

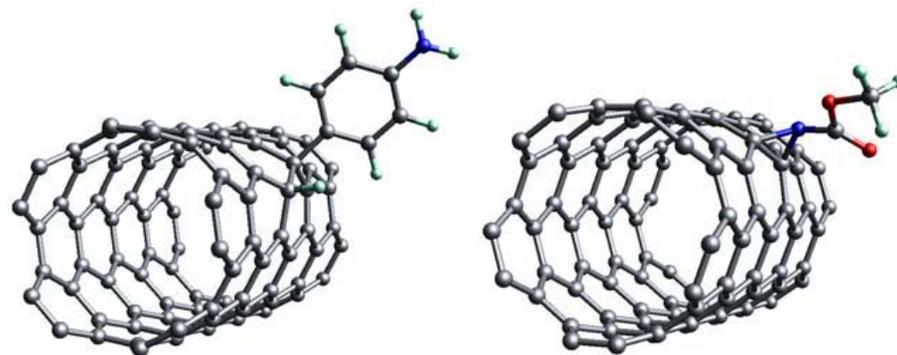
Functionalized Carbon Nanotubes

Organic chemistry meets electronics

engineer multiple functionalities in carbon nanotubes
with ligands attached to the sidewalls

Covalent Functionalization

- Robust attachment of sidewalls ligands
- Several groups have succeeded in covalent linking of hydrogen, fluorine, methyl, dichlorocarbene, nitrene, carboxyl
- Disrupted π -bonding detected by Raman spectroscopy



J.L.Bahr et al., J. Am. Chem. Soc., **123** 6536 (2001)

M.Holzinger et al., J. Am. Chem. Soc., **125** 8566 (2003)

A. Hirsch, Angew. Chem. Int. Ed., 41, 1853 (2002)

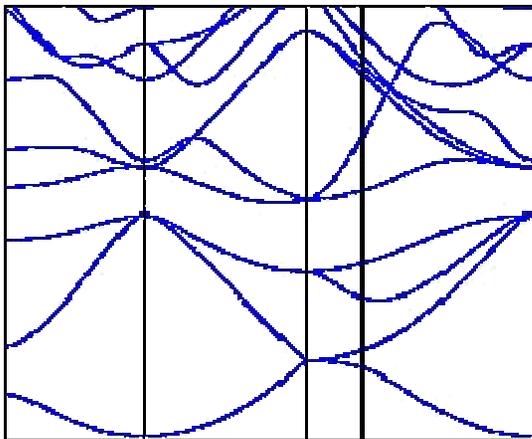
S. Niyogi et al, Acc. Chem. Res., 35, 1105 (2002)

J. L. Bahr et al, J. Mater. Chem., 12, 1952 (2002)

Generalized Wannier Functions for Composite Bands

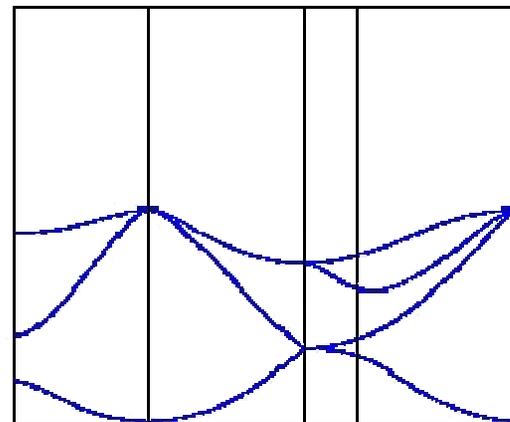
- $\{|\mathbf{R}n\rangle\}$ span the same space as $\{|\Psi_{nk}\rangle\}$
- $|\mathbf{R}n\rangle = w_n(\mathbf{r} - \mathbf{R})$ (translational images)
- $\langle \mathbf{R}n | \mathbf{R}'m \rangle = \delta_{n,m} \delta_{\mathbf{R},\mathbf{R}'}$
- “maximally” localized

Silicon

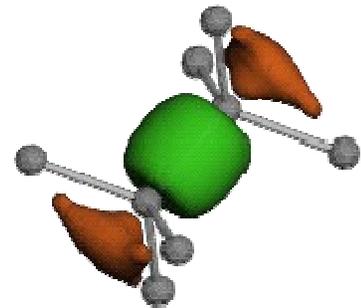


→
(Isolated group of bands)

Valence bands



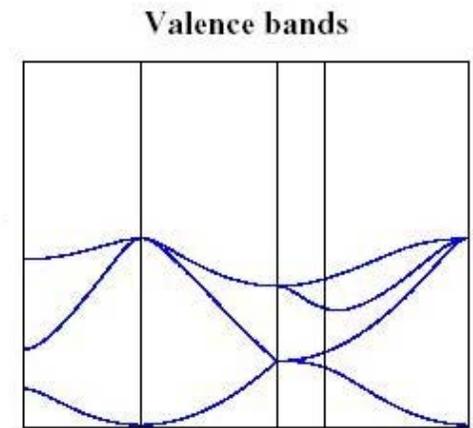
Wannier function



From Bloch Orbitals to Wannier Functions

Periodic $V_{\text{ext}} \Rightarrow \Psi_{n\mathbf{k}}(\mathbf{r}) = u_{n\mathbf{k}}(\mathbf{r}) e^{i\mathbf{k}\cdot\mathbf{r}}$

$$|\mathbf{R}n\rangle = \int_{BZ} \Psi_{n\mathbf{k}}(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{R}} d\mathbf{k}$$



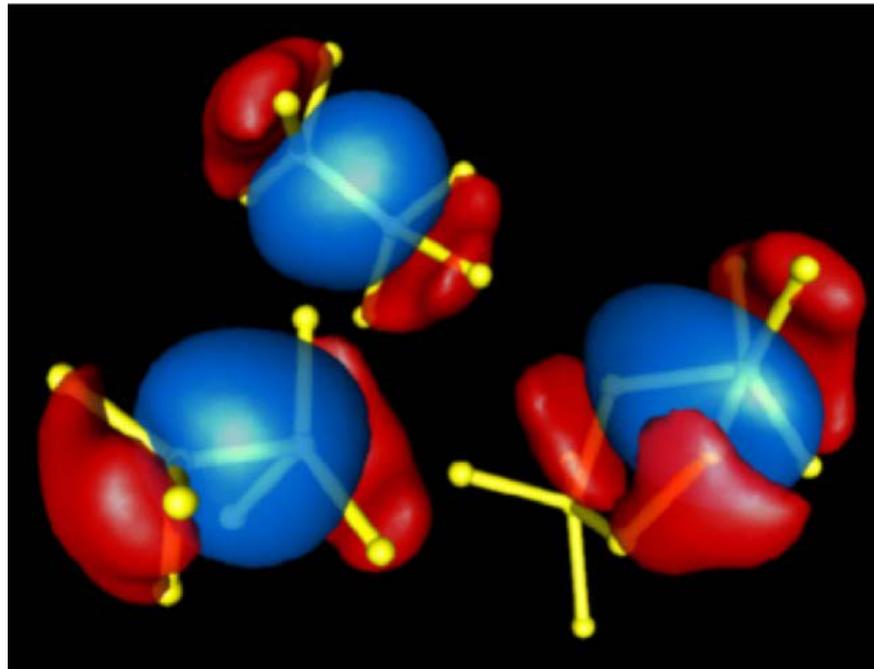
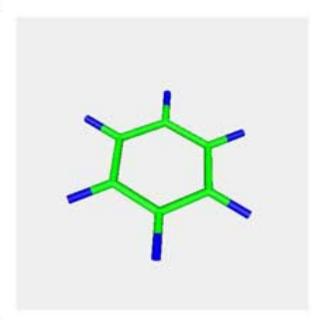
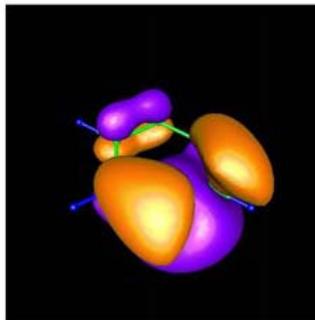
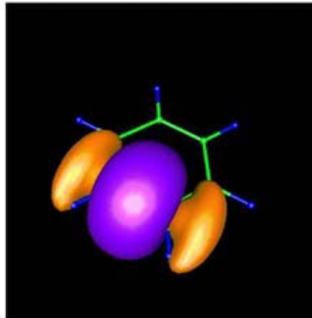
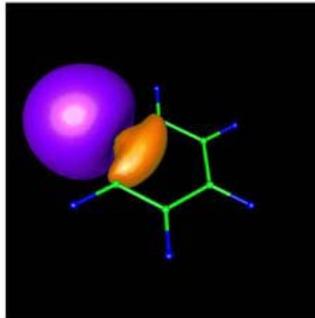
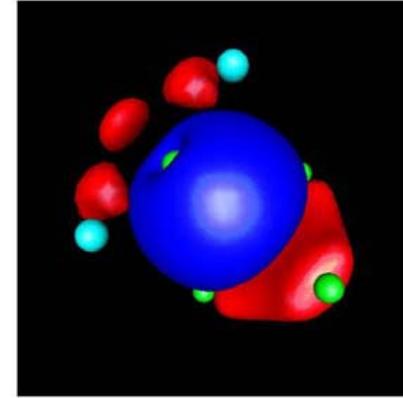
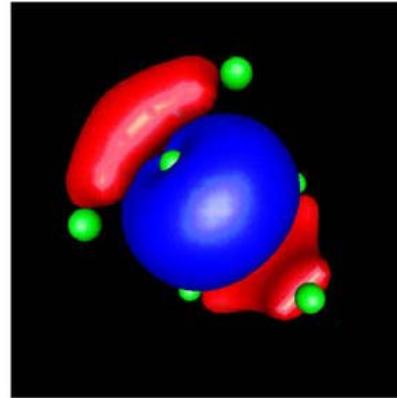
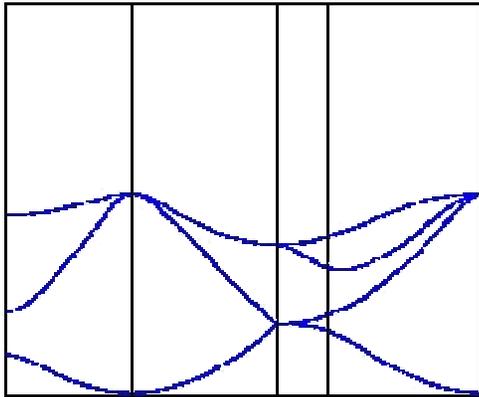
Gauge freedoms

- Arbitrary phase factor for every $n\mathbf{k}$ (Schrödinger)
- Arbitrary unitary rotations $U_{mn}^{(\mathbf{k})}$ for every \mathbf{k} (DFT)

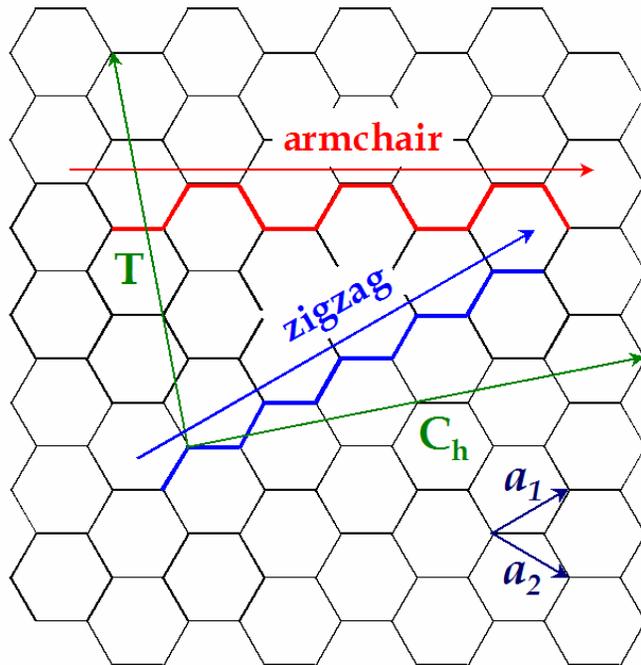
$$|\mathbf{R}n\rangle = \int_{BZ} \sum_m U_{mn}^{(\mathbf{k})} \Psi_{m\mathbf{k}}(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{R}} d\mathbf{k}$$

Silicon, GaAs, Amorphous Silicon, Benzene

Valence bands

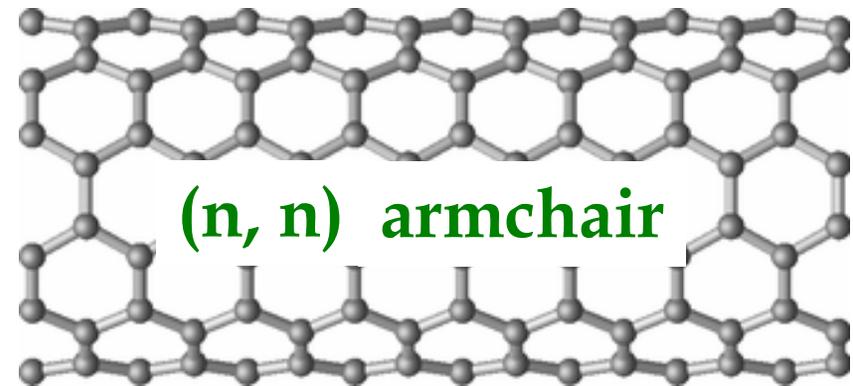


Pristine Carbon Nanotubes



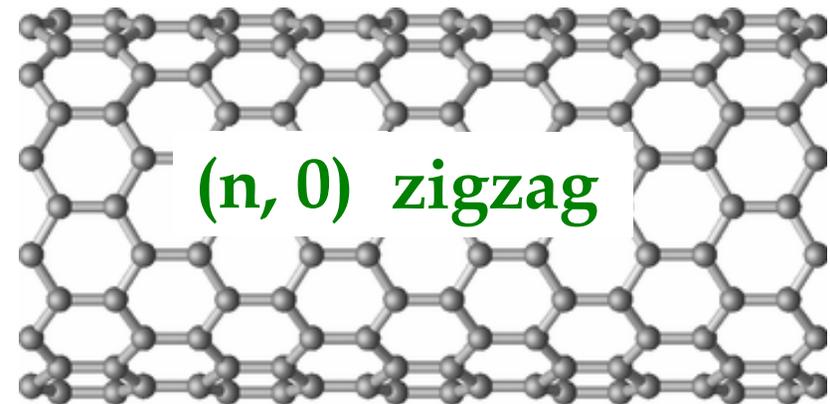
$$C_h = n a_1 + m a_2$$

- Armchair (n,n)
- Zig-Zag (n,0)
- Chiral (n,m) $n \neq m$



(n, n) armchair

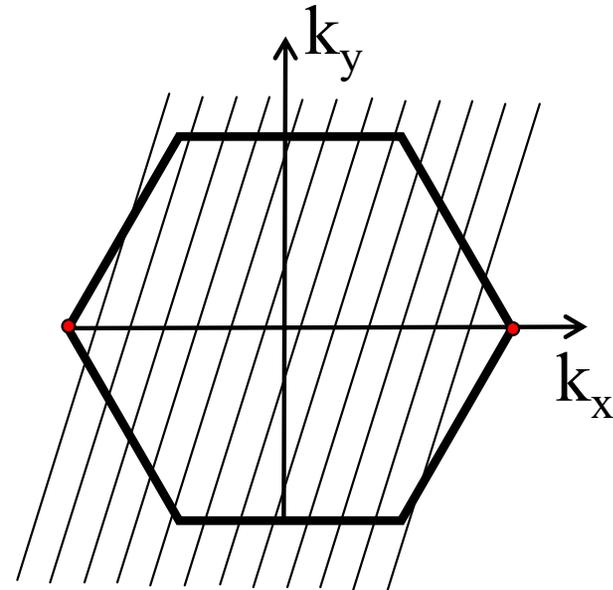
tube axis



(n, 0) zigzag

2-Dimensional Band Structure of Graphene

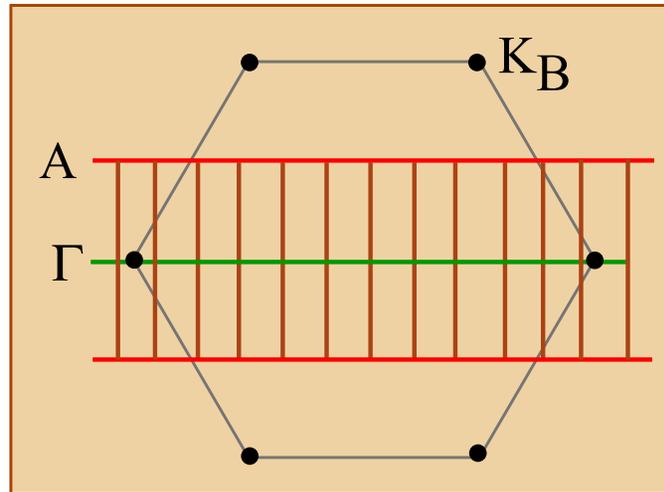
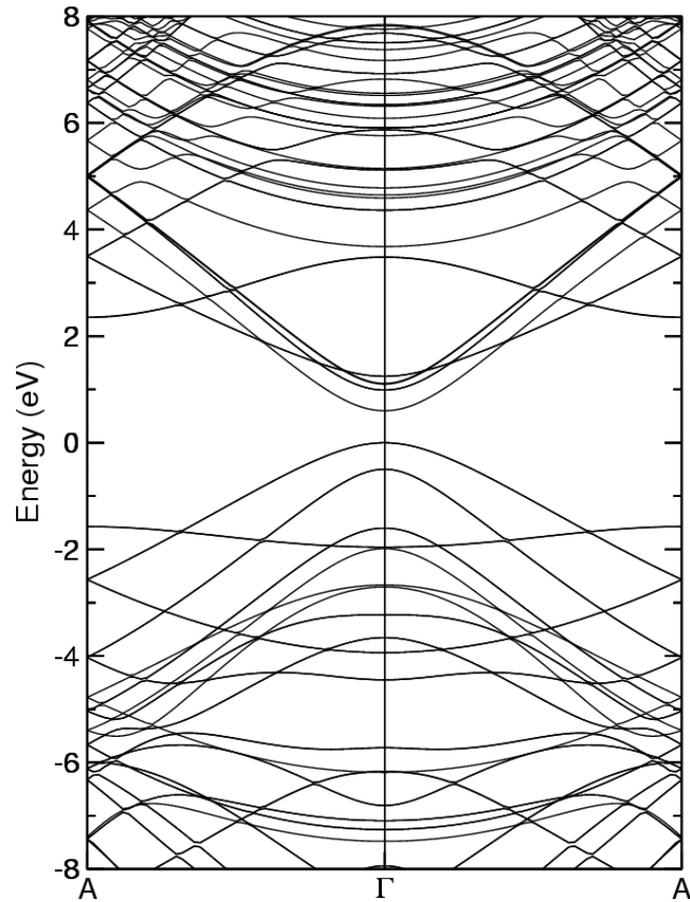
Image removed for copyright reasons.
See Souza, I., N. Marzari and D. Vanderbilt.
"Maximally localized Wannier functions for entangled
energy bands." *Physical Review B* 65 (2001).



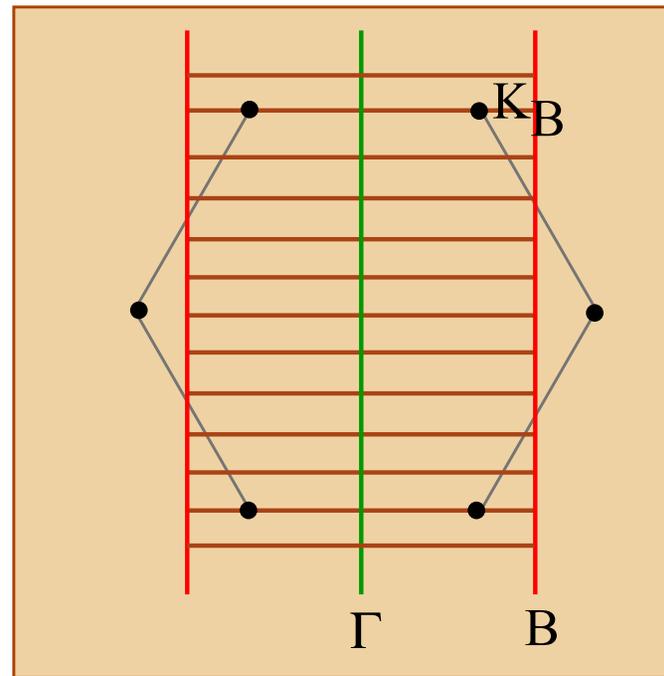
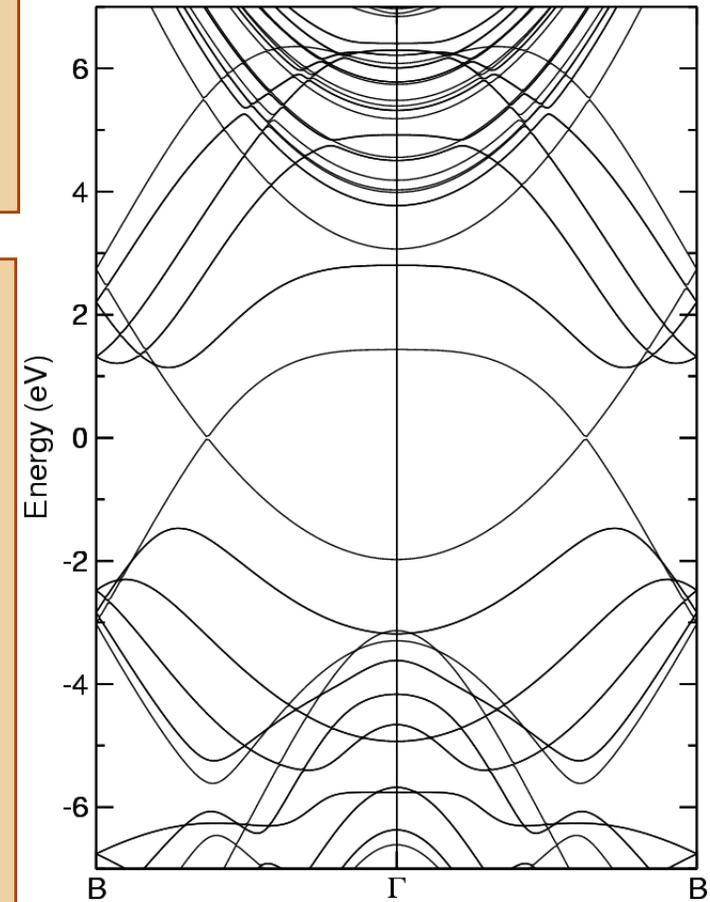
● **mod(3) rule: if $(n-m)$ is a multiple of 3 the tube is metallic**

Band Structure of (8,0) and (5,5) SWNT

(8,0) semiconducting

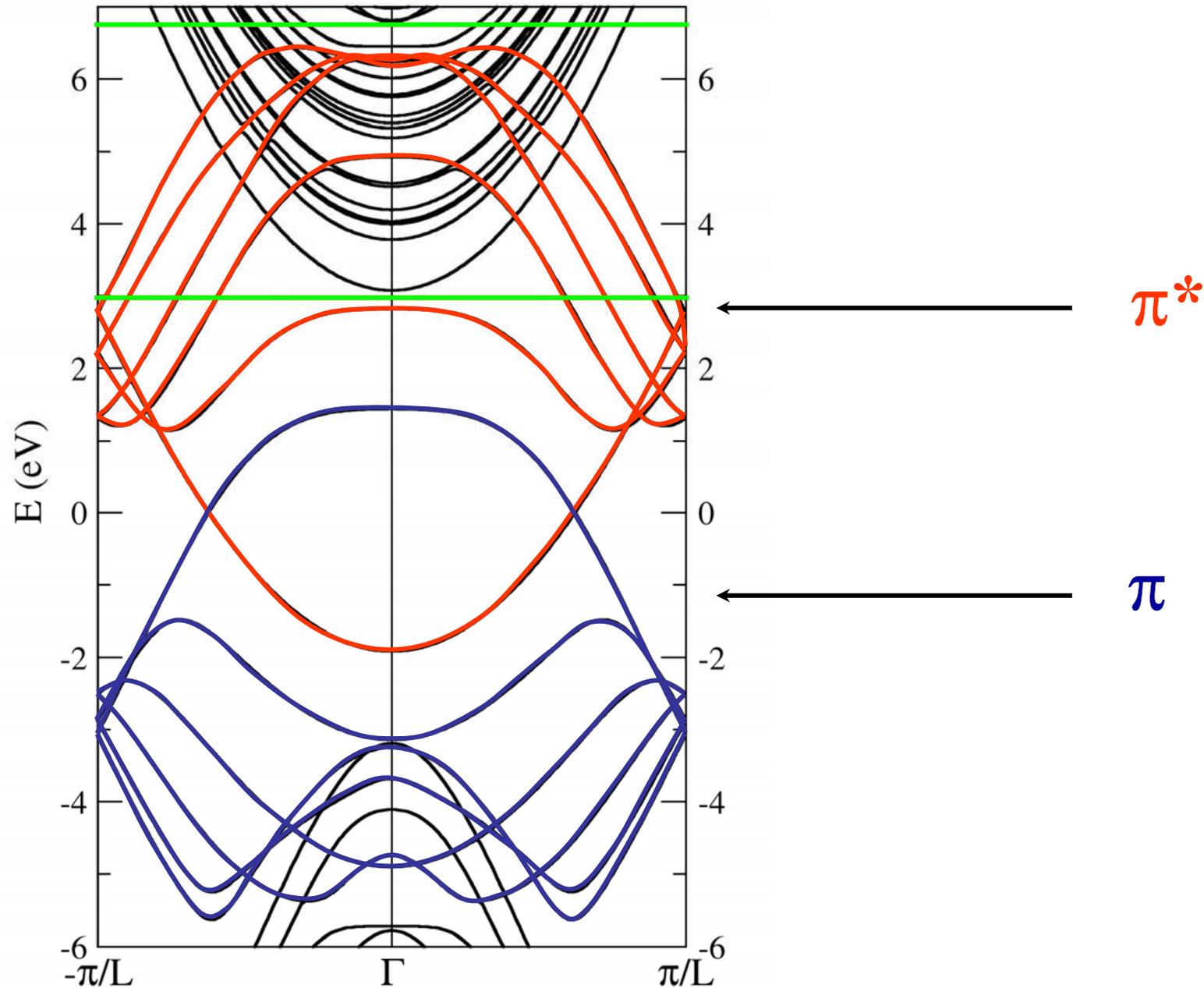


(5,5) metallic



Figures by MIT OCW.

Disentanglement: Conduction Bands in (5,5) SWNT



The LEGO Bricks of Electronic Structure

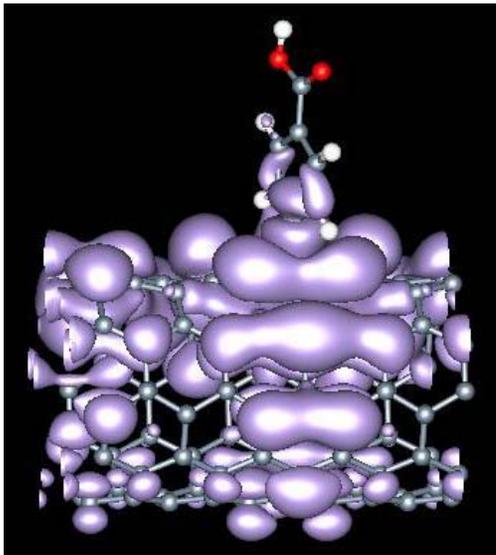
First-Principles
Molecular Dynamics:
Time Evolution of the
Electronic Ground State



Optimal Unitary
Transformation of the
Bloch Orbitals



Real Space
Maximally-Localized
Wannier Functions

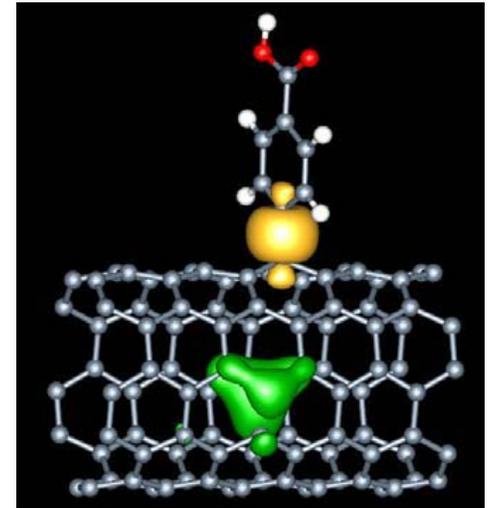


Minimization of the spread
functional

$$\Omega = \sum_n [\langle r^2 \rangle_n - \langle \mathbf{r} \rangle_n^2]$$

exploiting the arbitrariness of the
unitary transformations between the
Bloch orbitals

$$|\mathbf{R}n\rangle = \int_{BZ} \sum_m U_{mn}^{(\mathbf{k})} \Psi_{m\mathbf{k}}(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{R}} d\mathbf{k}$$



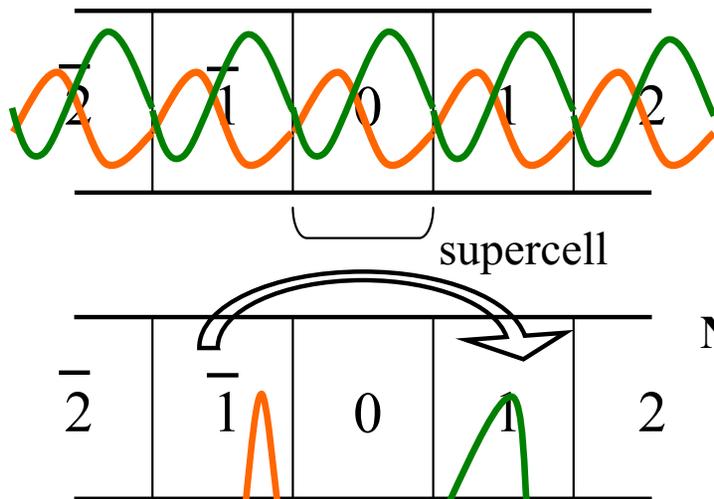
The LEGO Bricks of Electronic Structure

Sparse
Hamiltonian Matrix



Large Scale Calculations

Zero interaction between
Wannier orbitals whose centers
are further apart than decay
length



$$H = \begin{pmatrix} H_{01}^+ & H_{00} & H_{01} & 0 \\ & H_{01}^+ & H_{00} & H_{01} \\ & & H_{01}^+ & H_{00} & H_{01} \\ 0 & & & & \end{pmatrix}$$

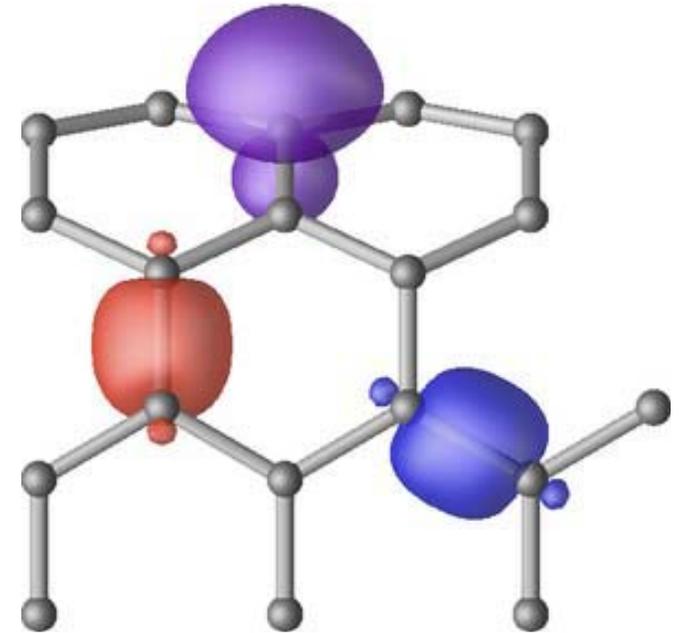
No interaction

Max-loc WFs \leftrightarrow “Exact” Tight-Binding

Compact mapping of Bloch states into local orbitals

$$|\mathbf{R}n\rangle \quad \omega_n(\mathbf{r}-\mathbf{R}) \quad \frac{V}{8\pi^3} \int_{BZ} e^{-i\mathbf{k}\cdot\mathbf{R}} \psi_{n\mathbf{k}}(\mathbf{r}) d\mathbf{k}$$

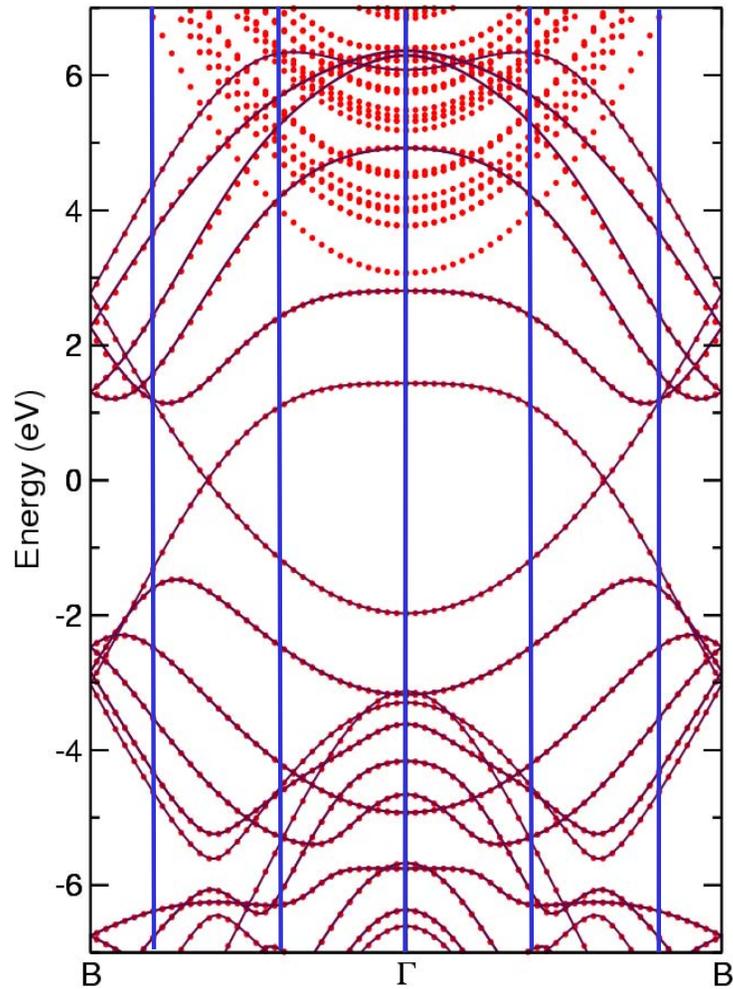
$$\psi_{n\mathbf{k}}(\mathbf{r}) \quad \frac{1}{\sqrt{N_R}} \sum_{\mathbf{R}} e^{i\mathbf{k}\cdot\mathbf{R}} \omega_n(\mathbf{r}-\mathbf{R})$$



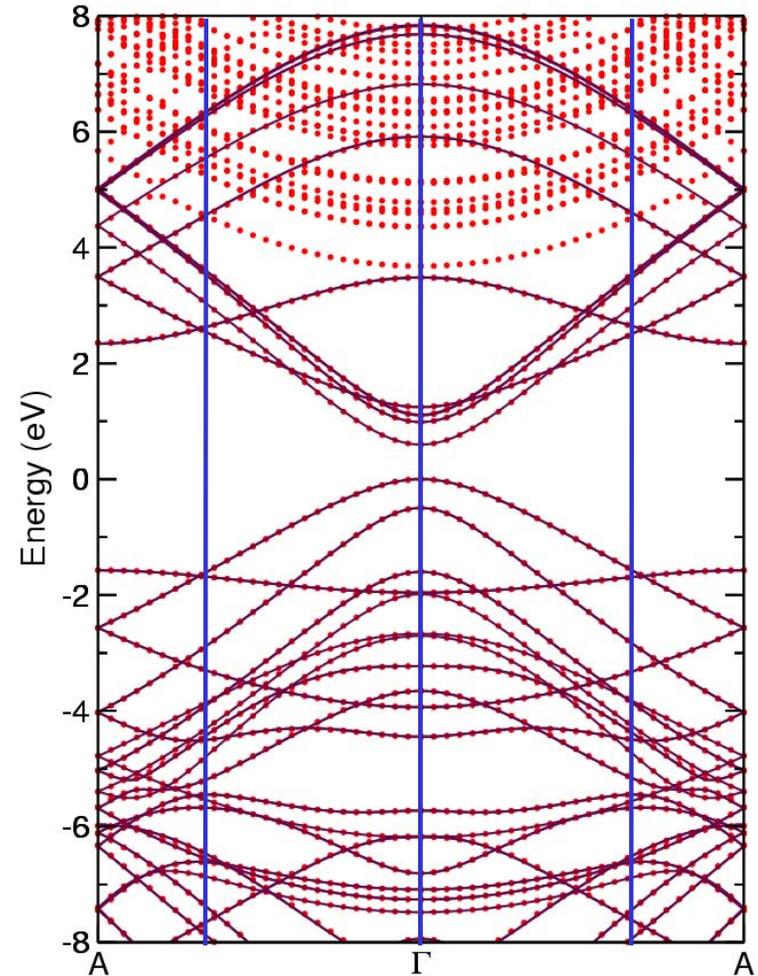
$$\langle \psi_{i\mathbf{k}} | \hat{H} | \psi_{j\mathbf{k}} \rangle \quad H_{ij}^{00} + e^{i\mathbf{k}\cdot\mathbf{R}} H_{ij}^{01} + e^{-i\mathbf{k}\cdot\mathbf{R}} H_{ij}^{0\bar{1}} \quad \Rightarrow \text{Diagonalize H Matrix}$$

Band Structure of (5,5) and (8,0) SWCNTs

(5,5) SWCNT



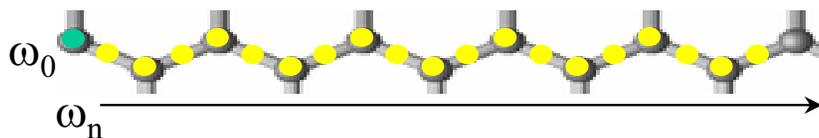
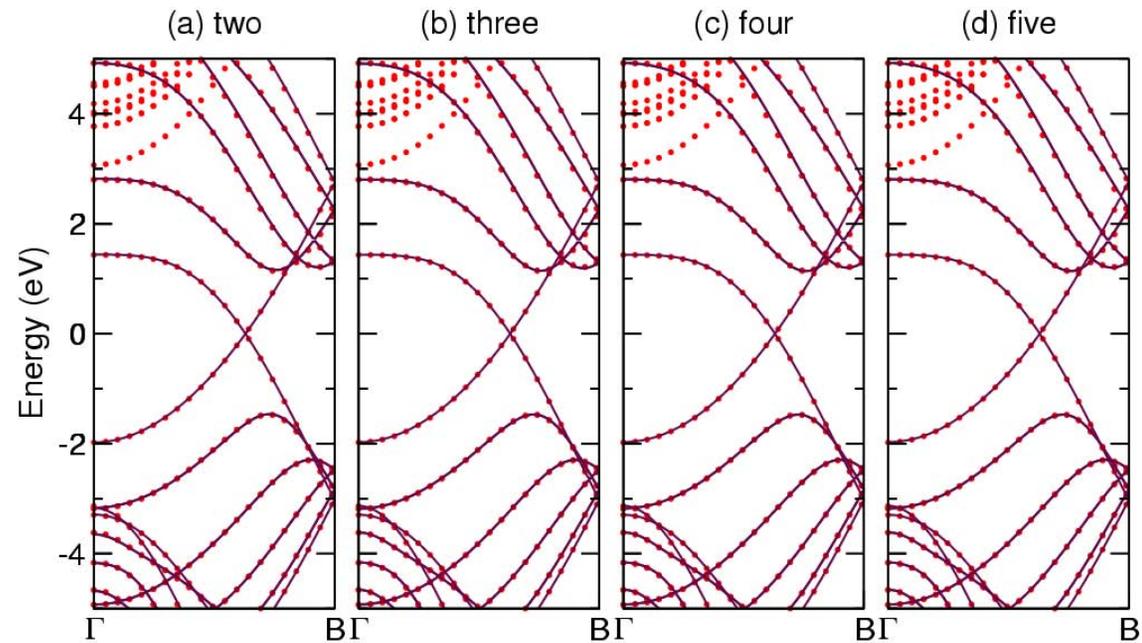
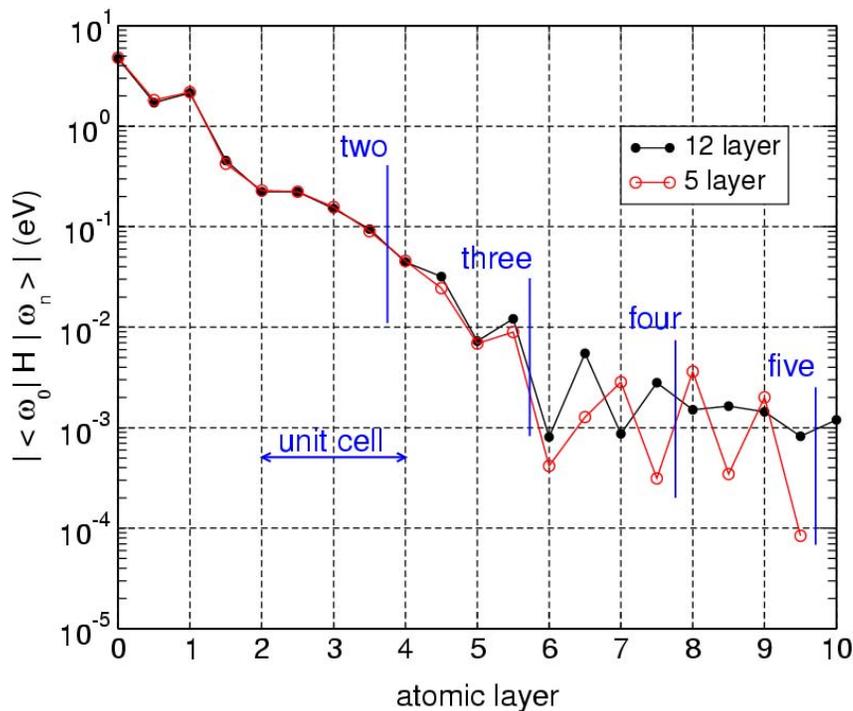
(8,0) SWCNT



Exponential Decay

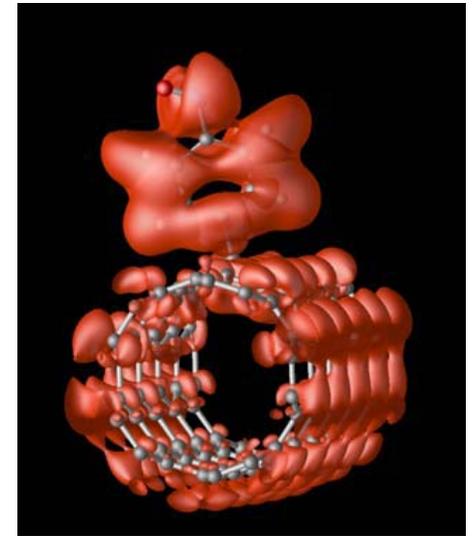
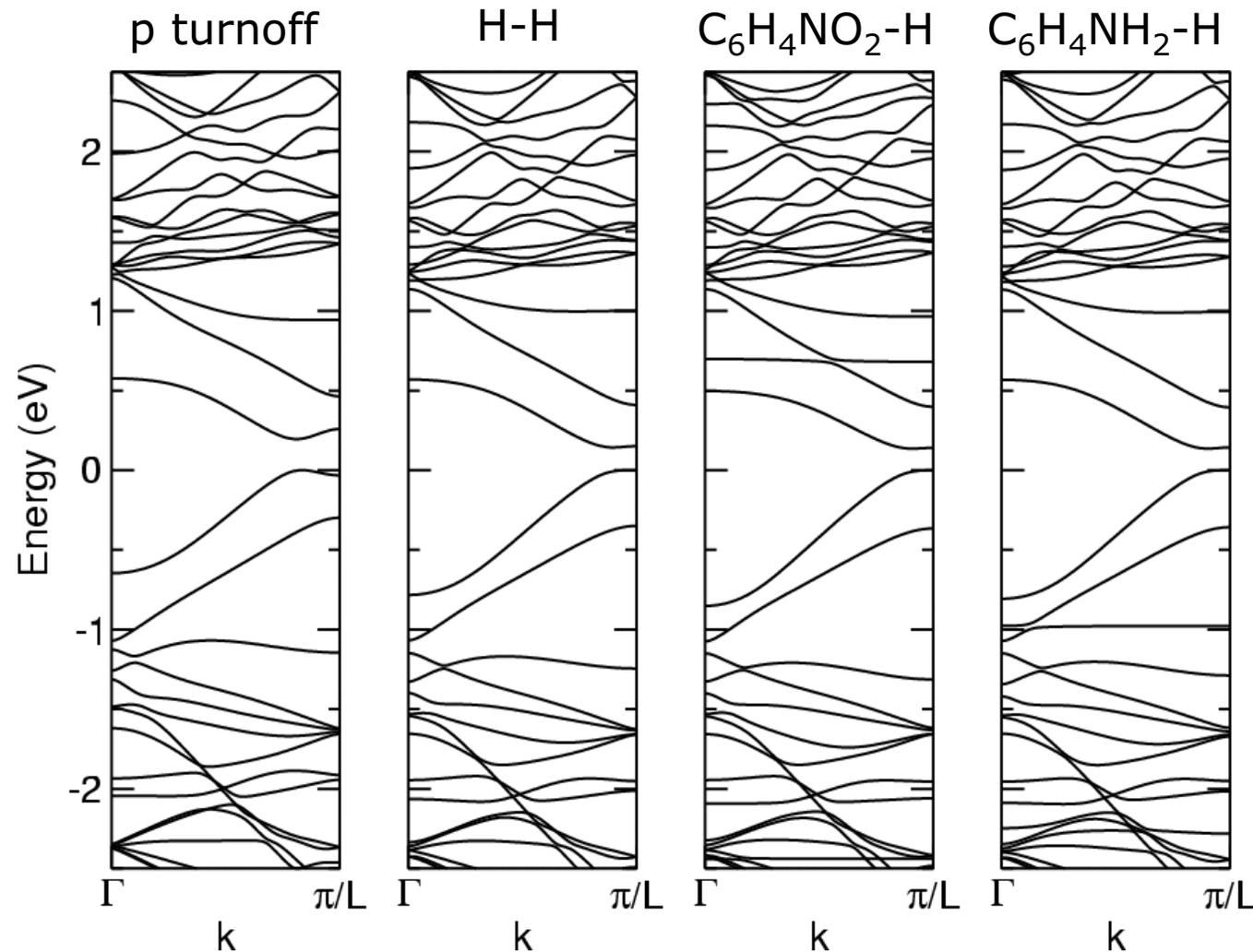
Even in a metal, we have smoothly connected manifolds – no relation with the physical decay of the density matrix

$$\langle \omega_0 | \hat{H} | \omega_n \rangle$$



Aryl / Hydrogen

- Band structure / Different functional groups
 - Band structure is not so much dependent on functional groups
 - Net effect of covalent functionalization can be mimicked by removing p-MLWFs \Rightarrow **localized** p orbitals

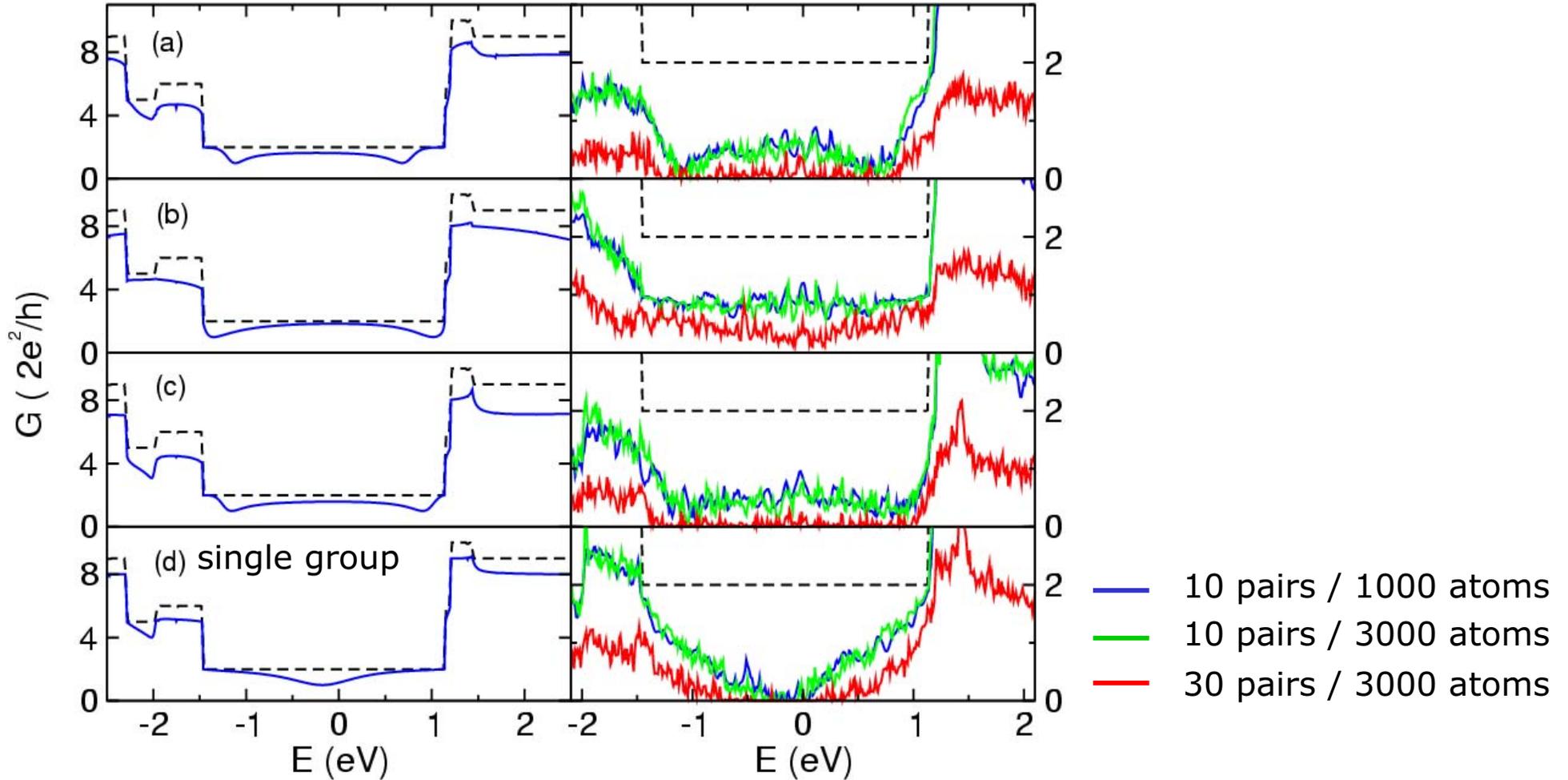


charge density injected when $-NO_2$ replaced by $-NH_2$

aminophenyl group is more electropositive than nitrophenyl group

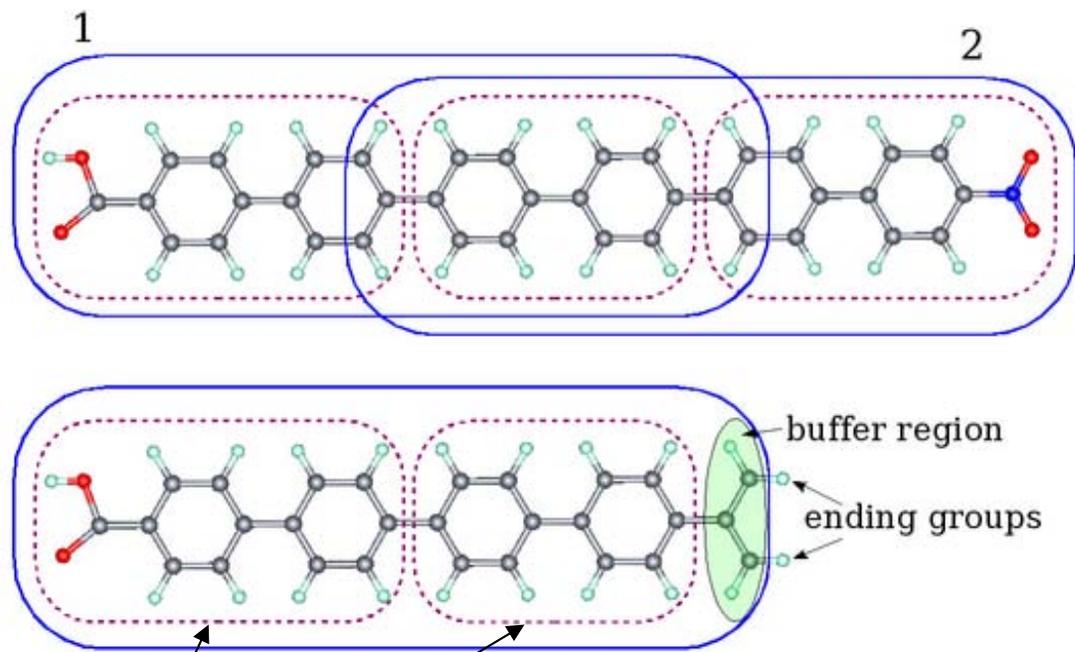
Aryl / Hydrogen

- Large scale simulations
 - Conductor part : 1000-3000 atoms / 12-37 nm
 - Average of five random configurations



Large Scale Calculations

- Parameterization



- principal layer

- Charge density convergence
- Same MLWFs and Hamiltonian matrix elements in overlapping region
- Electronic structure of long 1-D structure with a little cost

MLWFs Extraction From a Saturated Cluster

