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18.02 Multivariable Calculus  
Fall 2007

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## 18.02 Lecture 30. – Tue, Nov 27, 2007

Handouts: practice exams 4A and 4B.

Clarification from end of last lecture: we derived the diffusion equation from 2 inputs.  $u =$  concentration,  $\mathbf{F} =$  flow, satisfy:

- 1) from physics:  $\mathbf{F} = -k\nabla u$ ,
- 2) from divergence theorem:  $\partial u/\partial t = -\text{div } \mathbf{F}$ .

Combining, we get the diffusion equation:  $\partial u/\partial t = -\text{div } \mathbf{F} = +k\text{div}(\nabla u) = k\nabla^2 u$ .

### Line integrals in space.

Force field  $\mathbf{F} = \langle P, Q, R \rangle$ , curve  $C$  in space,  $d\vec{r} = \langle dx, dy, dz \rangle$

$$\Rightarrow \text{Work} = \int_C \mathbf{F} \cdot d\vec{r} = \int_C P dx + Q dy + R dz.$$

Example:  $\mathbf{F} = \langle yz, xz, xy \rangle$ .  $C$ :  $x = t^3$ ,  $y = t^2$ ,  $z = t$ .  $0 \leq t \leq 1$ . Then  $dx = 3t^2 dt$ ,  $dy = 2t dt$ ,  $dz = dt$  and substitute:

$$\int_C \mathbf{F} \cdot d\vec{r} = \int_C yz dx + xz dy + xy dz = \int_0^1 6t^5 dt = 1$$

(In general, express  $(x, y, z)$  in terms of a *single* parameter: 1 degree of freedom)

Same  $\mathbf{F}$ , curve  $C' =$  segments from  $(0, 0, 0)$  to  $(1, 0, 0)$  to  $(1, 1, 0)$  to  $(1, 1, 1)$ . In the  $xy$ -plane,  $z = 0 \implies \mathbf{F} = xy\hat{\mathbf{k}}$ , so  $\mathbf{F} \cdot d\vec{r} = 0$ , no work. For the last segment,  $x = y = 1$ ,  $dx = dy = 0$ , so  $\mathbf{F} = \langle z, z, 1 \rangle$  and  $d\vec{r} = \langle 0, 0, dz \rangle$ . We get  $\int_0^1 1 dz = 1$ .

Both give the same answer because  $\mathbf{F}$  is conservative, in fact  $\mathbf{F} = \nabla(xyz)$ .

Recall the fundamental theorem of calculus for line integrals:

$$\int_{P_0}^{P_1} \nabla f \cdot d\vec{r} = f(P_1) - f(P_0).$$

### Gradient fields.

$$\mathbf{F} = \langle P, Q, R \rangle = \langle f_x, f_y, f_z \rangle ?$$

Then  $f_{xy} = f_{yx}$ ,  $f_{xz} = f_{zx}$ ,  $f_{yz} = f_{zy}$ , so  $P_y = Q_x$ ,  $P_z = R_x$ ,  $Q_z = R_y$ .

Theorem:  $\mathbf{F}$  is a gradient field if and only if these equalities hold (assuming defined in whole space or simply connected region)

Example: for which  $a, b$  is  $axy\hat{\mathbf{i}} + (x^2 + z^3)\hat{\mathbf{j}} + (byz^2 - 4z^3)\hat{\mathbf{k}}$  a gradient field?

$$P_y = ax = 2x = Q_x \text{ so } a = 2; P_z = 0 = 0 = R_x; Q_z = 3z^2 = bz^2 = R_y \text{ so } b = 3.$$

Systematic method to find a potential: (carried out on above example)

$$f_x = 2xy, f_y = x^2 + z^3, f_z = 3yz^2 - 4z^3:$$

$$f_x = 2xy \Rightarrow f(x, y, z) = x^2y + g(y, z).$$

$$f_y = x^2 + g_y = x^2 + z^3 \Rightarrow g_y = z^3 \Rightarrow g(y, z) = yz^3 + h(z), \text{ and } f = x^2y + yz^3 + h(z).$$

$$f_z = 3yz^2 + h'(z) = 3yz^2 - 4z^3 \Rightarrow h'(z) = -4z^3 \Rightarrow h(z) = -z^4 + c, \text{ and } f = x^2y + yz^3 - z^4 + c.$$

Other method:  $f(x_1, y_1, z_1) = f(0, 0, 0) + \int_{P_0}^{P_1} \mathbf{F} \cdot d\vec{r}$  (use a curve that gives an easy computation, e.g. 3 segments parallel to axes).

**Curl:** encodes by how much  $\mathbf{F}$  fails to be conservative.

$$\text{curl} \langle P, Q, R \rangle = (R_y - Q_z)\hat{\mathbf{i}} + (P_z - R_x)\hat{\mathbf{j}} + (Q_x - P_y)\hat{\mathbf{k}}.$$

How to remember the formula? Use the del operator  $\nabla = \langle \partial/\partial x, \partial/\partial y, \partial/\partial z \rangle$ .

Recall from last week that  $\nabla \cdot \mathbf{F} = \langle \partial/\partial x, \partial/\partial y, \partial/\partial z \rangle \cdot \langle P, Q, R \rangle = P_x + Q_y + R_z = \text{div } \mathbf{F}$ .

$$\text{Now we have: } \nabla \times \mathbf{F} = \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ P & Q & R \end{vmatrix} = \text{curl } \mathbf{F}.$$

**Interpretation of curl for velocity fields:** curl measures *angular velocity*.

Example: rotation around  $z$ -axis at constant angular velocity  $\omega$  (trajectories are circles centered on  $z$ -axis):  $\mathbf{v} = \langle -\omega y, \omega x, 0 \rangle$ .

Then  $\nabla \times \mathbf{v} = \dots = 0\hat{\mathbf{i}} + 0\hat{\mathbf{j}} + (\omega + \omega)\hat{\mathbf{k}} = 2\omega\hat{\mathbf{k}}$ . So length of curl = twice angular velocity, and direction = axis of rotation.

A general motion can be complicated, but decomposes into various effects.

- curl measures the *rotation* component of a complex motion.

## 18.02 Lecture 31. – Thu, Nov 29, 2007

Handouts: PS11 solutions, PS12.

Stokes' theorem is the 3D analogue of Green's theorem for work (in the same sense as the divergence theorem is the 3D analogue of Green for flux).

$$\text{Recall } \text{curl } \mathbf{F} = \nabla \times \mathbf{F} = \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ P & Q & R \end{vmatrix} = \langle R_y - Q_z, P_z - R_x, Q_x - P_y \rangle.$$

**Stokes' theorem:** if  $C$  is a *closed curve*, and  $S$  *any* surface bounded by  $C$ , then

$$\oint_C \mathbf{F} \cdot d\vec{r} = \iint_S (\nabla \times \mathbf{F}) \cdot \hat{\mathbf{n}} dS.$$

Orientation: compatibility of an orientation of  $C$  with an orientation of  $S$  (changing orientation changes sign on both sides of Stokes).

Rule: if I walk along  $C$  in positive direction, with  $S$  to my left, then  $\hat{\mathbf{n}}$  is pointing up. (Various examples shown.)

Another formulation (right-hand rule): if thumb points along  $C$  (1-D object), index finger towards  $S$  (2-D object), then middle finger points along  $\hat{\mathbf{n}}$  (3-D object).

More examples shown.

**Example: Stokes vs. Green.** If  $S$  is a portion of  $xy$ -plane bounded by a curve  $C$  counterclockwise, then  $\oint_C \mathbf{F} \cdot d\vec{r} = \int_C P dx + Q dy$ , by Green this is equal to  $\iint_S (Q_x - P_y) dx dy = \iint_S \nabla \times \mathbf{F} \cdot \hat{\mathbf{k}} dx dy = \iint_S \nabla \times \mathbf{F} \cdot \hat{\mathbf{n}} dS$ , so Green and Stokes say the same thing in this example.

**Remark.** In Stokes' theorem we are free to choose any surface  $S$  bounded by  $C$ ! (e.g. if  $C$  = circle,  $S$  could be a disk, a hemisphere, a cone, ...)

**“Proof” of Stokes.**

- 1) if  $C$  and  $S$  are in the  $xy$ -plane then the statement follows from Green.
- 2) if  $C$  and  $S$  are in an arbitrary plane: this also reduces to Green in the given plane. Green/Stokes works in any plane because of *geometric invariance* of work, curl and flux under rotations of space. They can be defined in purely geometric terms so as not to depend on the coordinate system  $(x, y, z)$ ; equivalently, we can choose coordinates  $(u, v, w)$  adapted to the given plane, and work

with those coordinates, the expressions of work, curl, flux will be the familiar ones replacing  $x, y, z$  with  $u, v, w$ .

3) in general, we can decompose  $S$  into small pieces, each piece is nearly flat (slanted plane); on each piece we have approximately work = flux by Green's theorem. When adding pieces, the line integrals over the inner boundaries cancel each other and we get the line integral over  $C$ ; the flux integrals add up to flux through  $S$ .

**Example:** verify Stokes for  $\mathbf{F} = z\hat{i} + x\hat{j} + y\hat{k}$ ,  $C =$  unit circle in  $xy$ -plane (counterclockwise),  $S =$  piece of paraboloid  $z = 1 - x^2 - y^2$ .

Direct calculation:  $x = \cos\theta$ ,  $y = \sin\theta$ ,  $z = 0$ , so  $\oint_C \mathbf{F} \cdot d\vec{r} = \int_C z dx + x dy + y dz = \oint_C x dy = \int_0^{2\pi} \cos^2\theta d\theta = \pi$ .

By Stokes:  $\text{curl } \mathbf{F} = \langle 1, 1, 1 \rangle$ , and  $\hat{\mathbf{n}} dS = \langle -f_x, -f_y, 1 \rangle dx dy = \langle 2x, 2y, 1 \rangle dx dy$ .

$$\iint_S (\nabla \times \mathbf{F}) \cdot d\vec{S} = \iint_S \langle (2x + 2y + 1) \rangle dx dy = \iint_S 1 dx dy = \text{area}(\text{disk}) = \pi.$$

( $\iint x dx dy = 0$  by symmetry and similarly for  $y$ ).

## 18.02 Lecture 32. – Fri, Nov 30, 2007

### Stokes and path independence.

Definition: a region is simply connected if every closed loop  $C$  inside it bounds some surface  $S$  inside it.

Example: the complement of the  $z$ -axis is not simply connected (shown by considering a loop encircling the  $z$ -axis); the complement of the origin is simply connected.

Topology: uses these considerations to classify for example surfaces in space: e.g., the mathematical proof that a sphere and a torus are “different” surfaces is that the sphere is simply connected, the torus isn't (in fact it has two “independent” loops that don't bound).

Recall: if  $\mathbf{F}$  is a gradient field then  $\text{curl}(\mathbf{F}) = 0$ .

Conversely, Theorem: if  $\nabla \times \mathbf{F} = 0$  in a *simply connected* region then  $\mathbf{F}$  is conservative (so  $\int \mathbf{F} \cdot d\vec{r}$  is path-independent and we can find a potential).

Proof: Assume  $R$  simply connected,  $\nabla \times \mathbf{F} = 0$ , and consider two curves  $C_1$  and  $C_2$  with same end points. Then  $C = C_1 - C_2$  is a closed curve so bounds some  $S$ ;  $\int_{C_1} \mathbf{F} \cdot d\vec{r} - \int_{C_2} \mathbf{F} \cdot d\vec{r} = \oint_C \mathbf{F} \cdot d\vec{r} = \iint_S \nabla \times \mathbf{F} \cdot \hat{\mathbf{n}} dS = 0$ .

**Orientability.** We can apply Stokes' theorem to any surface  $S$  bounded by  $C$ ... provided that it has a well-defined normal vector! Counterexample shown: the Möbius strip. It's a one-sided surface, so we can't compute flux through it (no possible consistent choice of orientation of  $\hat{\mathbf{n}}$ ). Instead, if we want to apply Stokes to the boundary curve  $C$ , we must find a two-sided surface with boundary  $C$ . (pictures shown).

### Stokes and surface independence.

In Stokes we can choose any  $S$  bounded by  $C$ : so if a same  $C$  bounds two surfaces  $S_1, S_2$ , then  $\oint_C \mathbf{F} \cdot d\vec{r} = \iint_{S_1} \text{curl } \mathbf{F} \cdot \hat{\mathbf{n}} dS = \iint_{S_2} \text{curl } \mathbf{F} \cdot \hat{\mathbf{n}} dS$ ? Can we prove directly that the two flux integrals are equal?

Answer: change orientation of  $S_2$ , then  $S = S_1 - S_2$  is a closed surface with  $\hat{\mathbf{n}}$  outwards; so we can apply the divergence theorem:  $\iint_S \text{curl } \mathbf{F} \cdot \hat{\mathbf{n}} dS = \iiint_D \text{div}(\text{curl } \mathbf{F}) dV$ . But  $\text{div}(\text{curl } \mathbf{F}) = 0$ ,

always. (Checked by calculating in terms of components of  $\mathbf{F}$ ; also, symbolically:  $\nabla \cdot (\nabla \times \mathbf{F}) = 0$ , much like  $u \cdot (u \times v) = 0$  for genuine vectors).

#### Review for Exam 4.

We've seen three types of integrals, with different ways of evaluating:

1)  $\iiint f dV$  in rect., cyl., spherical coordinates (I re-explained the general setup and the formulas for  $dV$ ); applications: center of mass, moment of inertia, gravitational attraction.

2) surface integrals, flux. Setting up  $\iint_S \mathbf{F} \cdot \hat{\mathbf{n}} dS$ , by knowing formulas for  $\hat{\mathbf{n}} dS$ .

We have seen: planes parallel to coordinate planes (e.g.  $yz$ -plane:  $\hat{\mathbf{n}} = \pm \hat{\mathbf{i}}$ ,  $dS = dy dz$ ); spheres and cylinders ( $\hat{\mathbf{n}}$  = straight out/in from center or axis;  $dS = a dz d\theta$  for cylinders,  $a^2 \sin \phi d\phi d\theta$  for spheres); if we can express  $z = f(x, y)$ ,  $\hat{\mathbf{n}} dS = \pm \langle -f_x, -f_y, 1 \rangle dx dy$  (recall  $\langle \dots \rangle$  is not  $\hat{\mathbf{n}}$  and  $dx dy$  is not  $dS$ ); if  $S$  has a given normal vector  $\vec{N}$  (e.g. if  $S$  is given by  $g(x, y, z) = 0$ ),  $\hat{\mathbf{n}} dS = \pm \vec{N} / (\vec{N} \cdot \hat{\mathbf{k}}) dx dy$ .

3) line integrals  $\int_C \mathbf{F} \cdot d\vec{r}$  ( $= \int_C P dx + Q dy + R dz$ ), evaluate by parameterizing  $C$  and expressing in terms of a single variable.

While these various types of integrals are completely different in terms of interpretation and method of evaluation, we've seen some theorems that establish bridges between them:

a) ( $\iint$  vs  $\iiint$ ) divergence theorem:  $S$  closed surface,  $\hat{\mathbf{n}}$  outwards, then  $\iint_S \mathbf{F} \cdot \hat{\mathbf{n}} dS = \iiint_D (\text{div } \mathbf{F}) dV$ .

b) ( $\int$  vs  $\iint$ ) Stokes' theorem:  $C$  closed curve bounding  $S$  compatibly oriented, then  $\int_C \mathbf{F} \cdot d\vec{r} = \iint_S (\nabla \times \mathbf{F}) \cdot \hat{\mathbf{n}} dS$ .

Both sides of these theorems are integrals of the types discussed above, and are evaluated by the usual methods! (even if the integrand happens to be a div or a curl).

In fact, another conceptually similar bridge exists between no integral at all and line integral: the fundamental theorem of calculus,  $f(P_1) - f(P_0) = \int_C \nabla f \cdot d\vec{r}$ .

One more topic: given  $\mathbf{F}$  with  $\text{curl } \mathbf{F} = 0$ , finding a potential function.