

**MIT Department of Mechanical Engineering**  
**2.25 Advanced Fluid Mechanics**

**Stokes Second Problem ATP**

Stokes apparently had many problems. This Second Problem is identical to the First Problem, except that we replace (2) with  $u(y = 0, t) = U \cos(\omega t)$  — the plate now oscillates. Note that we are interested only in the steady periodic solution:  $u$  behaves as  $\cos(\omega t + \Phi^u)$  in time, where the phase  $\Phi^u$  is independent of  $t$ . (The initial condition (4) is thus irrelevant —it washes out.)

In the steady-periodic state the wall shear stress will be of the form

$$\tau_W = CU^{\alpha_1} \rho^{\alpha_2} \mu^{\alpha_3} \omega^{\alpha_4} \cos(\omega t + \Phi^\tau), \quad (1)$$

where the phase  $\Phi^\tau$  is independent of  $t$  and  $C$  is a non-dimensional constant. Find the exponents  $\alpha_1$ ,  $\alpha_2$ ,  $\alpha_3$  and  $\alpha_4$  by dimensional analysis.

*Hint:* (one approach): See Hint for Stokes' First Problem; make good use of the steady-periodic form of the solution.

**Solution:**

Let's start by non dimensionalizing the equations. Now write  $u^* = \frac{u}{U}$ ; thus divided by ( $U$ ).

$$\frac{\partial u^*}{\partial t} = \nu \frac{\partial^2 u^*}{\partial y^2} \quad 0 < y < \infty, \quad (2)$$

$$u^*(y = 0, t) = \cos(\omega t), \quad (3)$$

$$u^*(y \rightarrow \infty, t) \rightarrow 0, \quad (4)$$

$$u^*(y, t = 0) = 0, \quad (5)$$

and hence,

$$u^* = f(y, t, \nu, \omega), \quad (6)$$

(notice that no mass appears in the equations) so,

$$\Pi_1 = u^*, \quad (7)$$

$$\Pi_2 = ty^{\alpha_2} \omega^{\beta_2}, \quad (8)$$

$$\Pi_2 = \omega t, \quad (9)$$

$$\Pi_3^a = \nu y^{\alpha_3} \omega^{\beta_3}, \quad (10)$$

Solving the system of equations,

$$\alpha_3 = -2, \quad \beta_3 = -1, \quad (11)$$

then,

$$\Pi_3^a = \frac{\nu}{y^2 \omega}, \quad (12)$$

or,

$$\Pi_3 = \frac{y}{\sqrt{\frac{\nu}{\omega}}}. \quad (13)$$

Then, reexpressing the original function in terms of the non-dimensional parameters,

$$\Pi_1 = f_*(\Pi_2, \Pi_3), \quad (14)$$

or,

$$\frac{u}{U} = f_{**}\left(\frac{y}{\sqrt{\frac{\nu}{\omega}}}, \omega t\right), \quad (15)$$

Now, for the steady periodic behaviour  $\frac{u}{U}$  must be 'sinusoidal' in time, so

$$\frac{u}{U} = A\left(\frac{y}{\sqrt{\frac{\nu}{\omega}}}\right) \cos\left(\omega t + \Phi\left(\frac{y}{\sqrt{\frac{\nu}{\omega}}}\right)\right) \quad (16)$$

where,

$$A(0) = 1 \quad \text{and} \quad \Phi(0) = 0. \quad (17)$$

Furthermore,

$$\tau_W = -\mu \frac{\partial u}{\partial y} \Big|_{y=0} \quad (18)$$

$$\tau_W = -\mu U \left( A'(0) \frac{1}{\sqrt{\frac{\nu}{\omega}}} \cos(\omega t + \Phi(0)) - A(0) \frac{1}{\sqrt{\frac{\nu}{\omega}}} \sin(\omega t + \Phi(0)) \Phi'(0) \right) \quad (19)$$

$$\tau_W = -\mu U \frac{1}{\sqrt{\frac{\nu}{\omega}}} \left( A'(0) \cos(\omega t + \Phi(0)) - \Phi'(0) \sin(\omega t + \Phi(0)) \right). \quad (20)$$

Finally, the last equation can be reexpressed as:

$$\tau_W = -\mu U \frac{1}{\sqrt{\frac{\nu}{\omega}}} \left( A'(0)^2 + \Phi'(0)^2 \right)^{0.5} \cos(\omega t + \Phi(0)^\tau). \quad (21)$$

where,  $\Phi_0^\tau = \arctan \frac{\Phi'(0)}{A'(0)}$ ; and  $A'(0)$  and  $\Phi'(0)$  are 'universal constants'.

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