

Elements of Polymer Structure and Viscoelasticity

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Mechanics and Materials II 2.002

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Outline

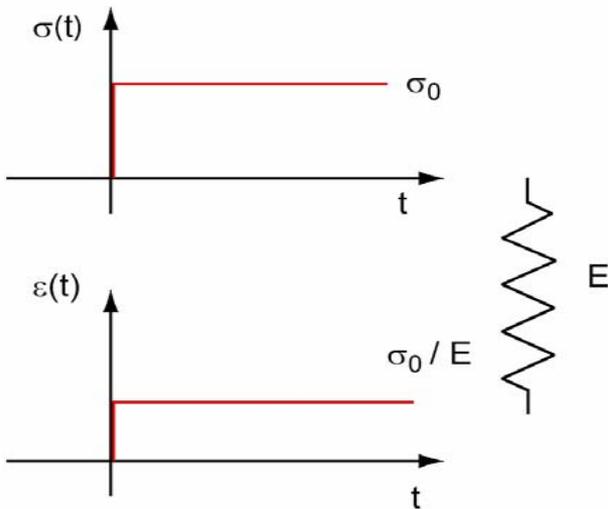
- Elements of polymer structure
 - Linear vs. branched;
 - Vinyl polymers and substitutions
 - Packing of polymer chains
 - Random/amorphous
 - Glass transition temperature, T_G
 - Semi-crystalline
 - Crystalline volume fraction; melting temperature
 - Amorphous T_G
- Elements of linear viscoelasticity
 - Creep and relaxation
 - Analogue models

Idealized Linear Elastic Response

Linear elasticity: $\sigma(t) = E \epsilon(t)$

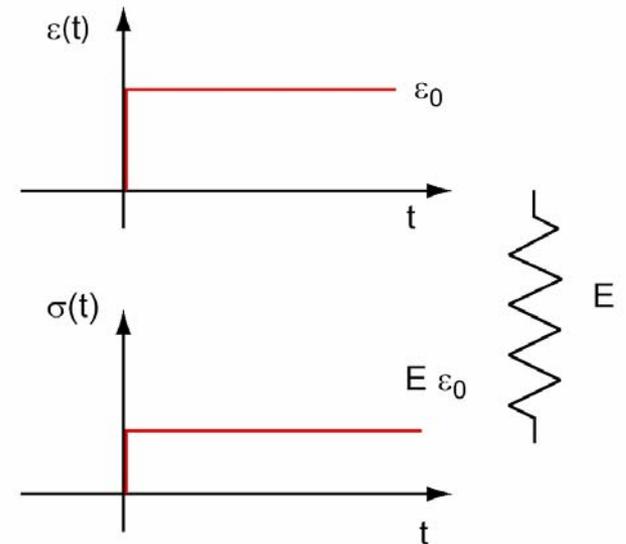
[NO] Creep

INPUT:



[NO] Relaxation

INPUT:



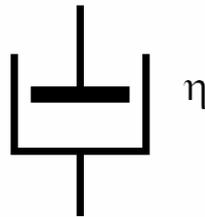
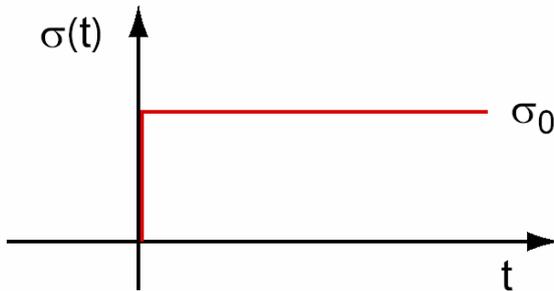
Note: units of E are stress.

Idealized Linear Viscous Response

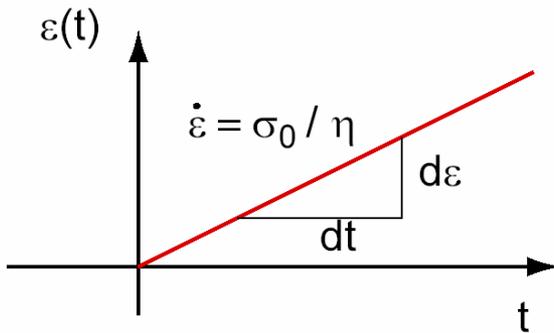
Linear viscosity: $\sigma(t) = \eta \frac{d\epsilon(t)}{dt} \equiv \eta \dot{\epsilon}(t)$

Creep:

INPUT:



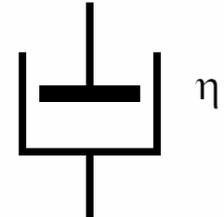
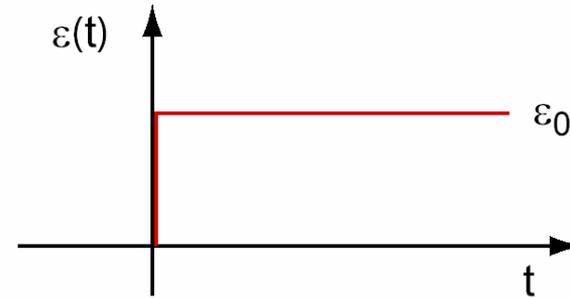
OUTPUT:



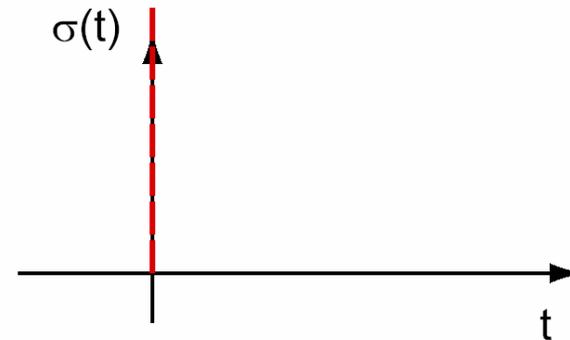
Note: units of η are stress \times time
e.g., sec \times MPa

Relaxation:

INPUT:



OUTPUT:

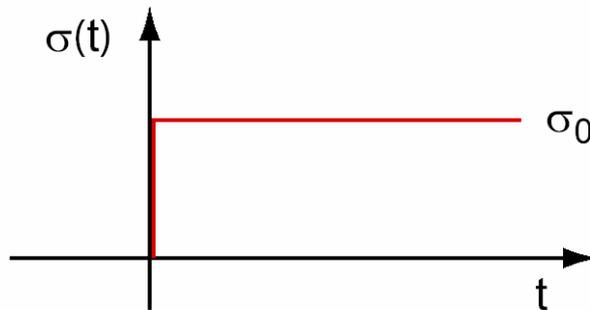


The strain rate (and therefore stress) becomes arbitrarily large during an infinitesimal time interval, and then, like the strain-rate, goes to zero

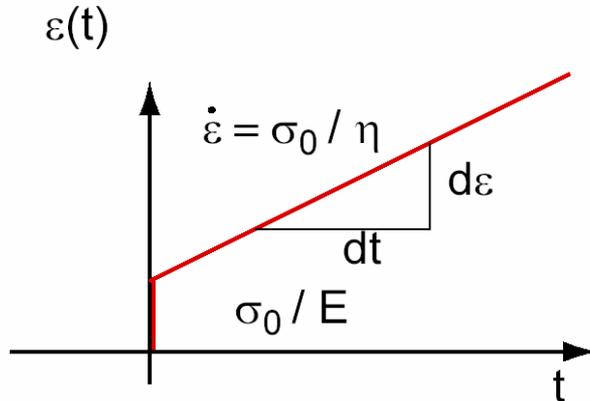
Maxwell Model: an Idealized Linear Viscoelastic Response

Creep:

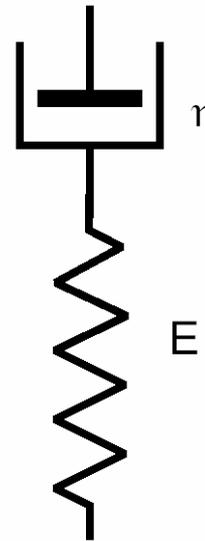
INPUT:



OUTPUT:



$$\begin{aligned}\epsilon(t) &= \epsilon_{\text{elastic}}(t) + \epsilon_{\text{viscous}}(t); \\ \dot{\epsilon}(t) &= \dot{\epsilon}_{\text{elastic}}(t) + \dot{\epsilon}_{\text{viscous}}(t) \\ &= \dot{\sigma}(t)/E + \sigma(t)/\eta\end{aligned}$$



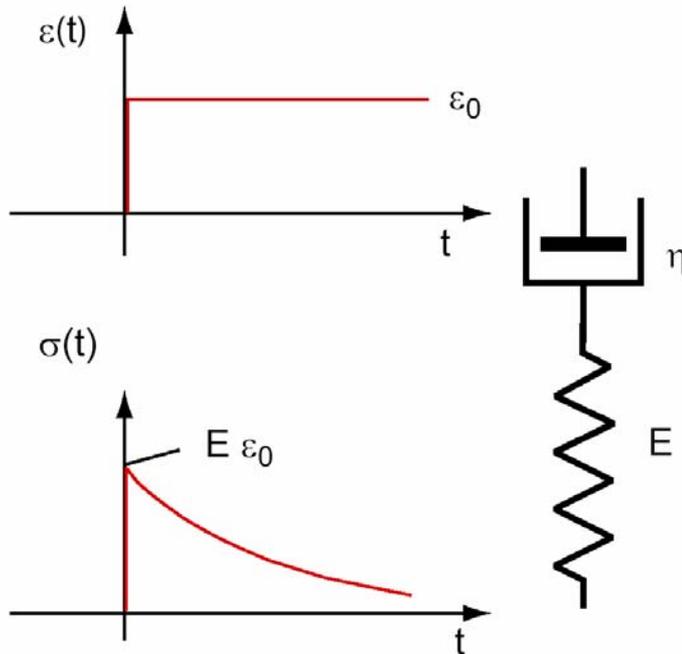
- For times near the finite stress jump, all strain occurs in the elastic element;
- During the hold period, all strain occurs in the viscous element

Maxwell Model: an Idealized Linear Viscoelastic Response

Relaxation:

INPUT:

$$\begin{aligned}\epsilon(t) &= \epsilon_{\text{elastic}}(t) + \epsilon_{\text{viscous}}(t); \\ \dot{\epsilon}(t) &= \dot{\epsilon}_{\text{elastic}}(t) + \dot{\epsilon}_{\text{viscous}}(t) \\ &= \dot{\sigma}(t)/E + \sigma(t)/\eta\end{aligned}$$



OUTPUT:

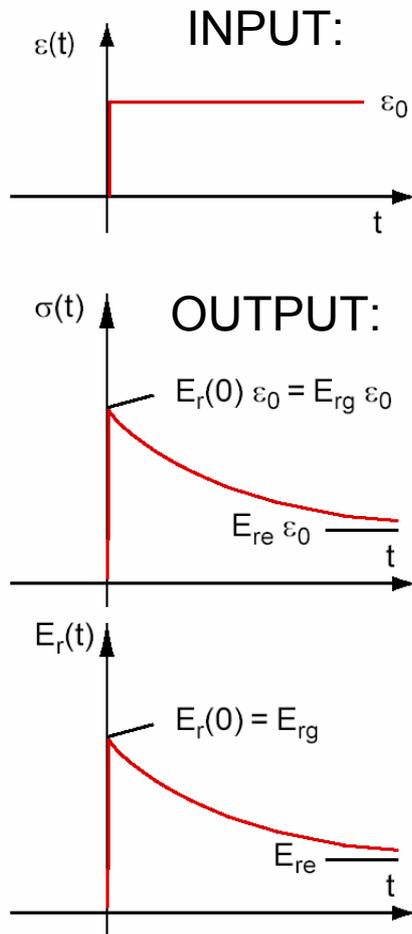
- For times > 0 , $d \epsilon(t)/dt=0$
- During the hold period, elastic strain is traded for viscous strain, and stress drops:

$$\begin{aligned}0 &= \frac{d\sigma(t)}{dt} + \frac{E}{\eta} \sigma(t) \Rightarrow \\ \sigma(t) &= \sigma(t=0) \exp -t/(\eta/E) \\ &= \epsilon_0 \times E \exp -t/(\eta/E)\end{aligned}$$

Characteristic relaxation time:

$$\tau = \eta / E$$

“Real” Polymer Relaxation (an Idealization)



Testing of real polymers under relaxation can be used to extract a time-dependent relaxation modulus, $E_r(t)$;
 Short-term response: E_{rg} ;
 Long-term response: E_{re}

$$E_r(t) \equiv \sigma(t)/\varepsilon_0;$$

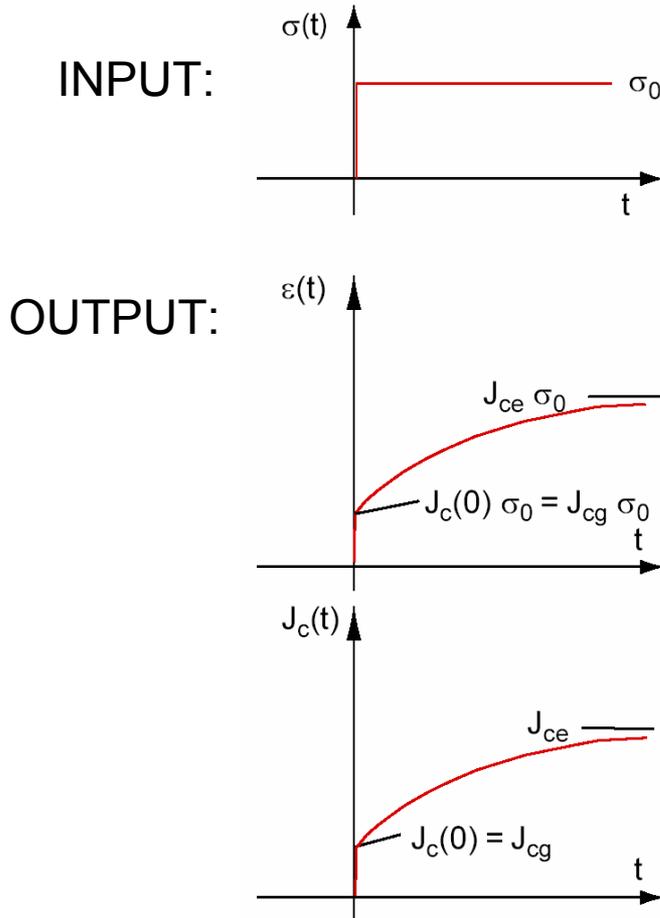
$$E_r(t \rightarrow 0^+) \equiv E_{rg} \quad (\text{glassy})$$

$$E_r(t \rightarrow \infty) \equiv E_{re} \quad (\text{equilibrium})$$

NOTE: provided $|\varepsilon_0|$ is sufficiently small (typically, less than 0.01), the relaxation modulus, $E_r(t)$, is approximately independent of ε_0 .

“Real” Polymer Creep: (an Idealization)

Testing of real polymers under suddenly-applied constant stress can be used to extract a time-dependent creep function, $J_c(t)$;
 Short-term response: J_{cg} ;
 Long-term response: J_{ce}



$$J_c(t) \equiv \epsilon(t) / \sigma_0;$$

$$J_c(t \rightarrow 0^+) \equiv J_{cg} \quad (\text{glassy})$$

$$J_c(t \rightarrow \infty) \equiv J_{ce} \quad (\text{equilibrium})$$

Note: units of $J_c(t)$: 1/ stress

NOTE: provided $|\epsilon(t)|$ remains sufficiently small (typically, less than 0.01), the creep function, $J_c(t)$, is approximately independent of σ_0 .

Relaxation Modulus, $E_r(t)$ and Creep Function, $J_c(t)$: Inverse Functions of Time?

Function	Dimensions	Trend
$J_c(t)$	1/stress	starts small; grows with time
$E_r(t)$	stress	starts large; decays with time

QUESTION: Are these inverse functions? Is $J_c(t) \times E_r(t) \equiv 1$ for all times?

•**ANSWERS:** In general, they are not precise inverses. However,

- Equilibrium and glassy values are nearly inverse:
- For intermediate times, t , the error in assuming that they are inverse is typically only a few per cent at most...

$$J_c(t \rightarrow 0^+) \doteq 1/E_r(t \rightarrow 0^+);$$

$$J_c(t \rightarrow \infty) \doteq 1/E_r(t \rightarrow \infty), \text{ but}$$

$$J_c(t) \times E_r(t) \neq 1.$$

Linearity of Response (an Idealization)

If creep response to stress jump σ_0 is

$$\epsilon(t) = J_c(t) \times \sigma_0,$$

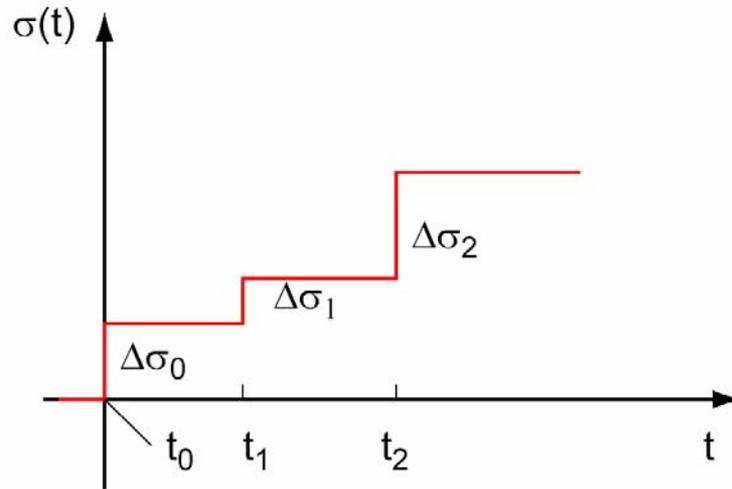
then the creep response to stress jump $\phi \times \sigma_0$
(where ϕ is a proportionality constant) is

$$\epsilon(t) = J_c(t) \times (\phi\sigma_0).$$

NOTE: similar linear scaling of stress relaxation response applies.

Superposition of Loading

Suppose that the stress history input consists of a sequence of stress jumps, $\Delta\sigma_i$, applied at successive times t_i , with $t_0=0$:

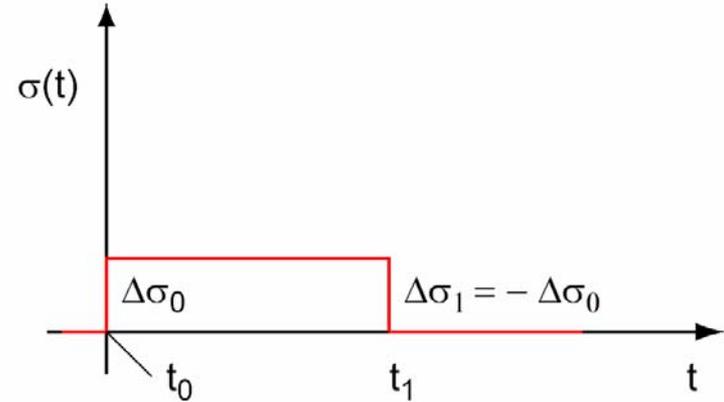


THEN, the resulting strain history is given by

$$\begin{aligned}\epsilon(t) &= \Delta\sigma_0 J_c(t) + \Delta\sigma_1 J_c(t - t_1) + \Delta\sigma_2 J_c(t - t_2) + \dots \\ &= \sum_{j=0} \Delta\sigma_j J_c(t - t_j)\end{aligned}$$

Special Case: Load/Unload

INPUT:

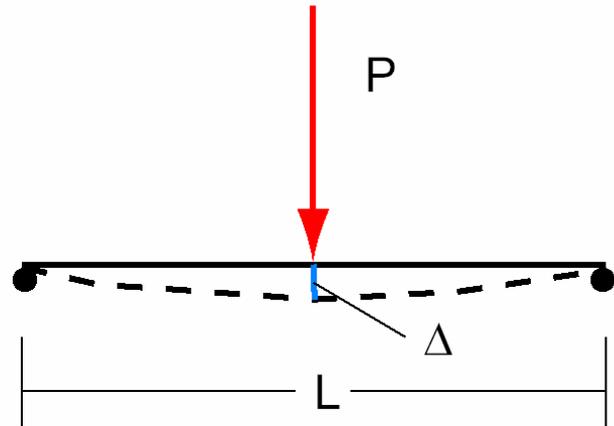


Correspondence Principle

- Suppose that a given load, P , produces displacement vector $\mathbf{u}(\mathbf{x})$ in a linear elastic body having Young's modulus E .
- The displacement vector depends on the position vector $\mathbf{x} = x \mathbf{e}_x + y \mathbf{e}_y + z \mathbf{e}_z$.
- The magnitudes of the displacement and strain components are proportional to P and inversely proportional to E .

EXAMPLE: Three-point
Mid-span bending:

$$\Delta = -v(x = L/2) = \frac{PL^3}{48IE}$$



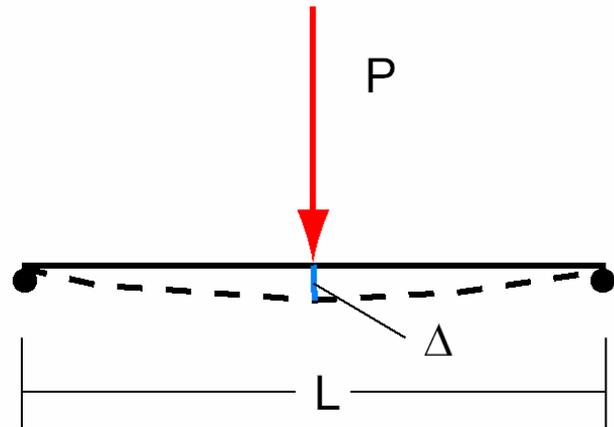
Correspondence Principle

- Now suppose that a given load jump, $P(t)$, is applied to a geometrically identical linear viscoelastic body having creep function $J_c(t)$.
- All stress components in the body are time-independent, and spatially vary precisely as they do in an identical linear elastic body subject to the same load.
- The loading produces time-dependent displacement vector $\mathbf{u}(\mathbf{x}, t)$ and corresponding strain components.
- The magnitudes of the displacement and strain components are proportional to both P and $J_c(t)$.

EXAMPLE: Three-point
Mid-span bending:

$$\Delta(t) = -v(x = L/2, t) = \frac{PL^3}{48I} \times J_c(t)$$

For suddenly-applied load,
replace “ $1/E$ ” with “ $J_c(t)$ ” in
an elastic solution.



Correspondence Principle

- Suppose that a given displacement jump, $\Delta(t)$, is applied to a geometrically identical linear viscoelastic body having stress relaxation modulus $E_r(t)$.
- All displacement and strain components everywhere in the body are time-independent, and precisely equal those in an identical linear elastic body subject to the same applied displacement.
- These boundary conditions produce time-dependent loads, $P(t)$, and stresses
- The magnitudes of the time-dependent load and of the stress components are proportional to both Δ and to $E_r(t)$.

EXAMPLE: Three-point
Mid-span bending:

$$P(t) = \frac{48I}{L^3} \times \Delta \times E_r(t)$$

For suddenly-applied displacement, replace “E” with “ $E_r(t)$ ” in an elastic solution.

