

Phase Transformations: Spinodal Decomposition

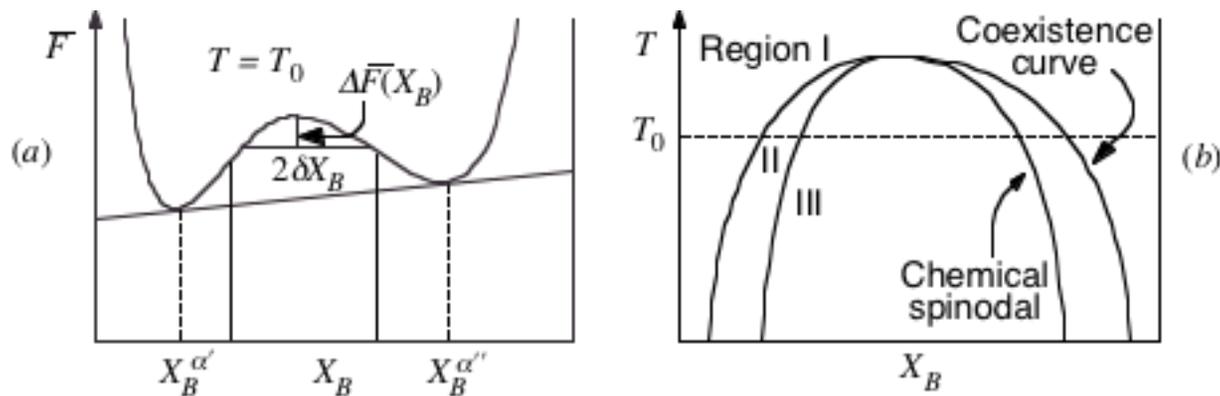


Today's topics:

- **A phase diagram with a spinodal region**
- **The free energy of a compositionally inhomogeneous solution**
- **Interdiffusion within the spinodal region**
- **The kinetics of spinodal decomposition**
- **Spinodal microstructures**

A phase diagram with a spinodal region

- The chemical spinodal is defined by $\frac{\partial^2 F}{\partial X_B^2} = 0$ and it coincides with the limit to metastability with respect to infinitesimal fluctuations of composition.

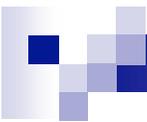


Spinodal decomposition becomes possible in region III where $\frac{\partial^2 F}{\partial X_B^2} < 0$



Free energy of inhomogeneous system

- **Thermodynamics of solutions normally describes systems which are compositionally homogeneous, at least on a local scale.**
- **Real materials often contain composition gradients or composition discontinuities (e.g., consider a binary alloy in a two-phase equilibrium state).**
- **It is possible to develop a description for the free energy of a phase in which there is an arbitrary compositional variation $c(x)$.**

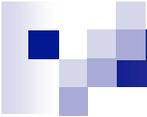


Free energy of inhomogeneous system...

- **Cahn and Hilliard (1958) used a variational approach to formulate a general expression for $F[c(x)]$ in a system of cross sectional area A with a one-dimensional compositional variation**

$$F[c(x)] = A \int_{-\infty}^{\infty} \left[f(c) + K \left(\frac{dc}{dx} \right)^2 \right] dx$$

K is a positive materials constant called the *gradient-energy coefficient* and $f(c)$ is the free energy for a homogeneous system.



Interdiffusion in the spinodal region

- In *any* system with composition gradients, gradient energy will modify the local driving force for diffusion. Practically speaking, gradient-energy effects are only important when the material contains a high density of regions with gradients (thin-film multilayers!).
- When gradient-energy effects are considered, the diffusion equation in 1-D becomes

$$\frac{\partial c}{\partial t} = \tilde{D} \left\{ \frac{\partial^2 c}{\partial x^2} - \left(2K / \frac{\partial^2 f}{\partial c^2} \right) \frac{\partial^4 c}{\partial x^4} \right\}$$

Solution to modified diffusion equation

- **The interdiffusivity is negative in the spinodal region**

$$\tilde{D} = M_c \frac{\partial^2 f}{\partial c^2}$$

(M_c is a mobility and is inherently positive).

- **In the spinodal region the solution to the diffusion equation has the form**

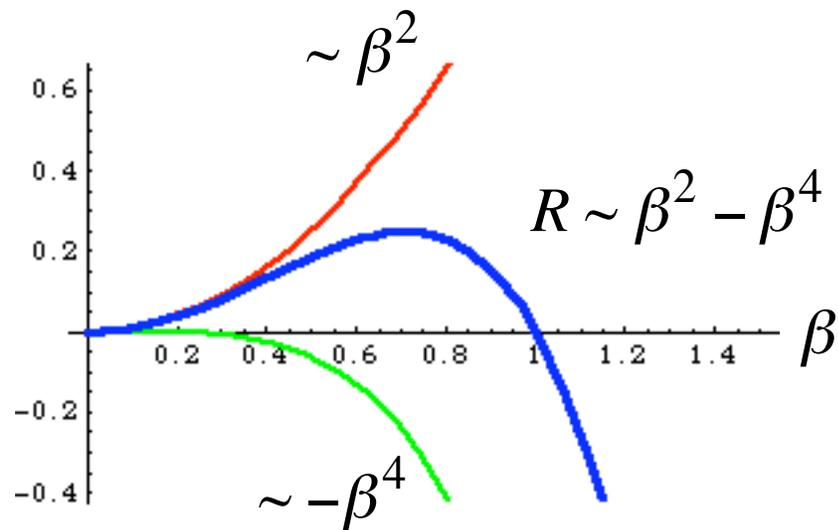
$$c = c_0 + e^{R(\beta)t} \cos(\beta x)$$

with

$$R(\beta) = -\tilde{D}\beta^2 \left[1 + \left(\frac{2K}{\partial^2 f / \partial c^2} \right) \beta^2 \right]$$

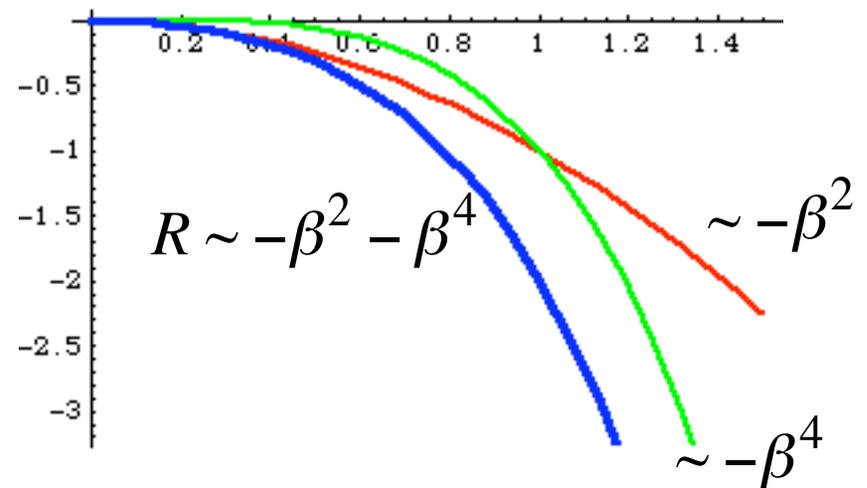
Solution to modified diffusion equation...

- The amplification factor $R(\beta)$ is a fourth-order polynomial with this form



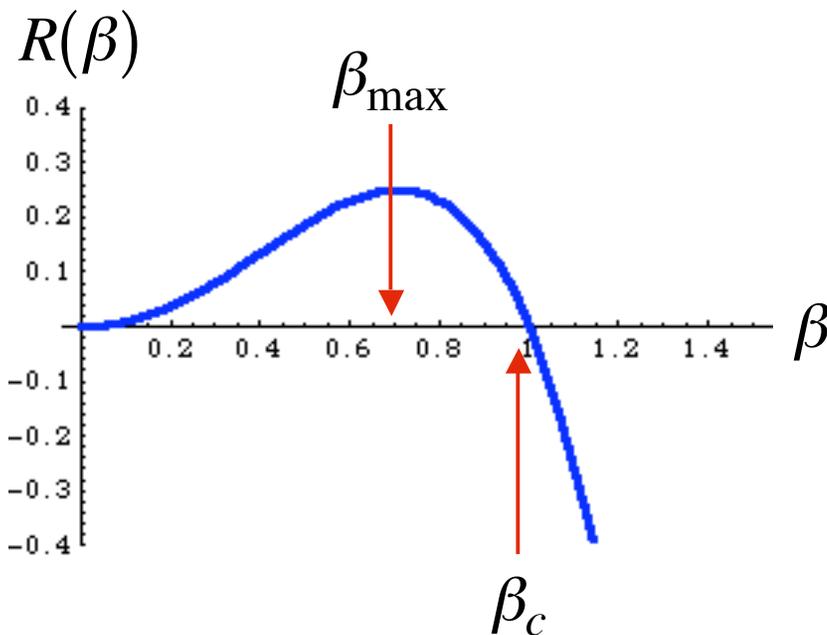
***Inside* spinodal
unstable waves grow**

***Outside* spinodal
all waves decay**



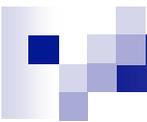
Solution to modified diffusion equation...

- Inside the spinodal, for small β , $R(\beta)$ is positive and it has a maximum at β_{\max} .



- Microstructure of developing structure is *periodic* with wavelength determined by

$$\lambda_{\max} = \frac{2\pi}{\beta_{\max}}$$



Solution to modified diffusion equation...

- β_{\max} is determined from the condition

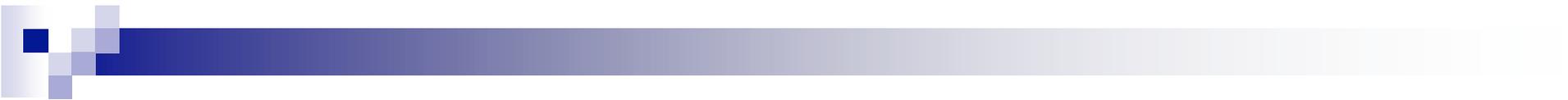
$$\frac{dR(\beta)}{d\beta} = 0 \quad \text{and thus} \quad 2\beta_{\max} + 8 \frac{K}{\partial^2 f / \partial c^2} \beta_{\max}^3 = 0$$

$$\text{so } \beta_{\max} = \sqrt{-\frac{\partial^2 f / \partial c^2}{4K}}$$

β_c is determined from the condition

$$R(\beta) = 0 \quad \text{and thus} \quad 1 + 2 \frac{K}{\partial^2 f / \partial c^2} \beta_c^2 = 0$$

$$\text{so } \beta_c = \sqrt{-\frac{\partial^2 f / \partial c^2}{2K}} \quad \text{and thus} \quad \beta_c = \sqrt{2} \beta_{\max}$$



Spinodal microstructures

- **The spinodal instability and the rapid growth of waves at β_{\max} leads to periodic modulations of composition. In elastically anisotropic alloys, the spinodal structure tends to be aligned along elastically soft directions.**

Figures removed due to copyright restrictions.

See Figures 18.16 and 18.17 in Balluffi, Robert W., Samuel M. Allen, and W. Craig Carter.

Kinetics of Materials. Hoboken, NJ: J. Wiley & Sons, 2005. ISBN: 0471246891.



Spinodal microstructures, cont'd

- **The wavelength of spinodal decomposition depends strongly on temperature because $\partial^2 f / \partial c^2$ is to first approximation a linearly decreasing function of temperature (becoming increasingly negative deeper inside the spinodal).**
- **As $\partial^2 f / \partial c^2$ becomes more negative, the wavelength of the spinodal structure gets smaller.**