority carrier diffusivity

N = majority carrier dopant concentration

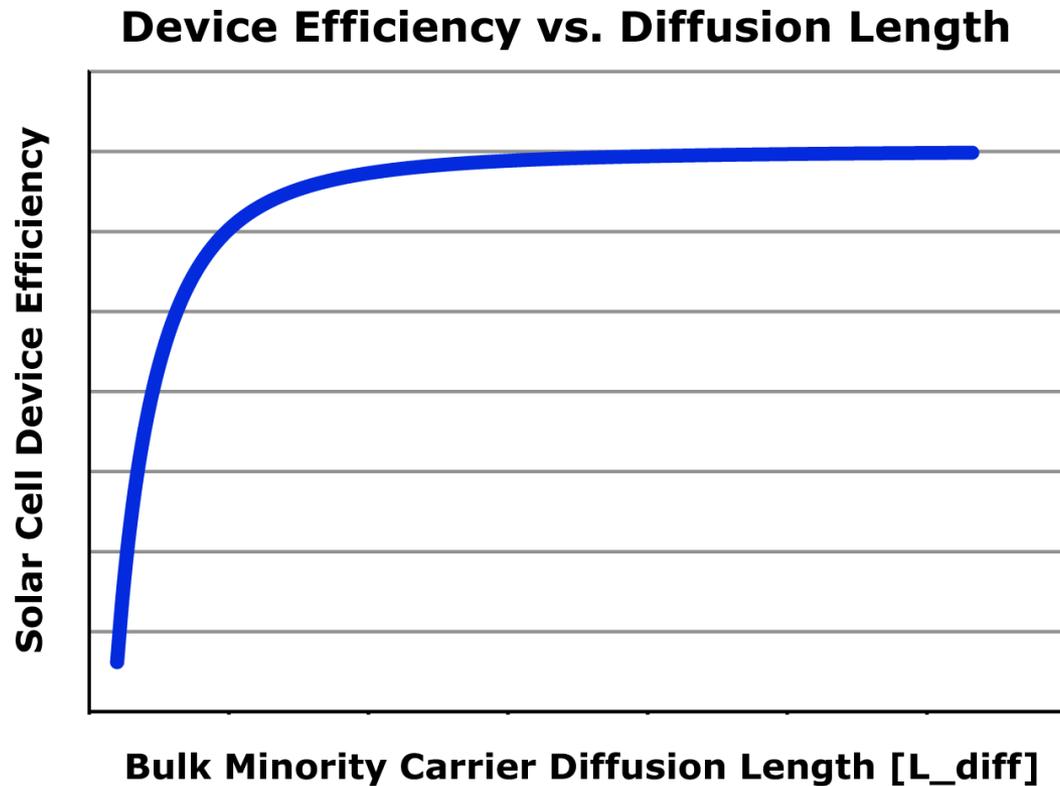
n_i = intrinsic carrier concentration

$$V_{\text{oc}} \propto \ln(L_{\text{diff}}^2) \propto 2 \ln(L_{\text{diff}}) \propto \ln(L_{\text{diff}})$$

Mathematical Formalism

Assuming a weak dependence of FF on L_{diff} , we have the following relationship:

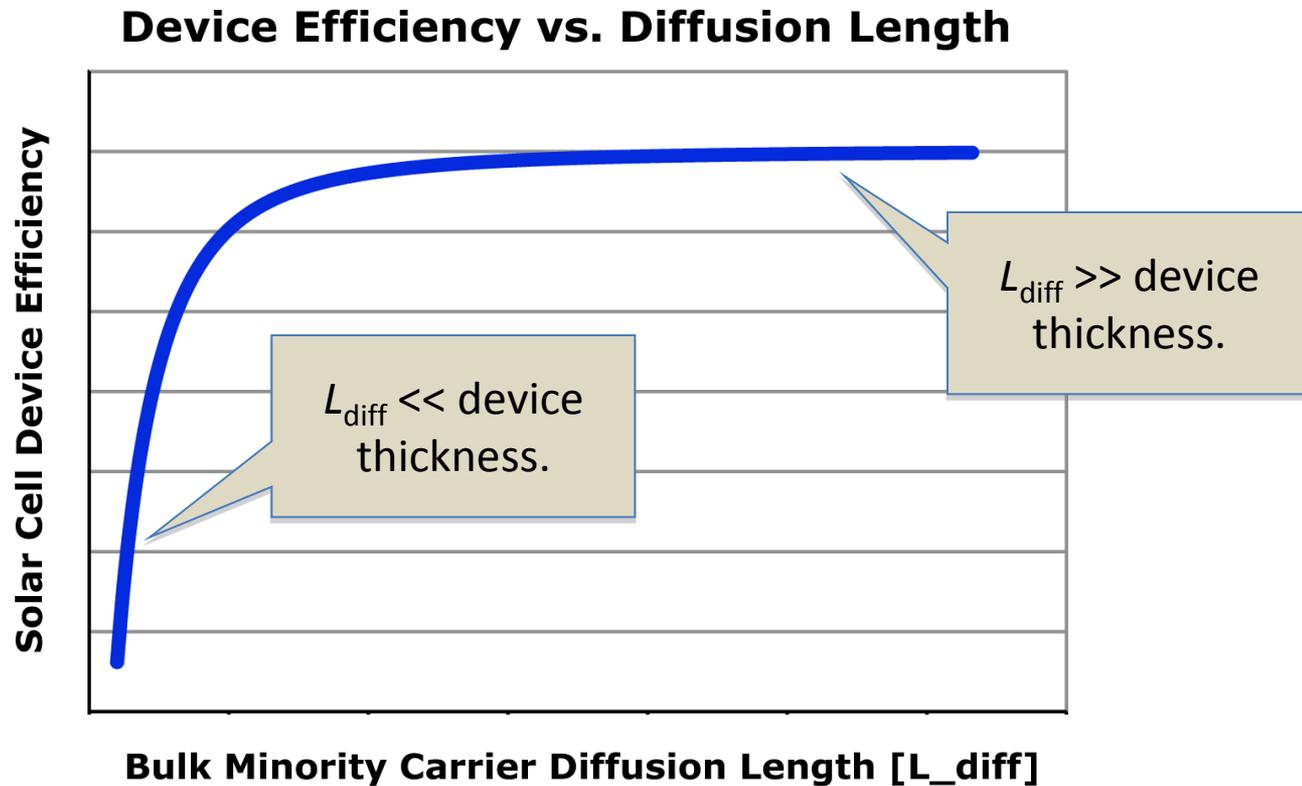
$$\eta \propto J_{\text{sc}} V_{\text{oc}} \propto L_{\text{diff}} \ln(L_{\text{diff}})$$



Mathematical Formalism

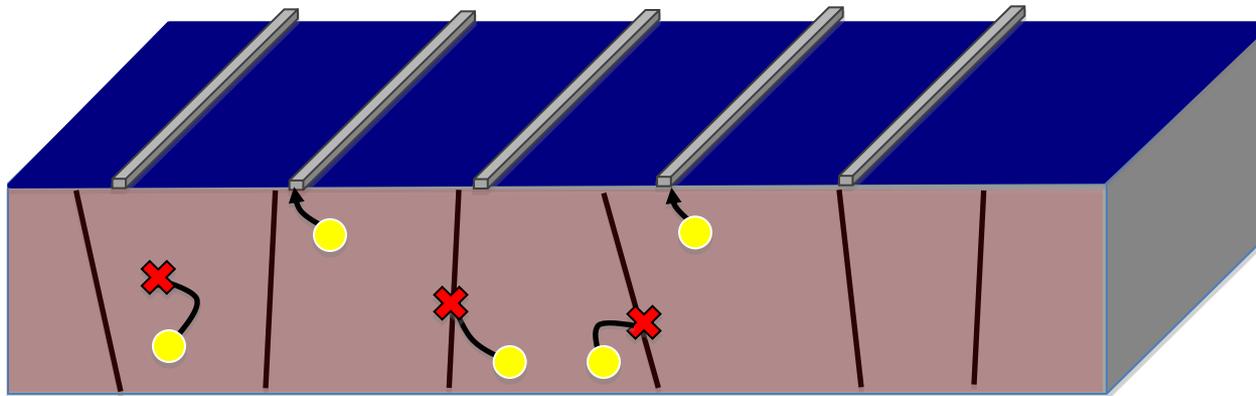
Assuming a weak dependence of FF on L_{diff} , we have the following relationship:

$$\eta \propto J_{\text{sc}} V_{\text{oc}} \propto L_{\text{diff}} \ln(L_{\text{diff}})$$



Minority Carrier Diffusion Length

When your material has short minority carrier diffusion length relative to absorber thickness, two engineering options:
Reduce absorber layer thickness (if light management permits) or increase diffusion length.

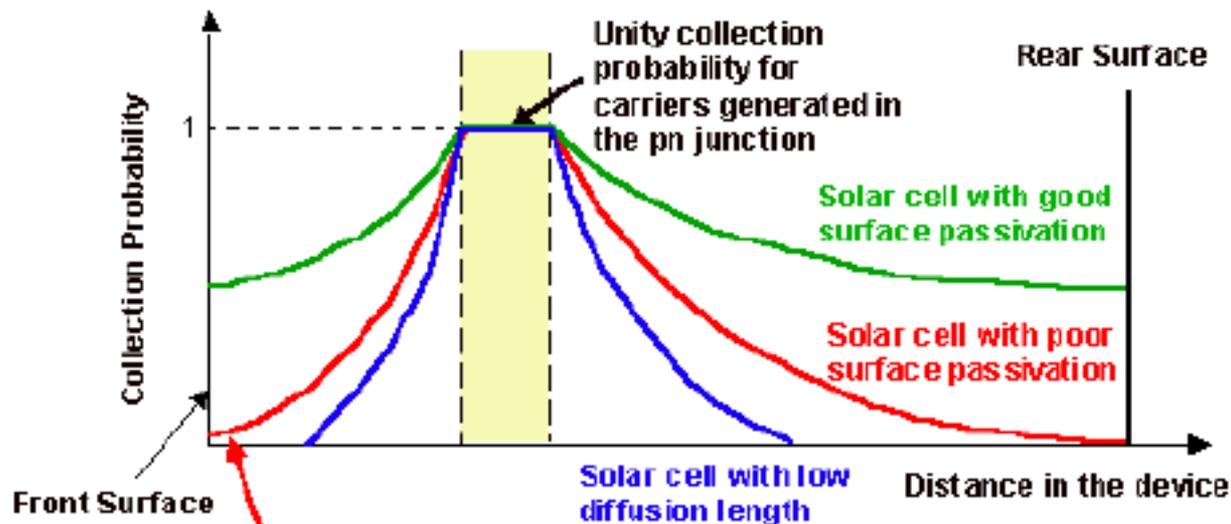


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 - Describe basic recombination mechanisms in semiconductor materials.
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4. Mobility:
 - Describe common mobility-limiting mechanisms (dopants, temperature, ionic semiconductors).

Collection Probability

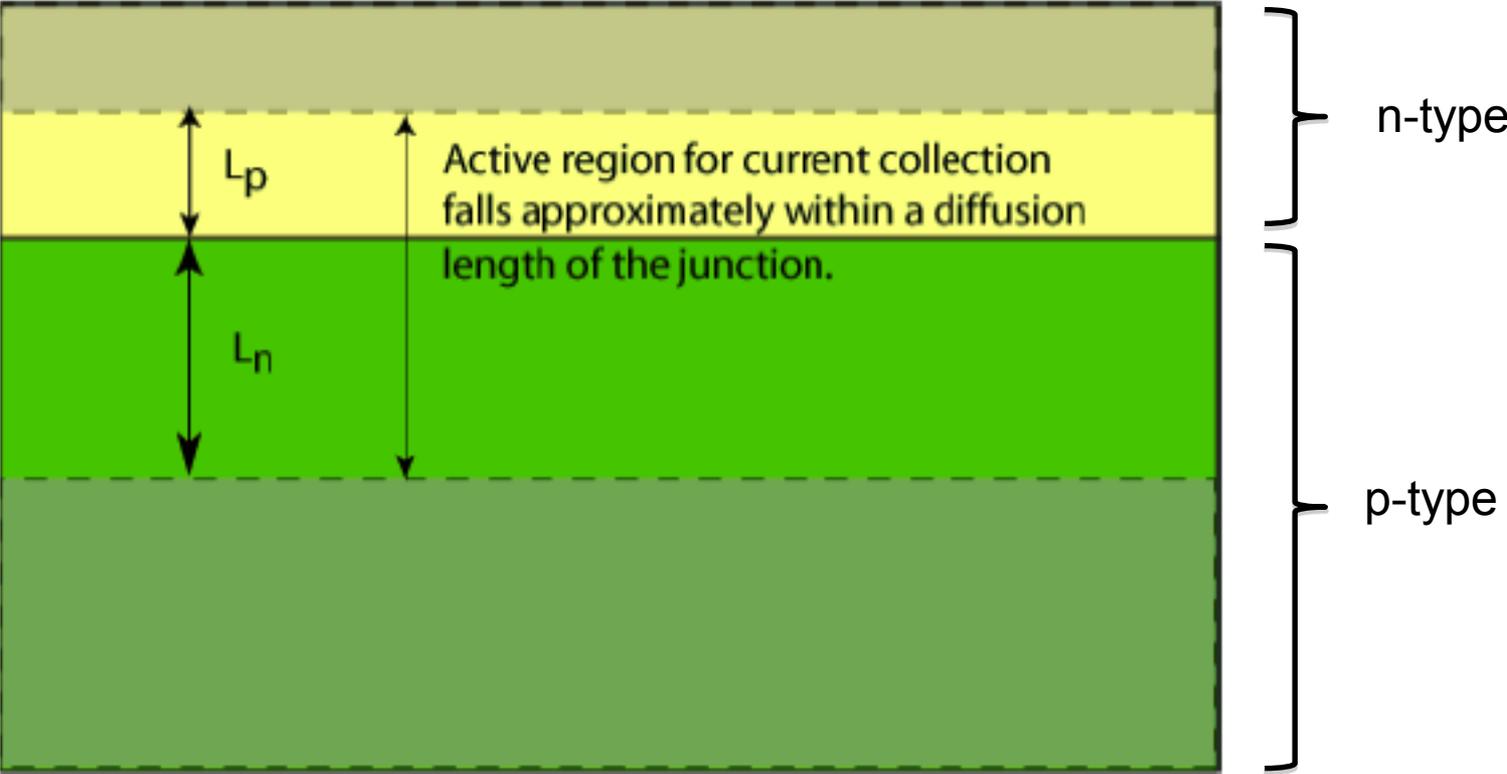
- Collection probability is the probability that a light generated carrier will reach the depletion region and be collected.
- Depends on where it is generated compared to junction and other recombination mechanisms, and the diffusion length.



With high surface recombination, the collection probability at the surface is low.

Courtesy of Christiana Honsberg. Used with permission.

Collection Probability



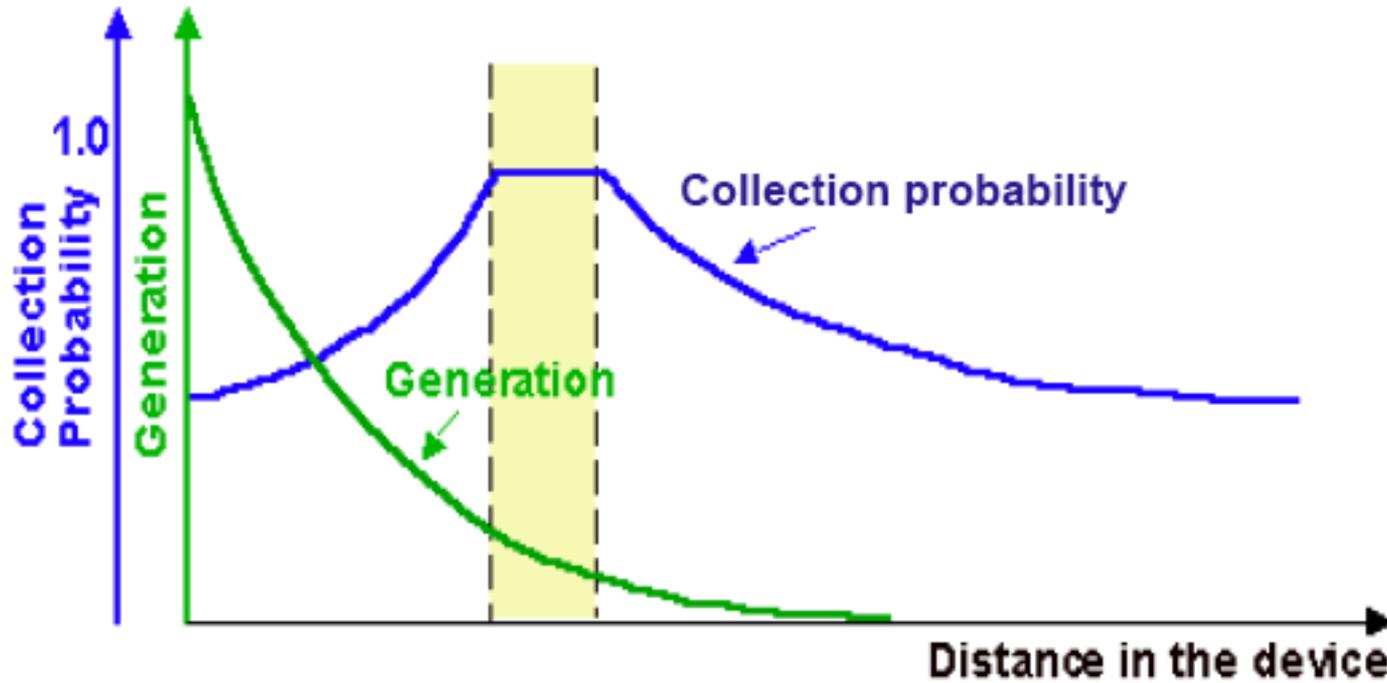
Collection probability is low further than a diffusion length away from junction

Courtesy of Christiana Honsberg. Used with permission.

Collection Probability

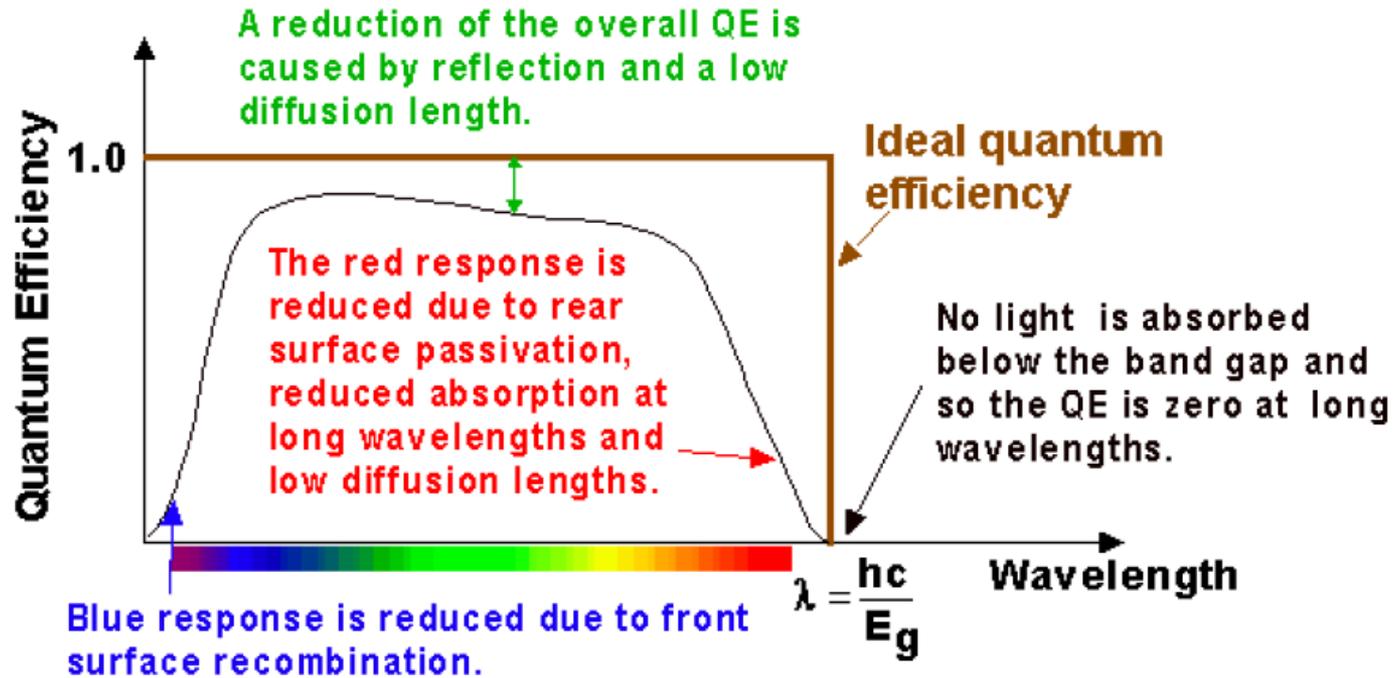
J_{sc} determined by generation rate and collection probability

$$J_L = q \int_0^W G(x) CP(x) dx$$



Courtesy of Christiana Honsberg. Used with permission.

Spectral Response (Quantum Efficiency)

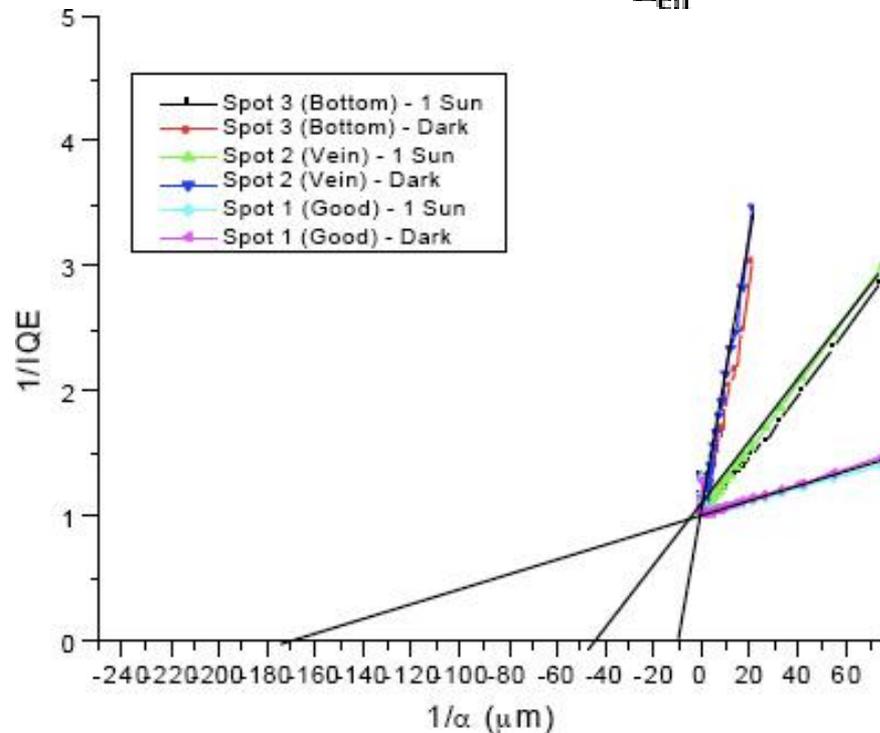


Courtesy of PVCDROM. Used with permission.

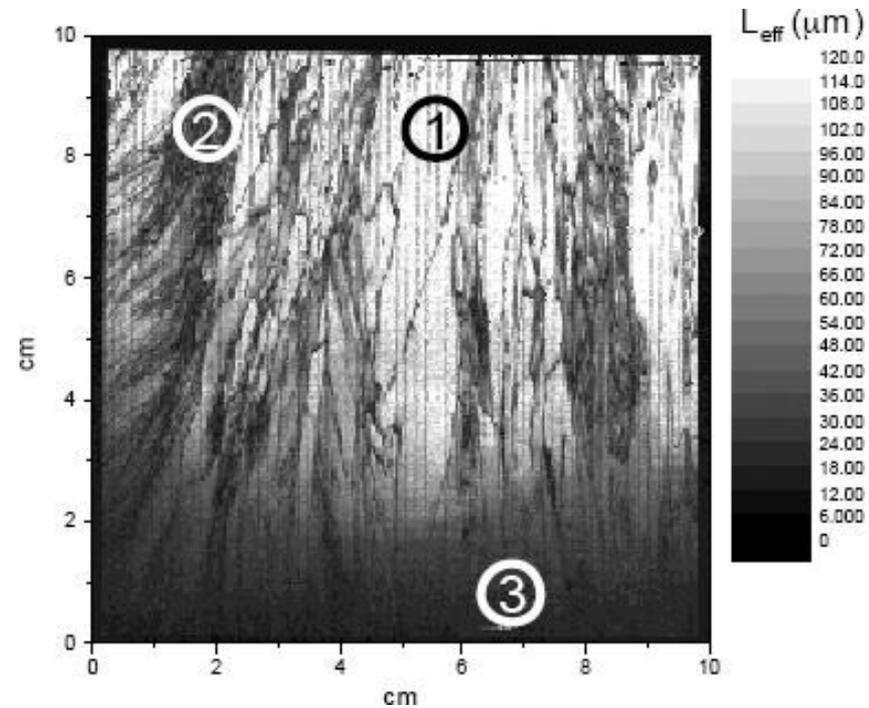
Minority Carrier Diffusion Length

At each point...

$$\text{IQE}^{-1} = 1 + \alpha^{-1} \frac{\cos \theta}{L_{\text{eff}}}$$



Mapped over an entire sample...



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See: P. A. Basore, *IEEE Trans. Electron. Dev.* **37**, 337 (1990).

What Limits Diffusion Length?

What Limits the Minority Carrier Diffusion Length?

Quick answer: Recombination-active defects, intrinsic mobility limitations, or absence of percolation pathways.

Thin films: Effect of bulk defects (GBs) on η .

Nanostructured: Effect of morphology on η .

Please see lecture video for visuals.

R.B. Bergmann, *Appl. Phys. A* **69**, 187 (1999)

F. Yang *et al.*, *ACS Nano* **2**, 1022 (2008)

Mathematical Formalism

Diffusion length is governed by “lifetime” and “mobility”

$$L_{\text{diff}} = \sqrt{D \tau_{\text{bulk}}}$$

$$D = \frac{k_B T}{q} \mu$$

L_{diff} = bulk diffusion length

D = diffusivity

τ_{bulk} = bulk lifetime

For carriers in an electric field

k_B = Boltzmann coefficient

T = temperature

q = charge

μ = mobility

Definition of “bulk lifetime”: The average time an excited carrier exists before recombining. (*i.e., temporal analogy to diffusion length.*)

Units = time.

τ_{bulk} of μs to ms is typical for indirect bandgap semiconductors, while τ_{bulk} of ns to μs is typical of direct bandgap semiconductors.

Definition of “mobility”: How easily a carrier moves under an applied field. (*i.e., ratio of drift velocity to the electric field magnitude.*)

Expressed in units of $\text{cm}^2/(\text{V}\cdot\text{s})$.

Mobilities of 10-100's cm^2/Vs typical for most crystalline semiconductors. Can be orders of magnitude lower for organic, amorphous, and ionic materials.

Mathematical Formalism

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τ_{bulk} = bulk lifetime

Carrier Lifetime τ_{bulk} = average time carriers in an electric field

k_B = Boltzmann coefficient

T = temperature

q = charge

μ = mobility

Carrier Mobility μ

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Excess Electron Carrier Concentration

Generally equal to
doping concentration

$$n = n_0 + \Delta n$$

$$\Delta n = G\tau$$

Excess Electron Carrier
Density

$$\left[\frac{\text{carriers}}{\text{cm}^3} \right]$$

Generation rate

$$\left[\frac{\text{carriers}}{\text{cm}^3 \cdot \text{sec}} \right]$$

Excited Carrier
Lifetime

[sec]

Minority, Majority Carriers in Silicon Under AM1.5G

$$\Delta n = G\tau = 10^{11} \text{ cm}^{-3}$$

AM1.5G $\sim 10^{16}$
 $\left[\frac{\text{carriers}}{\text{cm}^3 \cdot \text{sec}} \right]$

Excited Carrier
Lifetime
 $\sim 10\mu\text{s}$

if :

$$N_D = 10^{16} \text{ cm}^{-3}$$

then :

$$\Delta p \gg p_0 = \frac{n_i^2}{N_D} = 10^4 \text{ cm}^{-3}$$

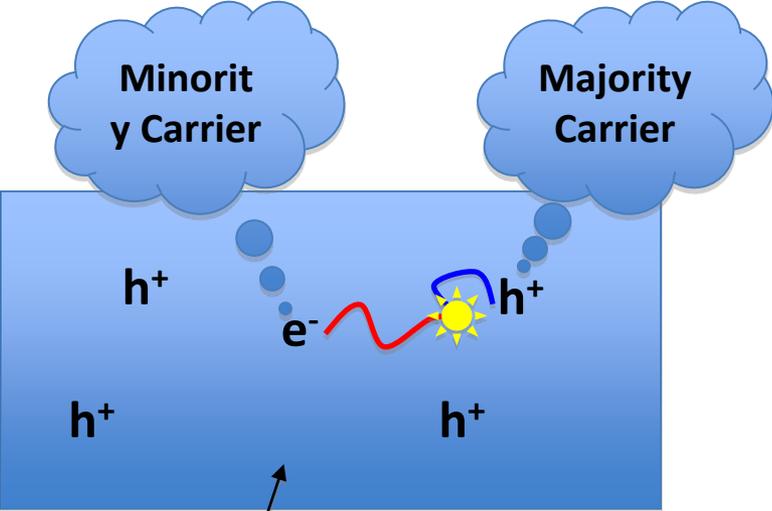
$$N_D \gg \Delta n$$

$$\Delta n = \Delta p = 10^{11} \text{ cm}^{-3}$$

$$n_i = 10^{10} \text{ cm}^{-3}$$

$$\Delta n = \Delta p > n_i$$

Carrier Lifetime and Recombination



Bulk Minority Carrier Lifetime:

$$\tau = \frac{\Delta n}{R}$$

Δn = Excess minority carrier concentration
 R = Recombination rate

First approximation of minority carrier lifetime, for low injection (e.g., illumination) conditions.

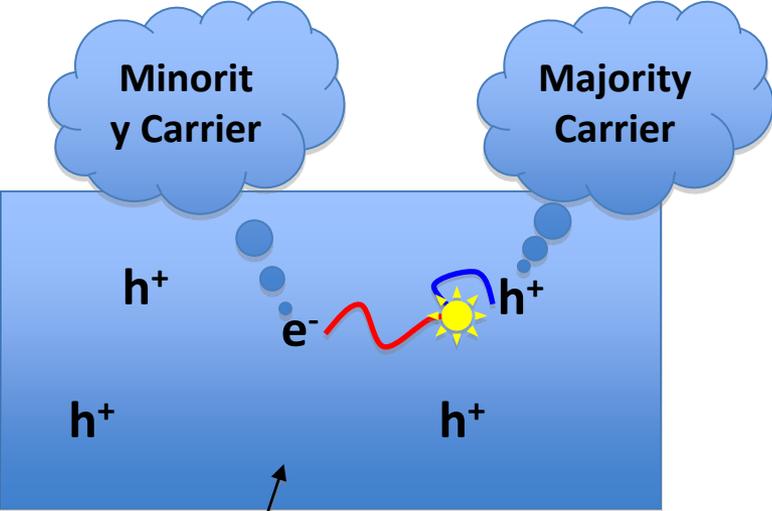
More Detailed Calculation:

$$\frac{1}{\tau_{\text{bulk}}} = \frac{1}{\tau_{\text{rad}}}$$

p-type silicon

Radiative recombination, can be derived from thermodynamics:
 $[\epsilon = \alpha]$

Carrier Lifetime and Recombination



p-type silicon

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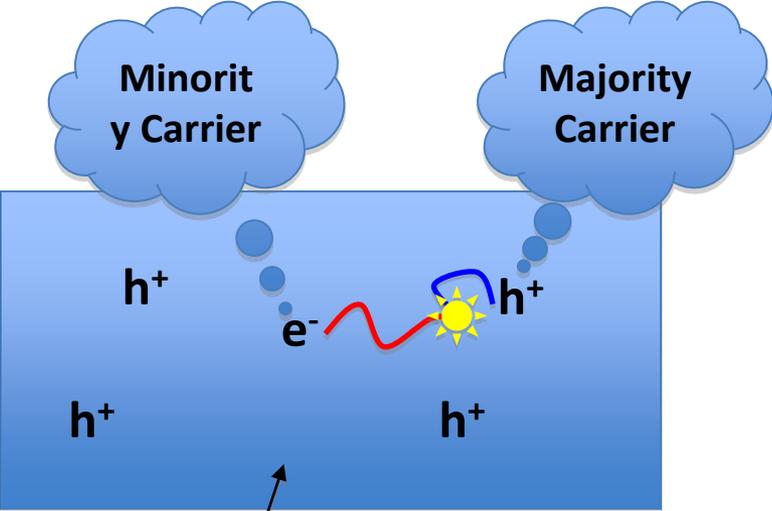
More Detailed Calculation:

$$\frac{1}{\tau_{\text{bulk}}} = \frac{1}{\tau_{\text{rad}}} + \frac{1}{\tau_{\text{Auger}}}$$

Dominant under very high injection conditions

$$\tau_{\text{Auger}} = \frac{1}{CN_A^2}$$

Carrier Lifetime and Recombination



Bulk Minority Carrier Lifetime:

$$\tau = \frac{\Delta n}{R}$$

Δn = Excess minority carrier concentration
 R = Recombination rate

First approximation of minority carrier lifetime, for low injection (e.g., illumination) conditions.

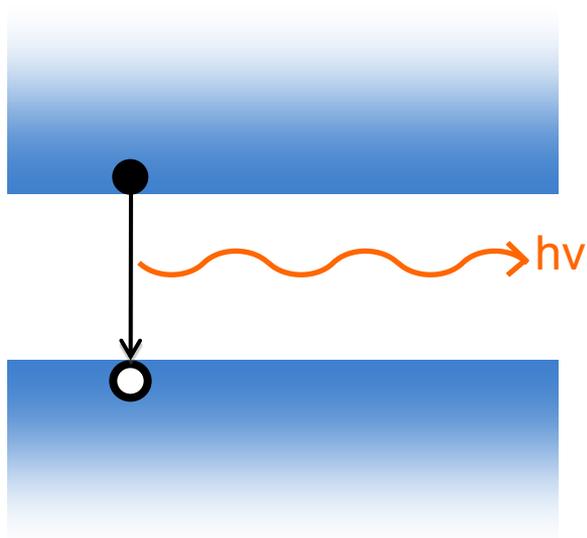
More Detailed Calculation:

$$\frac{1}{\tau_{\text{bulk}}} = \frac{1}{\tau_{\text{rad}}} + \frac{1}{\tau_{\text{Auger}}} + \frac{1}{\tau_{\text{SRH}}}$$



p-type silicon

Radiative Recombination



$$R = G = Bnp = Bn_i^2$$

Under equilibrium
dark conditions

$$R = B(np - n_i^2)$$

Net Recombination
non-equilibrium
 $= R - G_{\text{equilibrium}}$

$$n = n_0 + \Delta n$$

$$p = p_0 + \Delta p$$

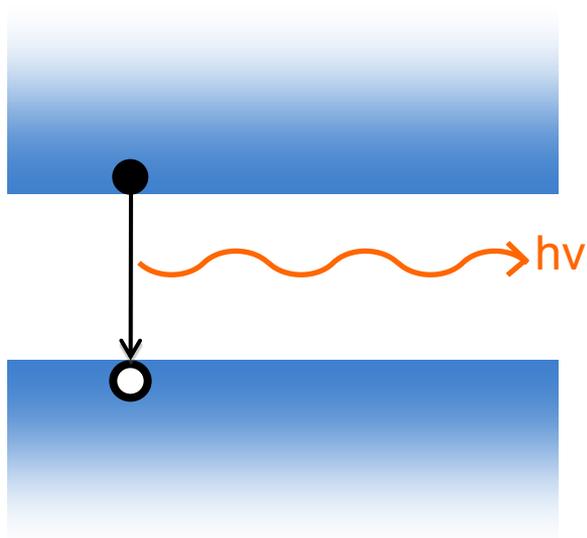
$$\tau_{\text{rad}} = \frac{1}{B(N_A + \Delta n)}$$

n-type (where $n_0 = N_D$)

$$\tau_{\text{rad}} = \frac{1}{B(N_D + \Delta p)}$$

p-type (where $p_0 = N_A$)

Radiative Recombination



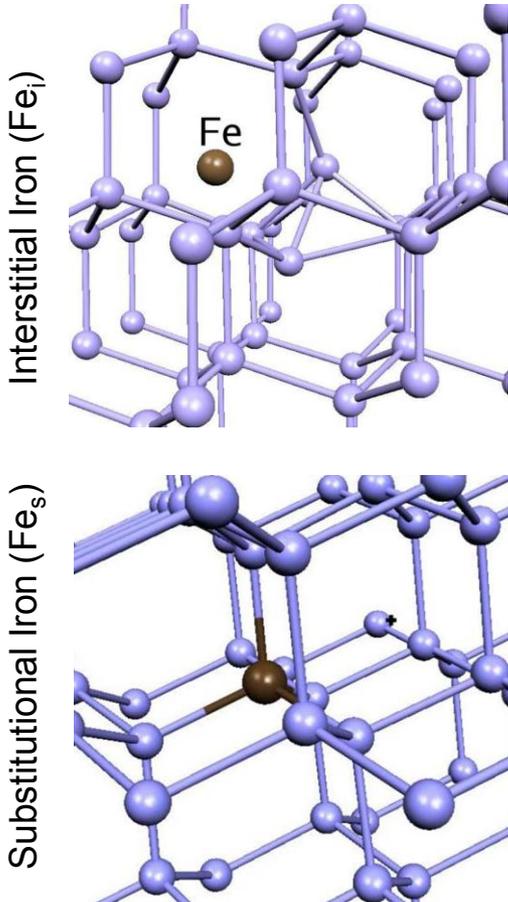
$$\tau_{rad, Si} = \frac{1}{\underbrace{(2 \times 10^{-15})}_B \underbrace{(10^{16})}_n} = 100ms$$

Radiative recombination is very slow in silicon, and is rarely the limiting lifetime in silicon-based solar cells. *However, radiative recombination is often the lifetime-limiting recombination pathway for high-quality “thin-film” materials, including GaAs.*

Defects and Carrier Recombination: τ_{SRH}

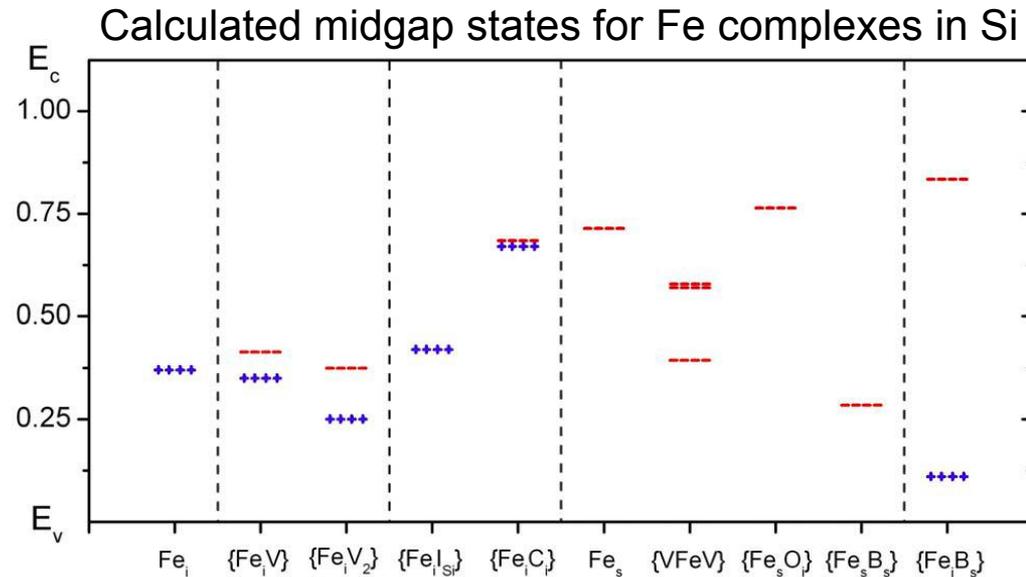
Defects can form in semiconductors.

Defect formation energy (ΔE_A) determines equilibrium concentration at a given temperature.



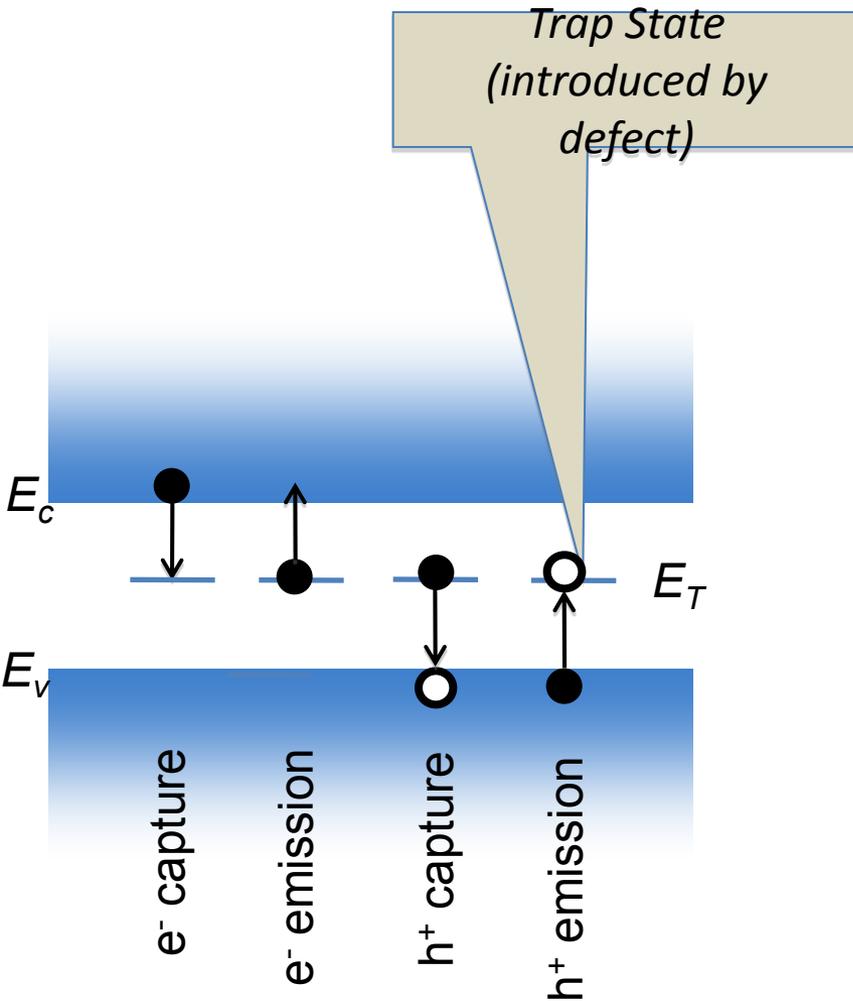
Defects can introduce midgap states.

Midgap states are characterized by their energy level(s) (E_D) and capture cross sections for electrons and holes (σ_e, σ_h)



S.K. Estreicher *et al.*, *Phys. Rev. B* **77**, 125214 (2008).

Defects and Carrier Recombination: τ_{SRH}



$$\tau_{SRH} = \frac{\tau_{n0}(n_0 + p_1 + \Delta n) + \tau_{p0}(p_0 + n_1 + \Delta n)}{n_0 + p_0 + \Delta n}$$

$$n_1 = N_c = e^{\left(\frac{E_T - E_c}{kT}\right)}$$

$$p_1 = N_v = e^{\left(\frac{E_v - E_T}{kT}\right)}$$

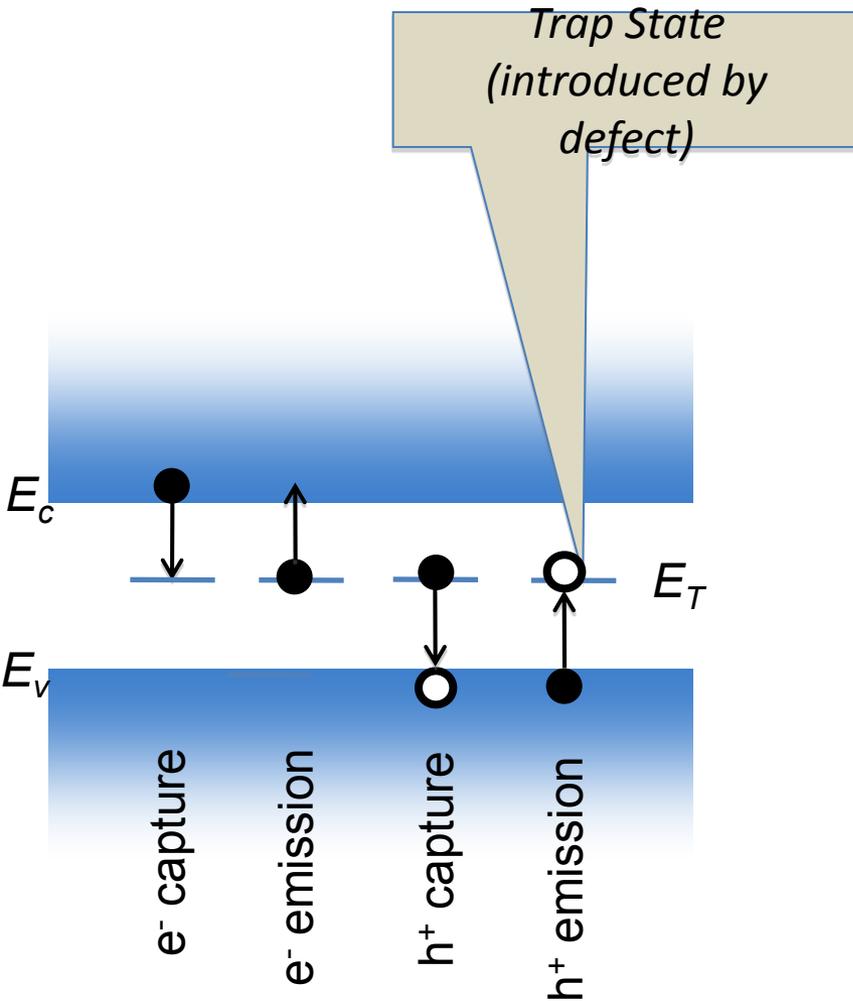
$$\tau_{p0} = \frac{1}{Nv_{th}\sigma_p}$$

$$\tau_{n0} = \frac{1}{Nv_{th}\sigma_n}$$

$N = \text{trap density}$

$\sigma = \text{capture cross-section}$

Defects and Carrier Recombination: τ_{SRH}



With deep traps (i.e. mid-gap) and low-injection conditions:

$$\tau_{SRH, n-type} = \tau_{p0} = \frac{1}{Nv_{th}\sigma_p}$$

$$\tau_{SRH, p-type} = \tau_{n0} = \frac{1}{Nv_{th}\sigma_n}$$

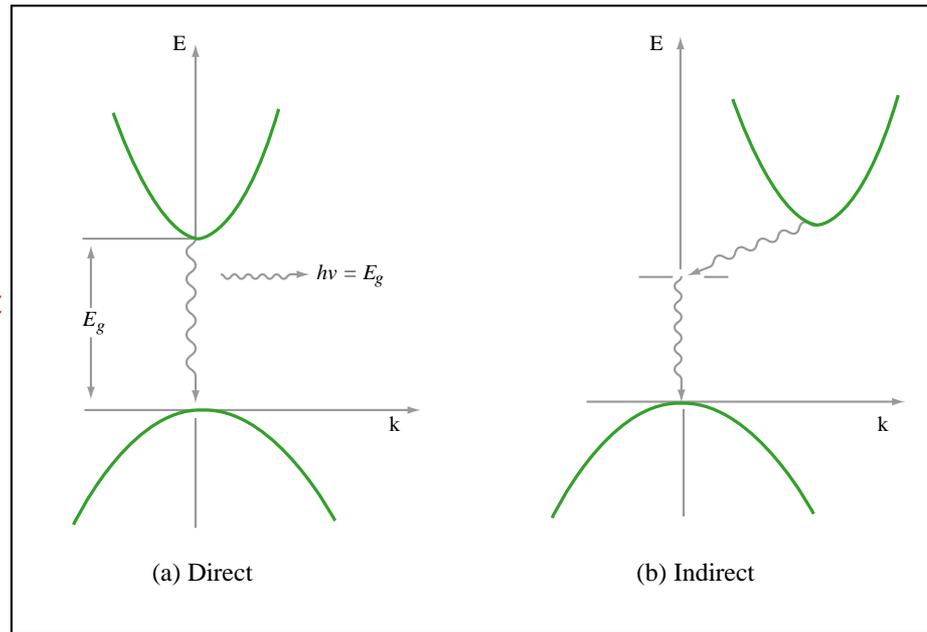
Deep traps and high-injection conditions:

$$\tau_{SRH} = \tau_{n0} + \tau_{p0}$$

Localized Defects Create Efficient Recombination Pathway

Direct Bandgap Semiconductor

Indirect Bandgap Semiconductor



Recombination efficient
(no phonon required)
 $\tau \sim \text{ns to } \mu\text{s}$

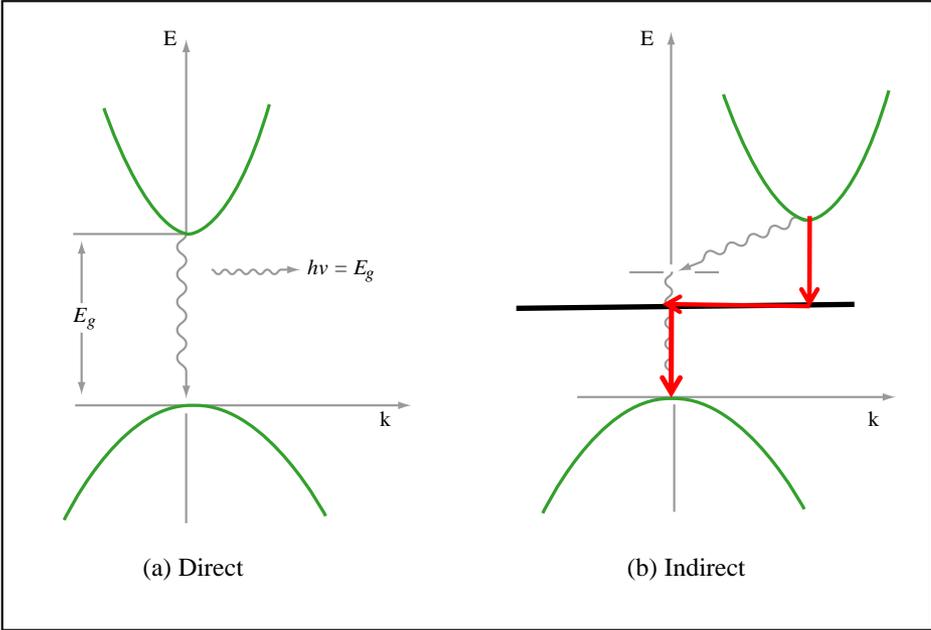
Recombination inefficient
(phonon required)
 $\tau \sim \text{ms}$

Image by MIT OpenCourseWare.

Localized Defects Create Efficient Recombination Pathway

Direct Bandgap Semiconductor

Indirect Bandgap Semiconductor



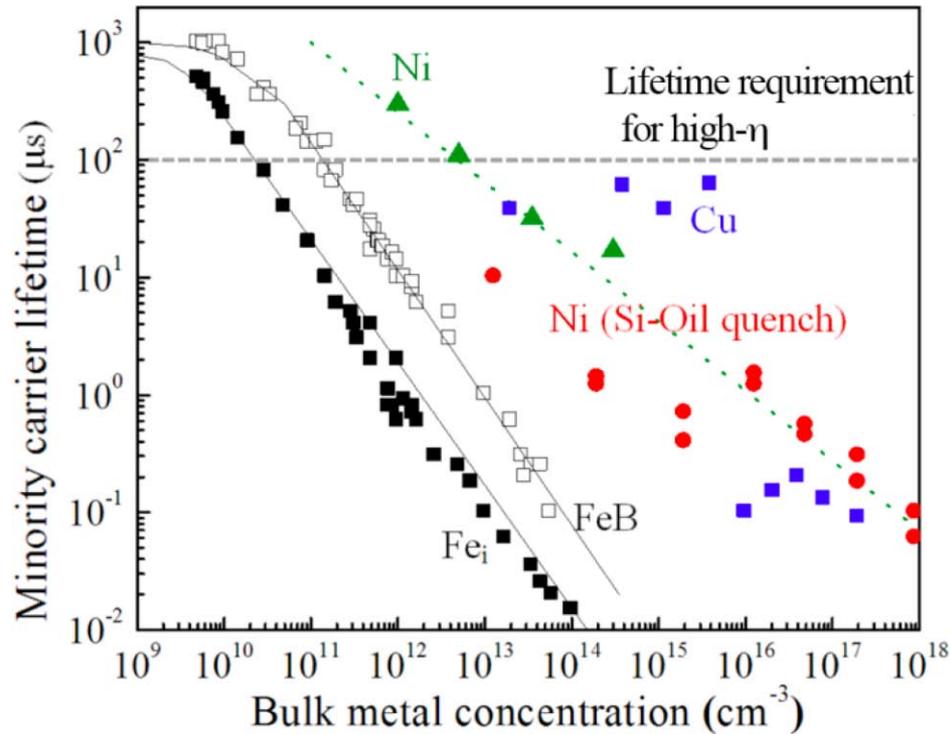
Recombination efficient
(no phonon required)
 $\tau \sim \text{ns to } \mu\text{s}$

Efficient recombination
via defect level!
 $\tau < \mu\text{s}$ with high
defect concentrations

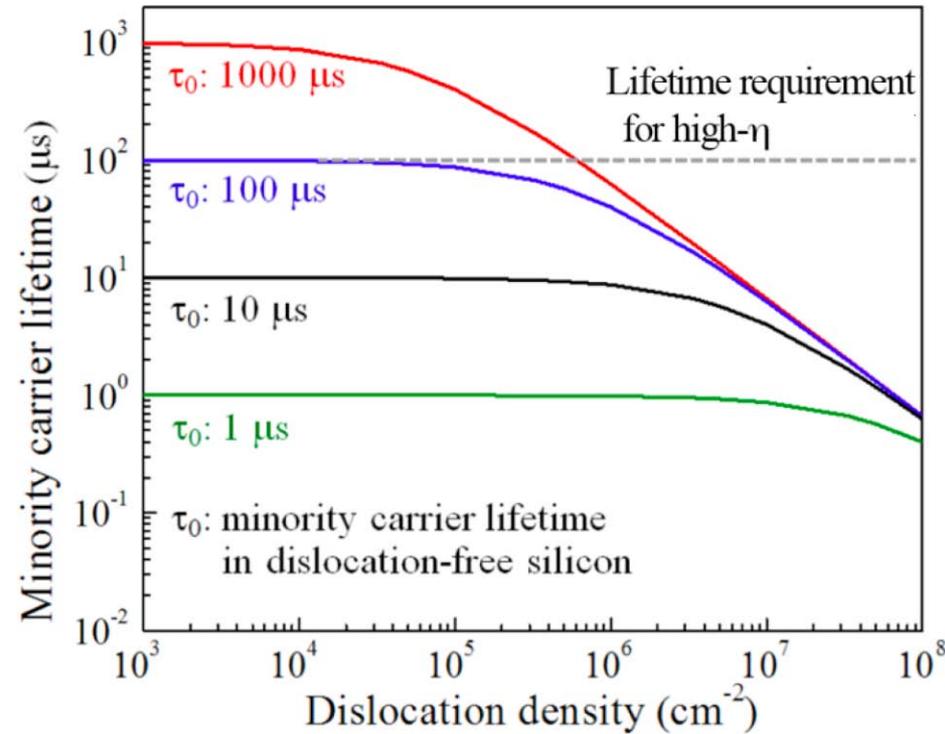
Image by MIT OpenCourseWare.

Defects Impact Minority Carrier Lifetime

Impurity Point Defects (c-Si)



Dislocations (c-Si)



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A.A. Istratov, *Materials Science & Engineering B*, **134**, 282 (2006)

C. Donolato, *Journal of Applied Physics* **84**, 2656 (1998)

Effect of Defects on Minority Carrier τ , L_{diff}

The distribution of defects matters as much as their total concentration!

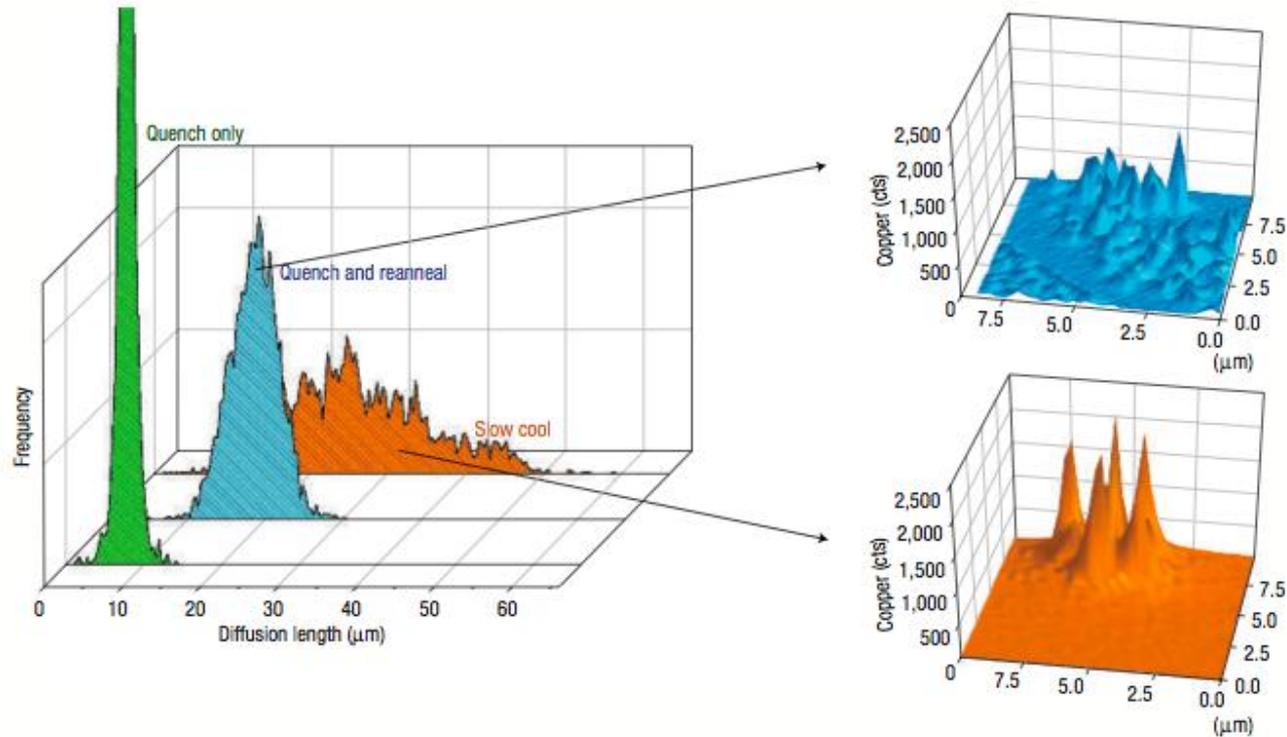
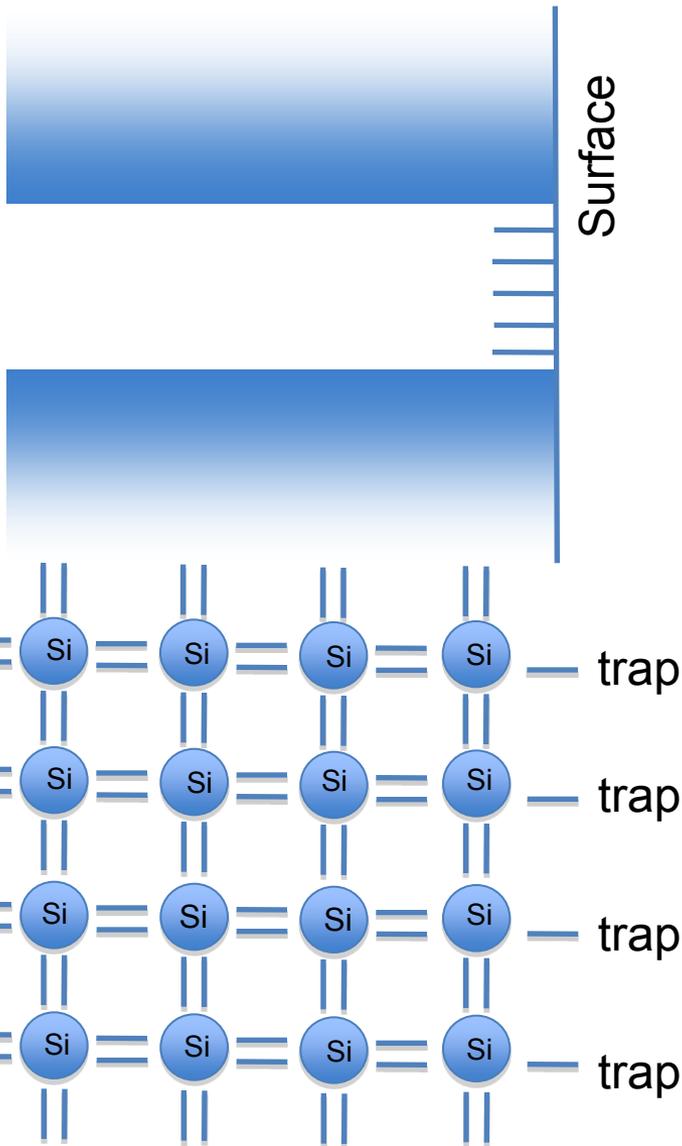


Figure 3 Effect of the distribution of metal defects on material performance. Material performance (minority-carrier diffusion length histograms, left) in three differently cooled samples (quench, quench and re-anneal, slow cool) is compared with size and spatial distributions of metal defects (high-resolution μ -XRF maps (right), XRF copper counts per second plotted against x and y coordinates in μm). The material with microdefects in lower spatial densities clearly outperforms materials with smaller nanodefects in higher spatial densities, despite the fact that all materials contain the same total amount of metals.

Reprinted by permission from Macmillan Publishers Ltd: Nature Materials.
Source: Buonassisi, T., et al. "Engineering Metal-Impurity Nanodefects for Low-Cost Solar Cells." *Nature Materials* 4, no. 9 (2005): 676-9. © 2005.

T. Buonassisi *et al.*, *Nature Mater.* 4, 676 (2005)

Surfaces Introduce Many Mid-Gap States



mid-gap surface states provide additional pathway for recombination.

They are often formed by dangling bonds on the surface.

$$\tau_{surf} \approx \frac{W}{2S} + \frac{1}{D} \left(\frac{W}{\pi} \right)^2$$

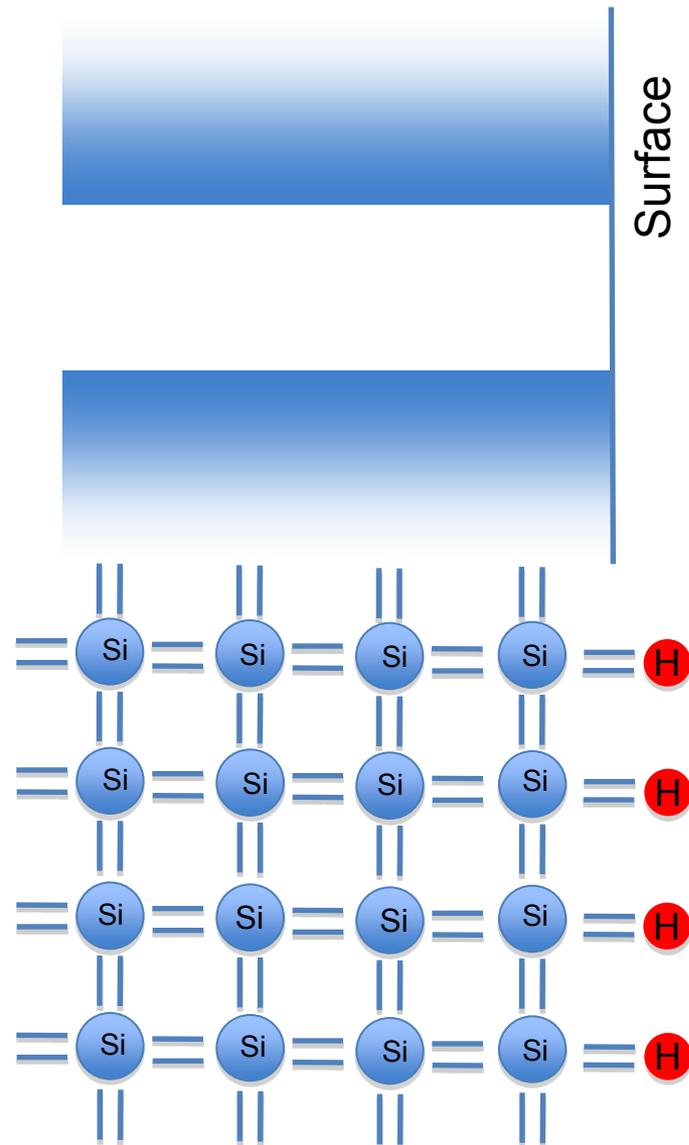
Sample thickness

Surface recombination velocity

Carrier diffusivity

Ref: Horányi et al., *Appl. Surf. Sci.* **63**, 306 (1993)

Proper Passivation Reduces Surface Recombination



Proper surface passivation ties up dangling bonds, reduces density of trap states.

“Perfect” passivation yields:

$$S \rightarrow 0, \quad \tau_{surf} \rightarrow \infty$$

Practically, S approaching 1 cm/s have been achieved on silicon.

Measuring Surface Recombination Velocity

1. Measure lifetime of samples of varying thickness, but with same bulk and surface properties. (The lifetime will be affected by both bulk and surface conditions, thus measurement will reveal an “effective” lifetime, or τ_{eff} .)

$$\frac{1}{\tau_{eff}} = \frac{1}{\tau_{bulk}} + \frac{1}{\tau_{surf}}$$

2. Any variation in lifetime should be due to a changing τ_{surf} , per below:

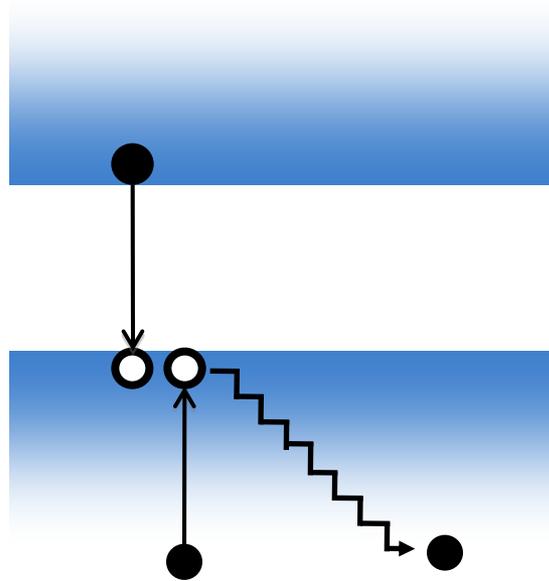
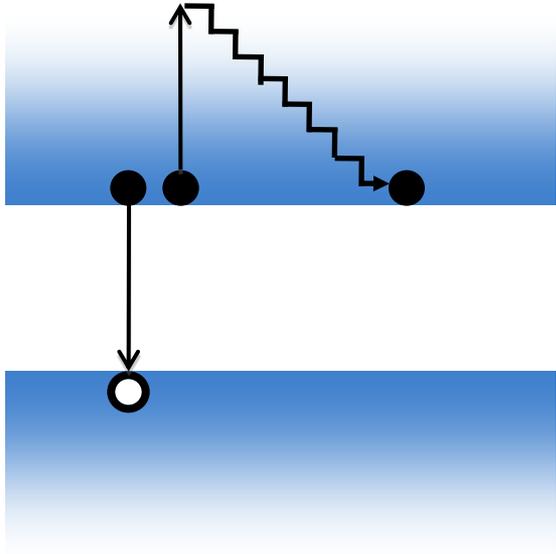
The diagram shows the equation $\tau_{surf} \approx \frac{W}{2S} + \frac{1}{D} \left(\frac{W}{\pi} \right)^2$. A callout box labeled "Sample thickness" points to the variable W . Another callout box labeled "Surface recombination velocity" points to the variable S . A third callout box labeled "Carrier diffusivity" points to the variable D .

$$\tau_{surf} \approx \frac{W}{2S} + \frac{1}{D} \left(\frac{W}{\pi} \right)^2$$

Ref: Horányi et al.,
Appl. Surf. Sci. **63**,
306 (1993)

3. With >2 lifetime measurements (sample thicknesses), can solve for τ_{surf} !

Auger Recombination, τ_{Auger}



$$R_{n\text{-type}} \sim pn^2$$

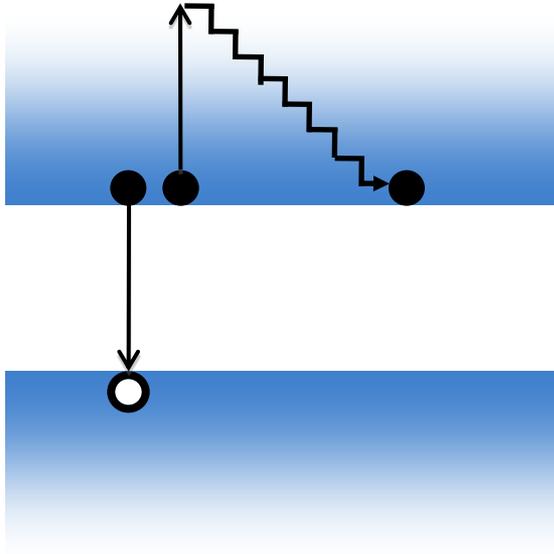
$$\tau_{n\text{-type}} \approx \frac{1}{Cn^2}$$

$$R_{p\text{-type}} \sim np^2$$

$$\tau_{p\text{-type}} \approx \frac{1}{Cp^2}$$

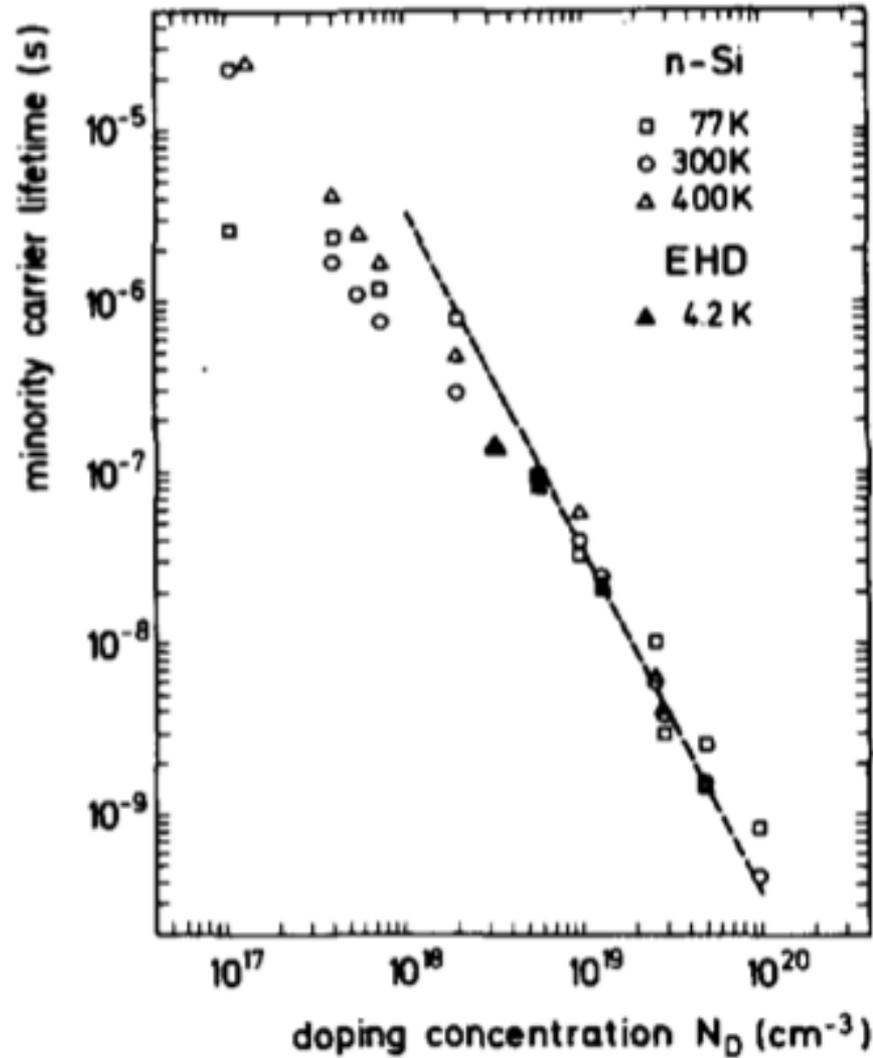
*Above equations true at sufficiently high doping densities ($\geq 10^{18} \text{cm}^{-3}$ in silicon)

Auger Recombination in *n*-type silicon



$$R_{n\text{-type}} \sim pn^2$$

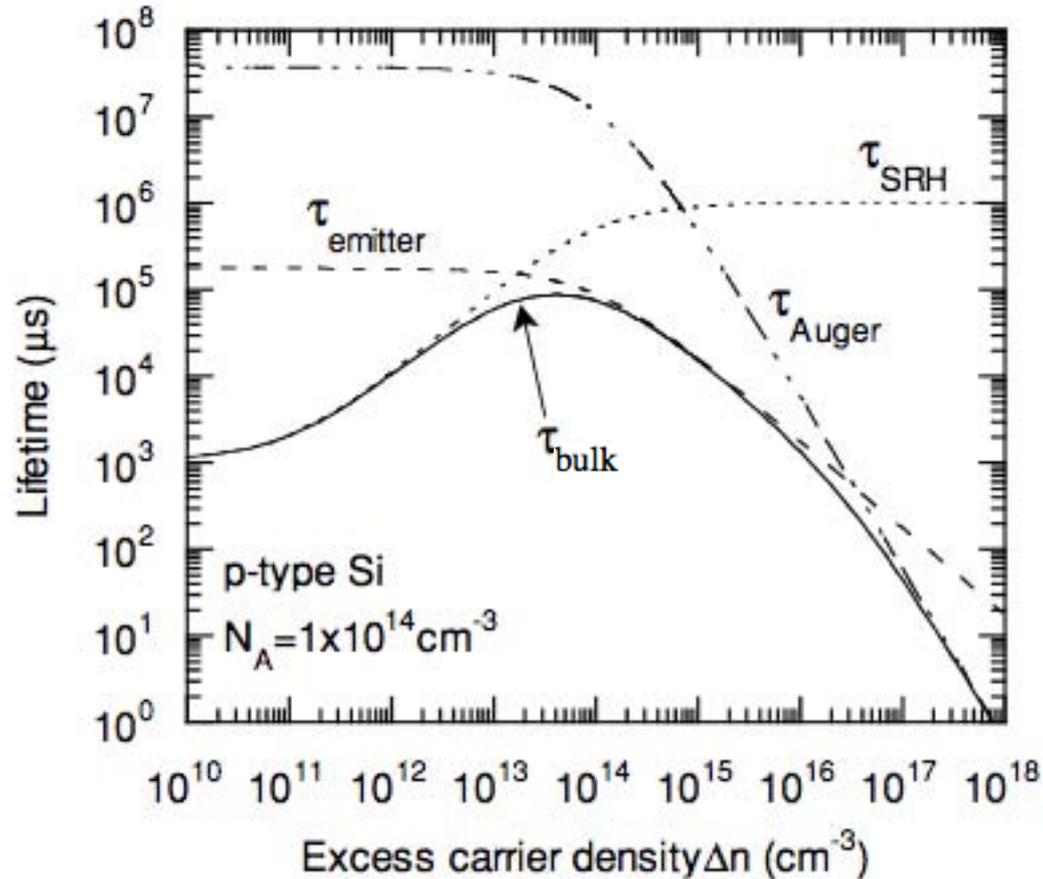
$$\tau_{n\text{-type}} \approx \frac{1}{Cn^2}$$



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J. Dziewior, W. Shmid, *Appl. Phys. Lett.*, **31**, p346 (1977)

Auger Recombination in *n*-type silicon



Courtesy of Daniel Macdonald. Used with permission.

$$\frac{1}{\tau_{\text{bulk}}} = \frac{1}{\tau_{\text{band}}} + \frac{1}{\tau_{\text{Auger}}} + \frac{1}{\tau_{\text{SRH}}}$$

Measuring Minority Carrier Lifetime

Recall that

$$\frac{1}{\tau_{\text{bulk}}} = \frac{1}{\tau_{\text{rad}}} + \frac{1}{\tau_{\text{Auger}}} + \frac{1}{\tau_{\text{SRH}}}$$

If defect-mitigated recombination is dominant, band-to-band radiative recombination will be suppressed:

$$\frac{1}{\tau_{\text{rad}}} \gg \frac{1}{\tau_{\text{SRH}}}$$

OR

$$\tau_{\text{rad}} \ll \tau_{\text{SRH}}$$

In fact, band-to-band (radiative recombination) and defect-mitigated (non-radiative recombination) are inversely proportional.

By imaging the band-to-band (radiative) recombination using a very sensitive CCD camera, we are able to quantitatively extract the minority carrier lifetime.

Photoluminescence Imaging (PLI)

Experimental Setup

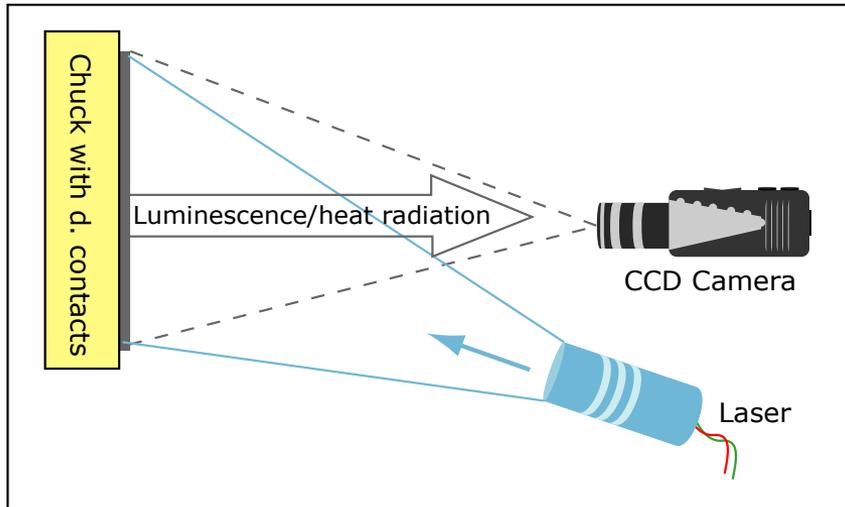


Image by MIT OpenCourseWare.

Measurement of Wafer

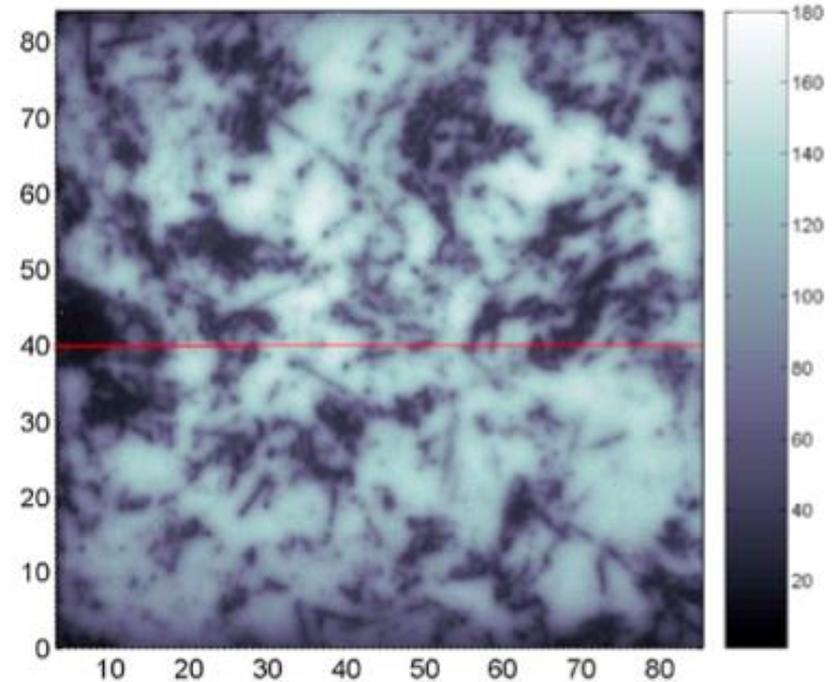


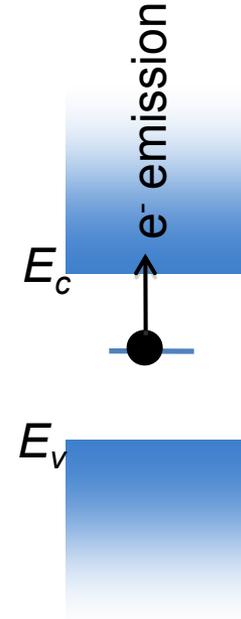
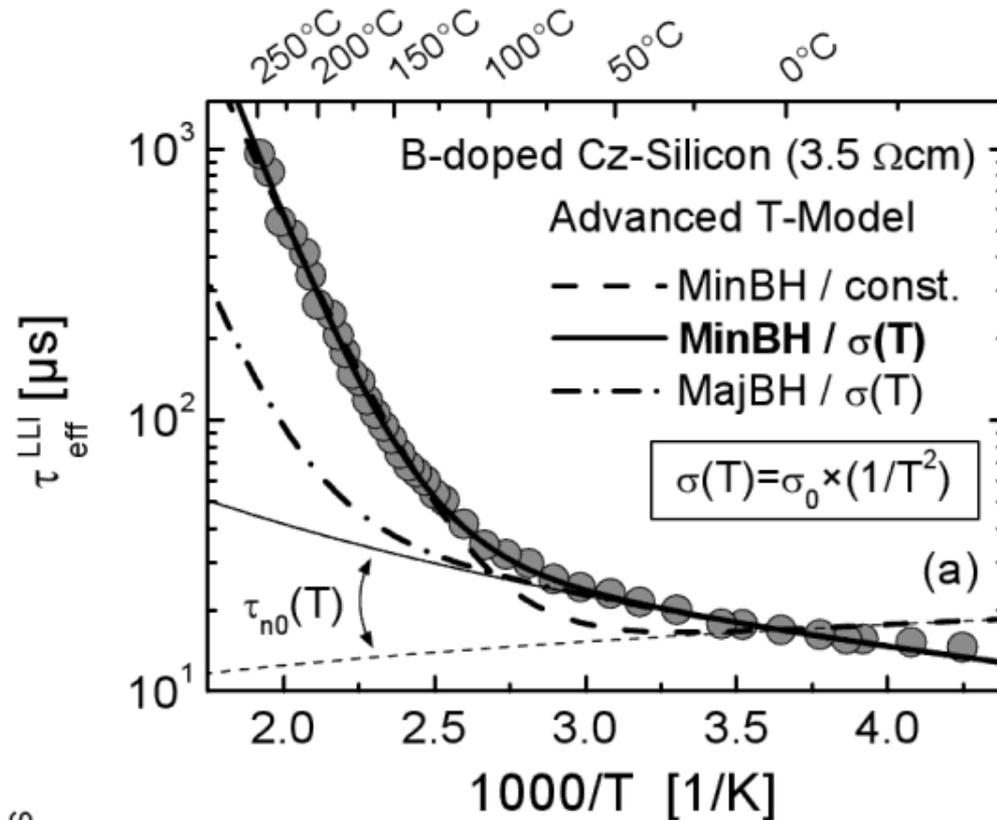
FIG. 1. (Color online) Lifetime distribution within an $8.5 \times 8.5 \text{ cm}^2$ area of a $302 \text{ }\mu\text{m}$ thick, $1.2 \text{ }\Omega \text{ cm}$, mc-Si, *p*-type wafer obtained from a PL image measured with a data acquisition time of 1.5 s and with a spatial resolution of $130 \text{ }\mu\text{m}$.

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M. Kasemann *et al.*, *Proc. IEEE PVSC*, San Diego, CA (2008).

T. Trupke *et al.*, *Appl. Phys. Lett.* **89**, 044107 (2006).

Temperature Changes σ_0



Trapped carriers are more likely to be (re-)emitted at higher thermal energies ($k_b T$). At lower T , trapped carriers reside in traps longer, facilitating recombination.

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Learning Objectives: Toward a 1D Device Model

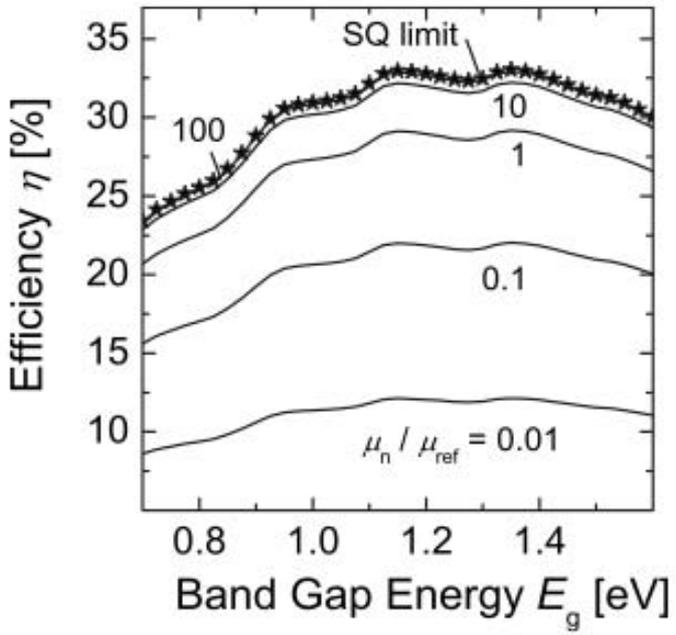
1. Describe what minority carrier diffusion length is, and calculate its impact on J_{sc} , V_{oc} . Describe how minority carrier diffusion length is affected by minority carrier lifetime and minority carrier mobility.
2. Describe how minority carrier diffusion length is measured.
3. Lifetime:
 - Describe basic recombination mechanisms in semiconductor materials.
 - Calculate excess carrier concentration as a function of carrier lifetime and generation rate. Compare to background (intrinsic + dopant) carrier concentrations.
4. **Mobility:**
 - **Describe common mobility-limiting mechanisms (dopants, temperature, ionic semiconductors).**

Effect of Reduced Mobility on Solar Cell Performance

Low carrier mobility can reduce device efficiency by several tens of percent relative (*as big of an impact as lifetime!*)

$$L_{\text{diff}} = \sqrt{D \tau_{\text{bulk}}}$$

$$D = \frac{k_B T}{q} \mu$$



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J. Mattheis *et al.*, *Phys. Rev. B.* **77**, 085203 (2008)

What Limits Mobility?

1. Defect Scattering
2. Trapping at stretched bonds
3. Incomplete percolation pathways
4. Phonon Scattering

Example of carrier trapping

L. Wagner *et al.*, *PRL* **101**, 265501 (2008)

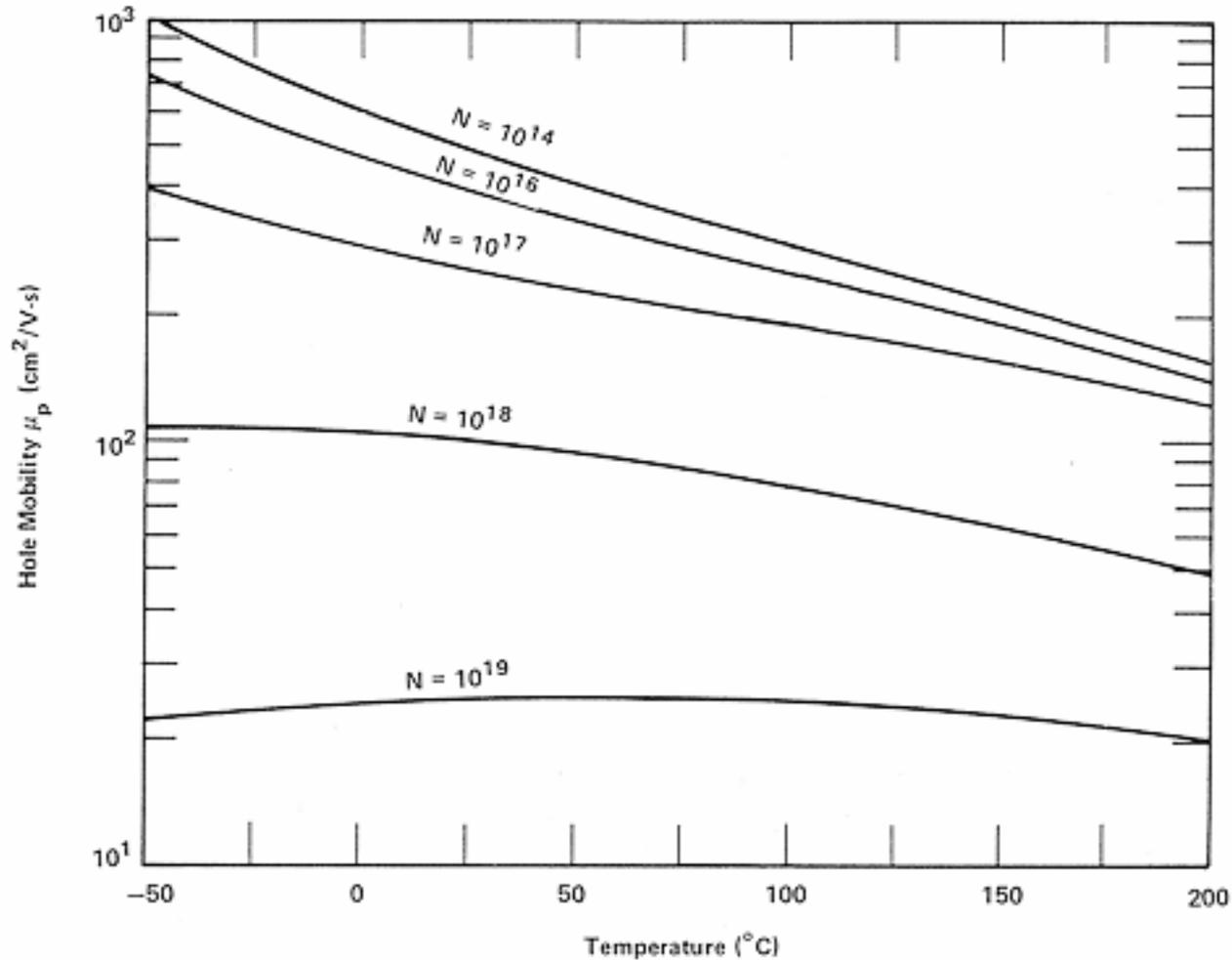
Example of complex percolation pathway

F. Yang *et al.*, *ACS Nano* **2**, 1022 (2008)

Diagrams removed due to copyright restrictions.
See lecture video.

Mobility and Carrier Concentration in Semiconductor

Increased concentration of ionized dopant atoms increases conductivity, but can reduce carrier mobility (due to scattering).



Percolation Pathways

Both μ and τ play a role in determining L_{diff} , and hence efficiency, for various device architectures.

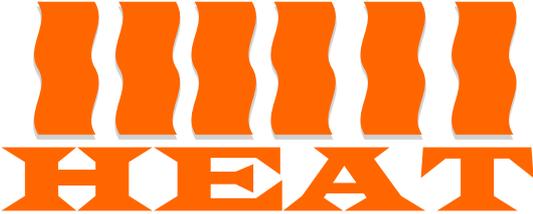
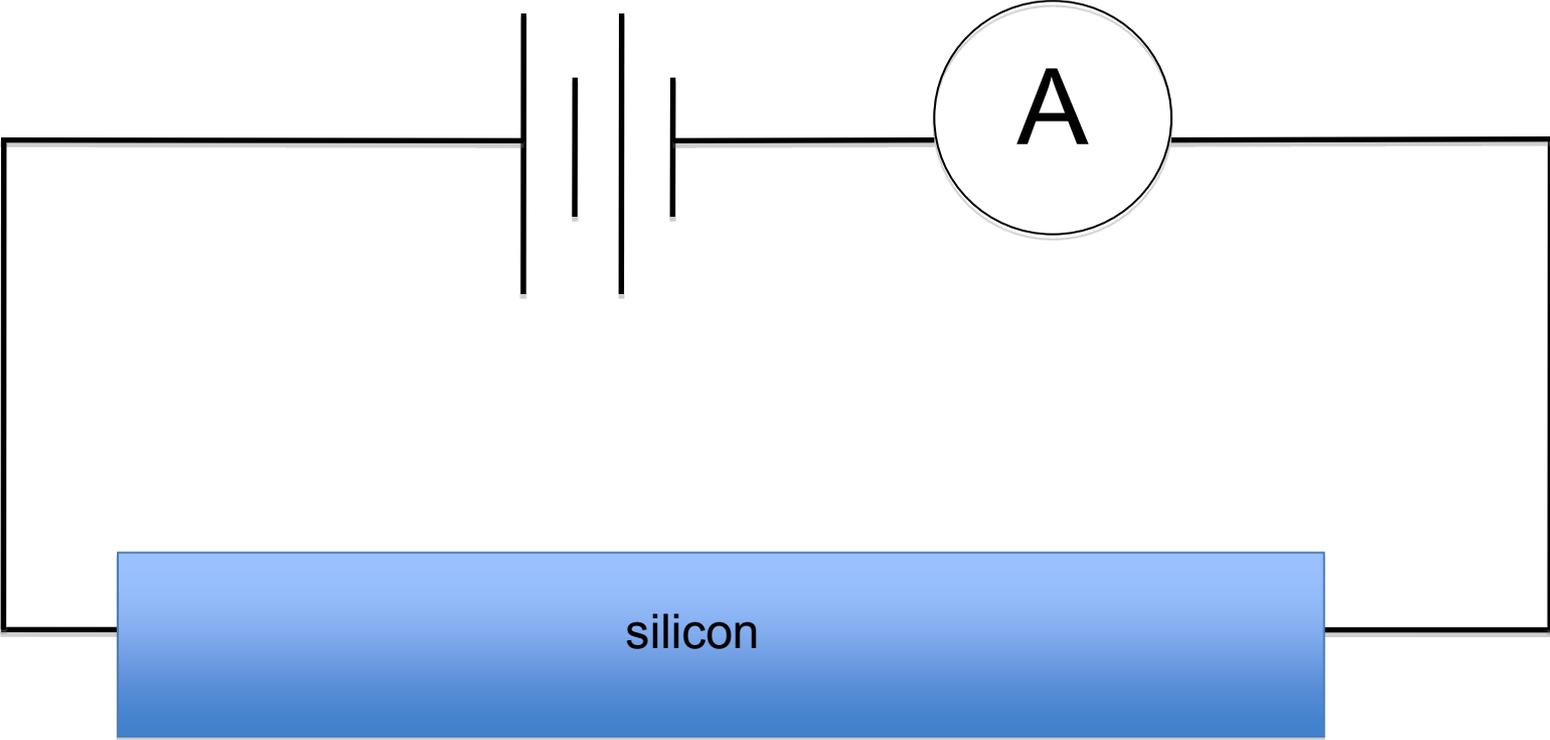
Table removed due to copyright restrictions.
See lecture video.

F. Yang *et al.*, *ACS Nano* **2**, 1022 (2008)

Simple Example

1. Well-behaved inorganic semiconductor (no effect of trapping at stretched bonds and incomplete percolation pathways).
2. Defect scattering dominant \rightarrow ionized dopants inside sample scatter carriers.

Demo: Conductivity of Heated Intrinsic and Doped Silicon



Demo Explained

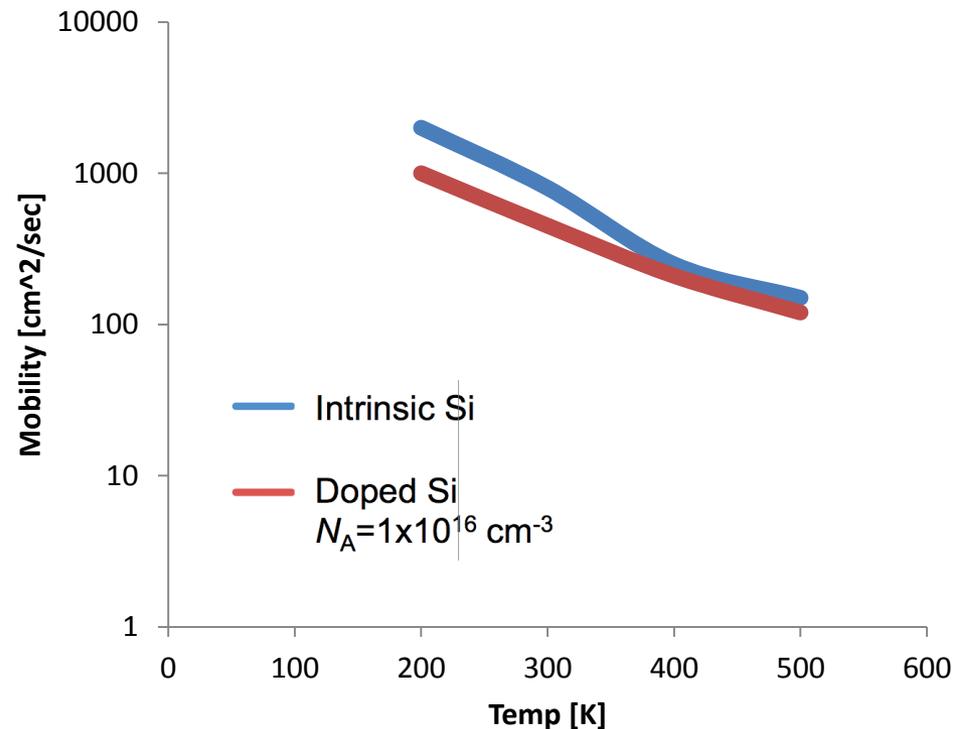
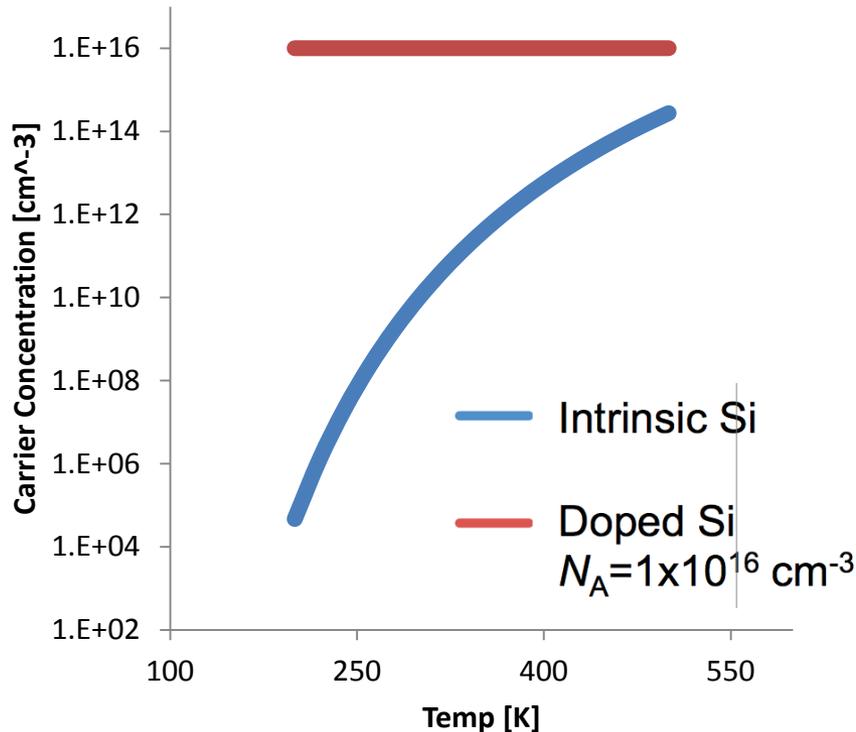
$$n = n_0 + \Delta n \approx N_D + \Delta n$$

For intrinsic Si, $\Delta n \gg N_D$.

For highly doped Si, $\Delta n \ll N_D$.

$$\sigma = \frac{1}{\rho} = q\mu n$$

Conductivity is the product of μ and n .



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