

Euler-Bernoulli Beams: Bending, Buckling, and Vibration

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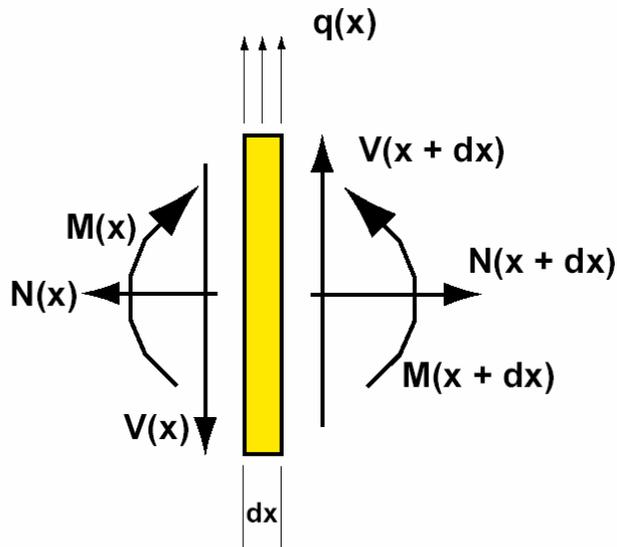
2.002 Mechanics and Materials II
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Linear Elastic Beam Theory

- Basics of beams
 - Geometry of deformation
 - Equilibrium of “slices”
 - Constitutive equations
- Applications:
 - Cantilever beam deflection
 - Buckling of beams under axial compression
 - Vibration of beams

Beam Theory: Slice Equilibrium Relations

- $q(x)$: distributed load/length
- $N(x)$: axial force
- $V(x)$: shear force
- $M(x)$: bending moment



Axial force balance:

$$0 = N(x + dx) - N(x) \Rightarrow N(x) = \text{constant}$$

Transverse force balance:

$$\begin{aligned} 0 &= q(x)dx + V(x + dx) - V(x) \\ &= q(x)dx + \left(V(x) + V'(x)dx + o(dx) \right) - V(x) \\ &= dx \left[V'(x) + q(x) \right] \Rightarrow \\ 0 &= V'(x) + q(x) \quad \text{CDL(3.11)} \end{aligned}$$

Moment balance about 'x+dx':

$$\begin{aligned} 0 &= V(x)dx + M(x + dx) - M(x) - (q(x)dx) dx \\ &= V(x)dx + \left(M(x) + M'(x)dx \right) - M(x) - (q(x)dx) dx/2 \\ &= dx \left[M'(x) + V(x) - q(x)dx/2 \right] \Rightarrow \\ 0 &= M'(x) + V(x) \quad \text{CDL(3.12)} \end{aligned}$$

Euler-Bernoulli Beam Theory: Displacement, strain, and stress distributions

Beam theory assumptions on spatial variation of displacement components:

$$u(x, y, z) = u_0(x) - yv'(x)$$

$$v(x, y, z) = v(x)$$

$$w(x, y, z) = 0$$

Axial strain distribution in beam:

$$\begin{aligned} \epsilon_{xx}(x, y, z) &\equiv \frac{\partial u(x, y, z)}{\partial x} \\ &= u'_0(x) - yv''(x) \\ &\equiv \epsilon_0(x) - y\kappa(x) \end{aligned}$$

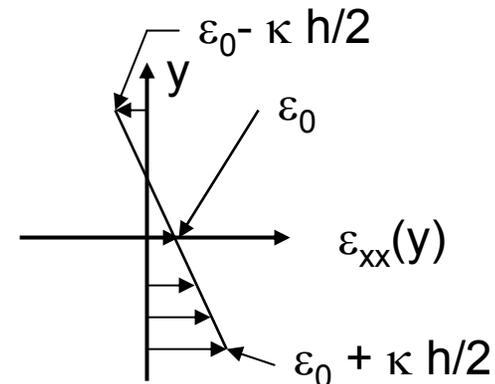
1-D stress/strain relation:

$$\sigma_{xx} = E\epsilon_{xx}$$

Stress distribution in terms of Displacement field:

$$\sigma_{xx}(x, y, z) = E (\epsilon_0(x) - y\kappa(x))$$

Axial **strain** varies linearly
Through-thickness at section 'x'



Slice Equilibrium: Section Axial Force $N(x)$ and Bending Moment $M(x)$ in terms of Displacement fields

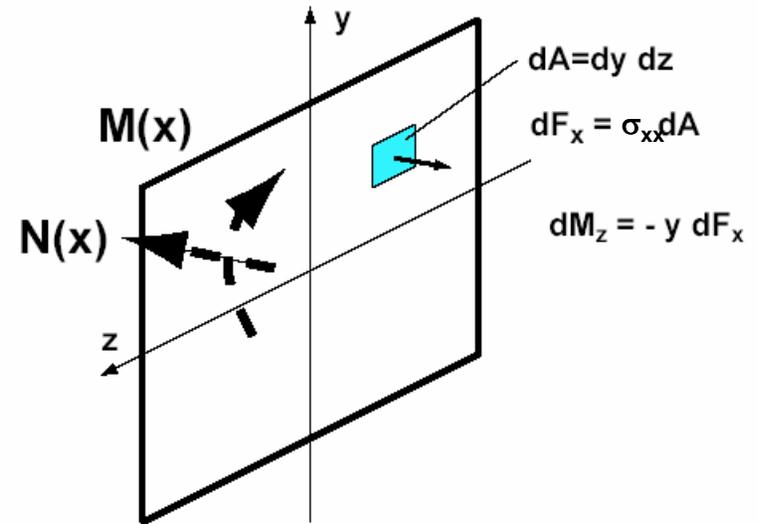
$N(x)$: x-component of force equilibrium
on slice at location 'x':

$$\begin{aligned} N(x) &\equiv \int \sigma_{xx}(x, y, z) dA(y, z) \\ &= \int E \{ \epsilon_0(x) - y\kappa(x) \} dA \\ &= EA\epsilon_0(x) - E\kappa(x) \int y dA. \end{aligned}$$

$M(x)$: z-component of moment equilibrium
on slice at location 'x':

$$\begin{aligned} M(x) &\equiv \int -y \sigma_{xx}(x, y, z) dA(y, z) \\ &= \int E \{ -y\epsilon_0(x) + y^2\kappa(x) \} dA \\ &= -E\epsilon_0(x) \int y dA + E\kappa(x) I \end{aligned}$$

where $I \equiv \int y^2 dA$ is area moment of inertia of cross section



Centroidal Coordinates

$$\bar{y} = \frac{1}{A} \int y dA$$

choice: $\bar{y} \equiv 0 \Rightarrow \int y dA = 0$

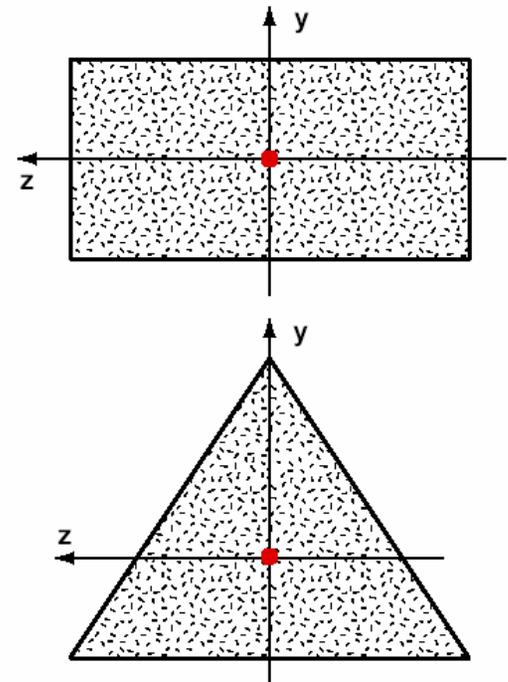
Simplifications:

$$N(x) = EA\epsilon_0(x) = EAu'_0(x)$$

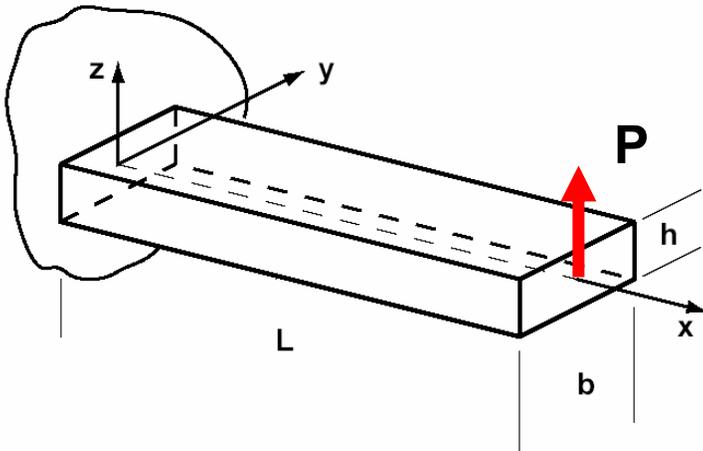
$$M(x) = EI\kappa(x) = EIv''(x)$$

Note: I is centroidal area moment of inertia:

$$I \equiv \int y^2 dA$$

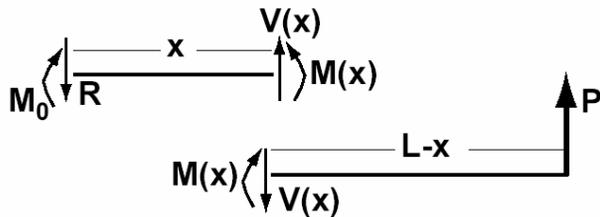


Tip-Loaded Cantilever Beam: Equilibrium



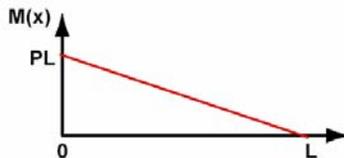
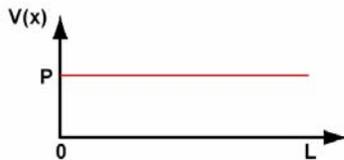
- statically determinant:
support reactions R , M_0
from equilibrium alone
- reactions “present”
because of $x=0$ geometrical
boundary conditions $v(0)=0$;
 $v'(0)=\phi(0)=0$

Free body diagrams:



$$0 = P - V(x)$$

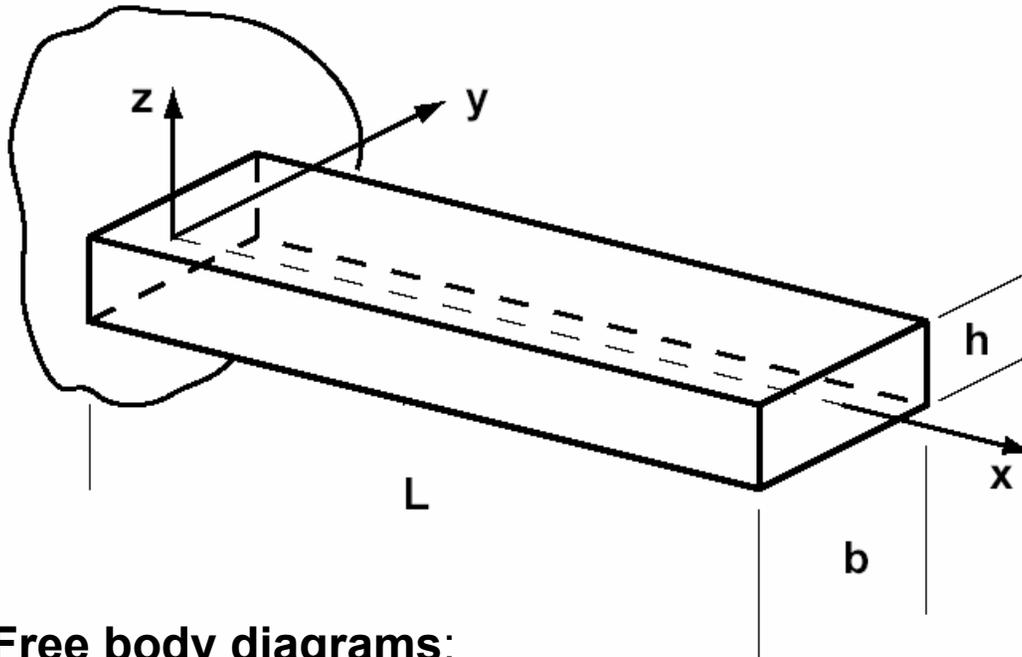
$$0 = M(x) - P(L - x)$$



- general equilibrium
equations (CDL 3.11-12)
satisfied

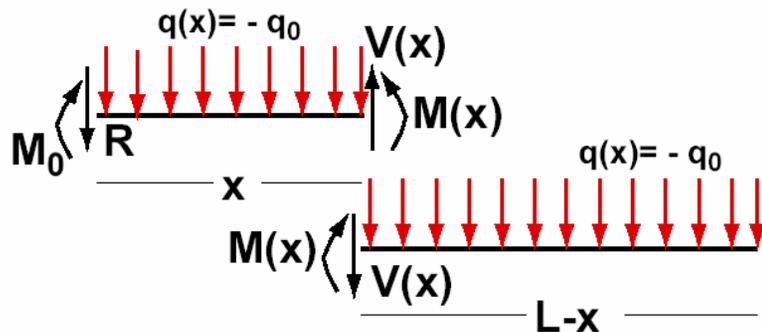
How to determine lateral displacement $v(x)$; especially at tip ($x=L$)?

Exercise: Cantilever Beam Under Self-Weight



- Weight per unit length: q_0
- $q_0 = \rho Ag = \rho b h g$

Free body diagrams:



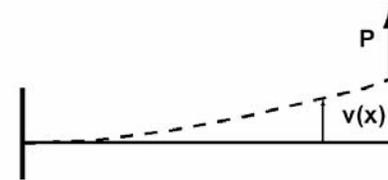
Find:

- Reactions: R and M_0
- Shear force: $V(x)$
- Bending moment: $M(x)$

Tip-Loaded Cantilever: Lateral Deflections

curvature / moment relations:

$$\begin{aligned}v''(x) &= \frac{1}{EI} M(x) \\ &= \frac{1}{EI} (P(L - x)) \Rightarrow \\ v'(x) &= \frac{P}{EI} \left(Lx - \frac{x^2}{2} + C_1 \right) \Rightarrow \\ v(x) &= \frac{P}{EI} \left(Lx^2/2 - x^3/6 + C_1x + C_2 \right)\end{aligned}$$



geometric boundary conditions

$$\begin{aligned}\phi(0) = v'(0) &= 0 \Rightarrow C_1 = 0 \\ v(0) &= 0 \Rightarrow C_2 = 0 \\ v(x) &= \frac{Px^2}{6EI} (3L - x)\end{aligned}$$

tip deflection and rotation:

$$\begin{aligned}\Delta \equiv v(L) &= \frac{PL^3}{3EI} \\ \Phi \equiv v'(L) &= \frac{PL^2}{2EI}\end{aligned}$$

stiffness and modulus:

$$\begin{aligned}k \equiv \frac{P}{\Delta} &= \frac{3EI}{L^3} \\ E &= \frac{kL^3}{3I} = \frac{PL^3}{3I\Delta}\end{aligned}$$

Tip-Loaded Cantilever: Axial Strain Distribution

strain field (no axial force):

$$\begin{aligned}\epsilon_{xx}(x, y) &= -yv''(x) \\ &= -\frac{yM(x)}{EI}\end{aligned}$$

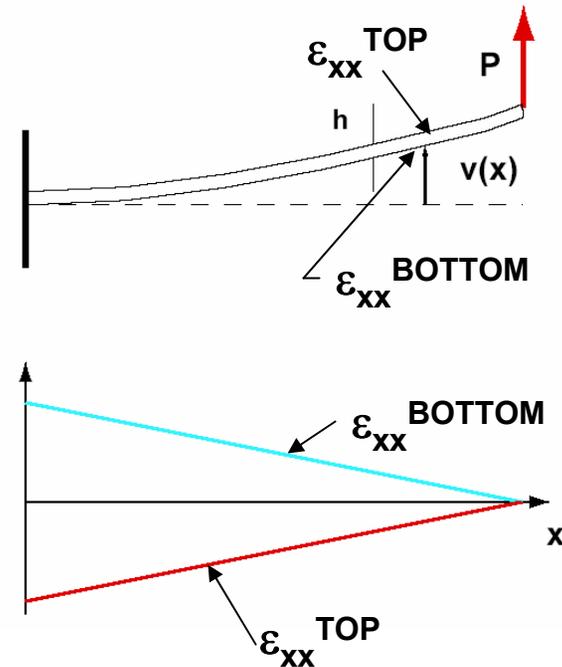
top/bottom axial strain distribution:

$$\begin{aligned}\epsilon_{xx}^{TOP}(x) &= -\frac{6P(L-x)}{bh^2E} \quad (y = h/2) \\ \epsilon_{xx}^{BOTTOM}(x) &= \frac{6P(L-x)}{bh^2E} \quad (y = -h/2)\end{aligned}$$

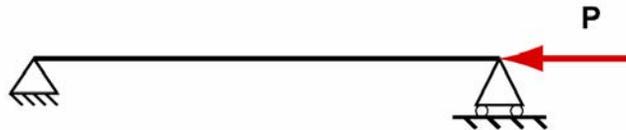
$$I_{\text{rectangle}} = \frac{bh^3}{12}$$

strain-gauged estimate of E:

$$E = \frac{6P(L-x)}{bh^2\epsilon_{xx}^{BOTTOM}(x)} = \frac{6P(L-x)}{bh^2|\epsilon_{xx}^{TOP}(x)|}$$

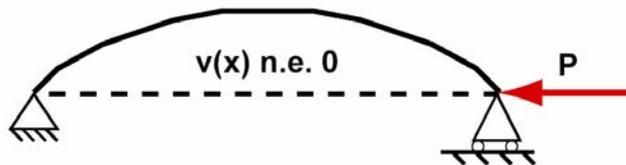


Euler Column Buckling: Non-uniqueness of deformed configuration



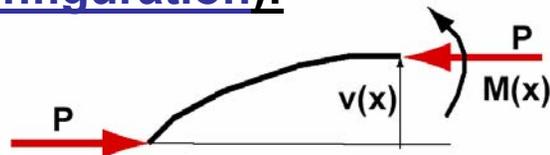
One solution:
 $N(x) = -P = \text{constant};$
 $v(x) = 0; u_0(x) = -Px/AE \quad \circ$

When might a buckled shape exist?



free body diagram

(note: evaluated in deformed configuration):



$$\sum M_z = 0 \Rightarrow M(x) + Pv(x) = 0$$

moment/curvature:

$$M(x) = EI\kappa(x) = EIv''(x)$$

ode for buckled shape:

$$\begin{aligned} 0 &= M(x) + Pv(x) \\ &= EIv''(x) + Pv(x) \\ 0 &= v''(x) + \frac{P}{EI}v(x) \\ &\equiv v''(x) + k^2v(x) \end{aligned}$$

Note: linear 2nd order ode;
Constant coefficients (but
parametric: $k^2 = P/EI$

Euler Column Buckling, Cont.

ode for buckled shape:

$$0 = v''(x) + k^2 v(x)$$

general solution to ode:

$$v(x) = C_1 \sin kx + C_2 \cos kx$$

boundary conditions:

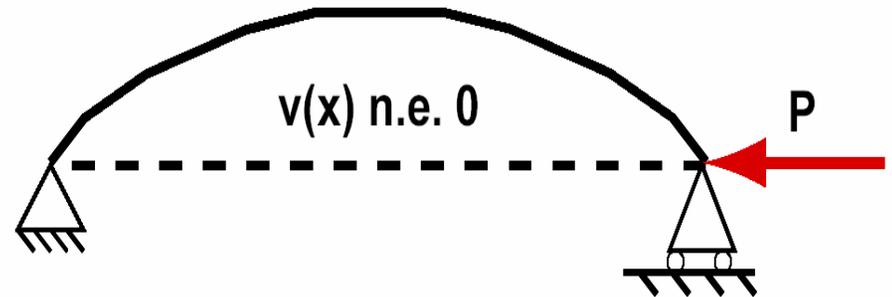
$$v(0) = 0 \Rightarrow C_2 = 0$$

$$v(L) = 0 \Rightarrow C_1 \sin kL = 0 \Rightarrow$$

$$C_1 = 0 \text{ (trivial)} \quad \text{or} \quad \sin kL = 0$$

buckling-based estimate of E:

$$E_{\text{pinned/pinned}} = \frac{P_{\text{crit}} L^2}{\pi^2 I}$$



parametric consequences:

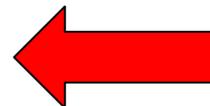
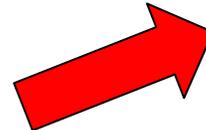
non-trivial buckled shape only when

$$\sin kL = 0 \Rightarrow kL = n\pi$$

$$k^2 = (n\pi/L)^2$$

$$P = EI k^2 = \frac{n^2 EI \pi^2}{L^2}$$

$$P_{\text{crit}}(n=1) = \frac{\pi^2 EI}{L^2}$$



Euler Column Buckling: General Observations

- buckling load, P_{crit} , is proportional to EI/L^2
- proportionality constant depends strongly on boundary conditions at both ends:
 - the more kinematically restrained the ends are, the larger the constant and the higher the critical buckling load (see Lab 1 handout)
- safe design of long slender columns requires adequate margins with respect to buckling
- buckling load may occur at a compressive stress value ($\sigma=P/A$) that is less than yield stress, σ_y

Euler-Bernoulli Beam Vibration

assume time-dependent lateral motion:

$$v(x, t) = \bar{v}(x) \sin \omega t$$

lateral velocity of slice at 'x':

$$\frac{\partial v(x, t)}{\partial t} \equiv \dot{v}(x, t) = \omega \bar{v}(x) \cos \omega t$$

lateral acceleration of slice at 'x':

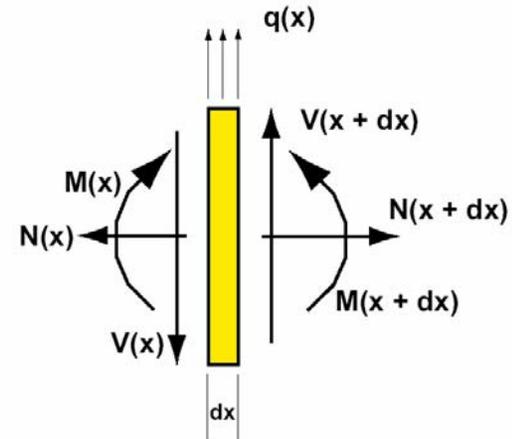
$$\frac{\partial^2 v(x, t)}{\partial t^2} \equiv \ddot{v}(x, t) = -\omega^2 \bar{v}(x) \sin \omega t$$

mass of dx-thickness slice:

$$dm = dx \rho A$$

linear momentum balance (Newton):

$$\begin{aligned} \sum F_y &= dm \ddot{v}(x, t) \Rightarrow \\ 0 &= dx \left(M''(x, t) - \rho \omega^2 A \bar{v}(x) \sin \omega t \right) \end{aligned}$$



net lateral force (q(x,t)=0):

$$\sum F_y = dx \frac{\partial V(x, t)}{\partial x} \equiv dx V'(x, t)$$

moment balance:

$$\begin{aligned} 0 &= \frac{\partial M(x, t)}{\partial x} + V(x, t) \\ &\equiv M'(x, t) + V(x, t) \Rightarrow \\ 0 &= M''(x, t) + V'(x, t) \end{aligned}$$

Euler-Bernoulli Beam Vibration, Cont.(1)

linear momentum balance:

$$0 = M''(x, t) - \rho\omega^2 A\bar{v}(x) \sin \omega t$$

moment/curvature:

$$M(x, t) = EI \kappa(x, t) = EI \bar{v}''(x) \sin \omega t$$

ode for mode shape, $v(x)$, and vibration frequency, ω :

$$\begin{aligned} 0 &= M''(x, t) - \rho\omega^2 A\bar{v}(x) \sin \omega t \\ &= \sin \omega t \left(EI \bar{v}''''(x) - \rho\omega^2 A\bar{v}(x) \right) \\ &= \sin \omega t \left(\bar{v}''''(x) - \frac{\rho\omega^2 A}{EI} \bar{v}(x) \right) \\ &\equiv \sin \omega t \left(\bar{v}''''(x) - \beta^4 \bar{v}(x) \right) \Rightarrow \\ 0 &= \bar{v}''''(x) - \beta^4 \bar{v}(x) \end{aligned}$$

general solution to ode:

$$\bar{v}(x) = A_1 \sin \beta x + A_2 \cos \beta x + A_3 \sinh \beta x + A_4 \cosh \beta x$$

Euler-Bernoulli Beam Vibration, Cont(2)

general solution to ode:

$$\bar{v}(x) = A_1 \sin \beta x + A_2 \cos \beta x + A_3 \sinh \beta x + A_4 \cosh \beta x$$



pinned/pinned boundary conditions:

$$\bar{v}(0) = 0 \Rightarrow A_2 + A_4 = 0$$

$$\bar{v}''(0) = 0 \Rightarrow \beta^2 (-A_2 + A_4) = 0$$

$$\bar{v}(L) = 0 \Rightarrow A_1 \sin \beta L + A_3 \sinh \beta L = 0$$

$$\bar{v}''(L) = 0 \Rightarrow \beta^2 (-A_1 \sin \beta L + A_3 \sinh \beta L) = 0$$



pinned/pinned restricted solution:

$$\beta \neq 0; \quad A_2 = A_3 = A_4 = 0;$$

$$A_1 \sin \beta L = 0 \Rightarrow$$

$$A_1 = 0 \text{ (trivial), \quad OR}$$

$$\sin \beta L = 0 \Rightarrow \beta L = n\pi$$

Solution (n=1, first mode):

A_1 : 'arbitrary' (but small)
vibration amplitude

$$\beta_1 = \pi/L \Rightarrow$$

$$v_{(n=1)}(x, t) = A_1 \sin(\pi x/L) \sin \omega_1 t$$

$$\beta_1^4 = (\pi/L)^4 = \omega_1^2 \rho A / EI \Rightarrow$$

$$\omega_1 = \sqrt{\frac{EI \pi^4}{\rho A L^4}}$$

τ_1 : period of first mode:

$$\tau_1 = \frac{2\pi}{\omega_1}$$

