

## Orthogonality Relations

We now explain the basic reason why the remarkable Fourier coefficient formulas work. We begin by repeating them from the last note:

$$\begin{aligned} a_0 &= \frac{1}{L} \int_{-L}^L f(t) dt, \\ a_n &= \frac{1}{L} \int_{-L}^L f(t) \cos\left(n\frac{\pi}{L}t\right) dt, \\ b_n &= \frac{1}{L} \int_{-L}^L f(t) \sin\left(n\frac{\pi}{L}t\right) dt. \end{aligned} \tag{1}$$

The key fact is the following collection of integral formulas for sines and cosines, which go by the name of **orthogonality relations**:

$$\frac{1}{L} \int_{-L}^L \cos\left(n\frac{\pi}{L}t\right) \cos\left(m\frac{\pi}{L}t\right) dt = \begin{cases} 1 & n = m \neq 0 \\ 0 & n \neq m \\ 2 & n = m = 0 \end{cases}$$

$$\frac{1}{L} \int_{-L}^L \cos\left(n\frac{\pi}{L}t\right) \sin\left(m\frac{\pi}{L}t\right) dt = 0$$

$$\frac{1}{L} \int_{-L}^L \sin\left(n\frac{\pi}{L}t\right) \sin\left(m\frac{\pi}{L}t\right) dt = \begin{cases} 1 & n = m \neq 0 \\ 0 & n \neq m \end{cases}$$

**Proof of the orthogonality relations:** This is just a straightforward calculation using the periodicity of sine and cosine and either (or both) of these two methods:

Method 1: use  $\cos at = \frac{e^{iat} + e^{-iat}}{2}$ , and  $\sin at = \frac{e^{iat} - e^{-iat}}{2i}$ .

Method 2: use the trig identity  $\cos(\alpha)\cos(\beta) = \frac{1}{2}(\cos(\alpha + \beta) + \cos(\alpha - \beta))$ , and the similar trig identities for  $\cos(\alpha)\sin(\beta)$  and  $\sin(\alpha)\sin(\beta)$ .

**Using the orthogonality relations to prove the Fourier coefficient formula**

Suppose we know that a periodic function  $f(t)$  has a Fourier series expansion

$$f(t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos\left(n\frac{\pi}{L}t\right) + b_n \sin\left(n\frac{\pi}{L}t\right) \tag{2}$$

How can we find the values of the coefficients? Let's choose one coefficient, say  $a_2$ , and compute it; you will easily how to generalize this to any other coefficient. The claim is that the right-hand side of the Fourier coefficient formula (1), namely the integral

$$\frac{1}{L} \int_{-L}^L f(t) \cos\left(2\frac{\pi}{L}t\right) dt.$$

is in fact the coefficient  $a_2$  in the series (2). We can replace  $f(t)$  in this integral by the series in (2) and multiply through by  $\cos\left(2\frac{\pi}{L}t\right)$ , to get

$$\frac{1}{L} \int_{-L}^L \frac{a_0}{2} \cos\left(2\frac{\pi}{L}t\right) + \sum_{n=1}^{\infty} \left( a_n \cos\left(n\frac{\pi}{L}t\right) \cos\left(2\frac{\pi}{L}t\right) + b_n \sin\left(n\frac{\pi}{L}t\right) \cos\left(2\frac{\pi}{L}t\right) \right) dt$$

Now the orthogonality relations tell us that almost every term in this sum will integrate to 0. In fact, the only non-zero term is the  $n = 2$  cosine term

$$\frac{1}{L} \int_{-L}^L a_2 \cos\left(2\frac{\pi}{L}t\right) \cos\left(2\frac{\pi}{L}t\right) dt$$

and the orthogonality relations for the case  $n = m = 2$  show this integral is equal to  $a_2$  as claimed.

**Why the denominator of 2 in  $\frac{a_0}{2}$  ?**

Answer: it is in fact just a convention, but the one which allows us to have the same Fourier coefficient formula for  $a_n$  when  $n = 0$  and  $n \geq 1$ . (Notice that in the  $n = m$  case for cosine, there is a factor of 2 only for  $n = m = 0$ .)

**Interpretation of the constant term  $\frac{a_0}{2}$ .**

We can also interpret the constant term  $\frac{a_0}{2}$  in the Fourier series of  $f(t)$  as the average of the function  $f(t)$  over one full period:  $\frac{a_0}{2} = \frac{1}{2L} \int_{-L}^L f(t) dt$ .

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