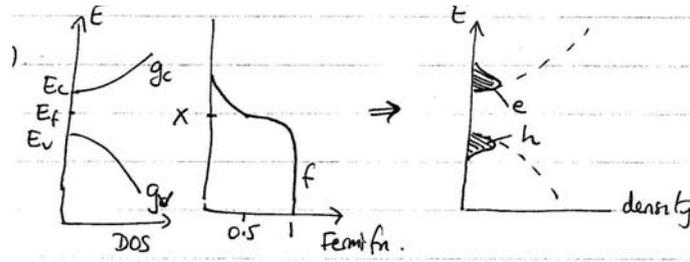


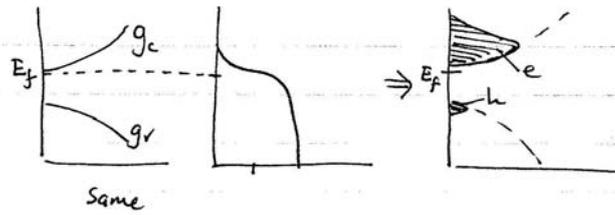
Sample Exam 1 - Solutions

Problem 1

- a.
i.



- ii.



- b.

$$n_i = \sqrt{N_c N_v} \exp \frac{-E_g}{2kT}$$

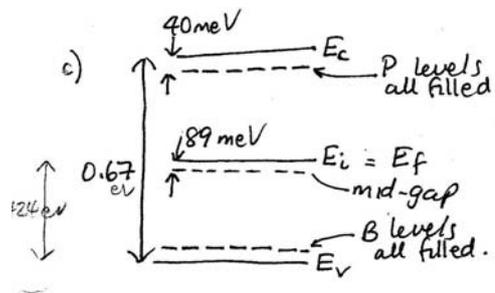
For Si: $n_i = 10^{10} \text{cm}^{-3}$, $E_g = 1.12 \text{eV}$

$$\begin{aligned} \text{Therefore: } \sqrt{N_c N_v} &= 10^{10} \exp(1.12/2 \times 0.0258) \\ &= 2.167 \times 10^{19} \text{cm}^{-3} \end{aligned}$$

$$\begin{aligned} \text{For Ge: } n_i &= 2.167 \times 10^{19} \exp -(0.67/2 \times 0.0258) \\ &= 6.1 \times 10^{13} \text{cm}^{-3} \end{aligned}$$

This is 6000 times larger because band gap is smaller.

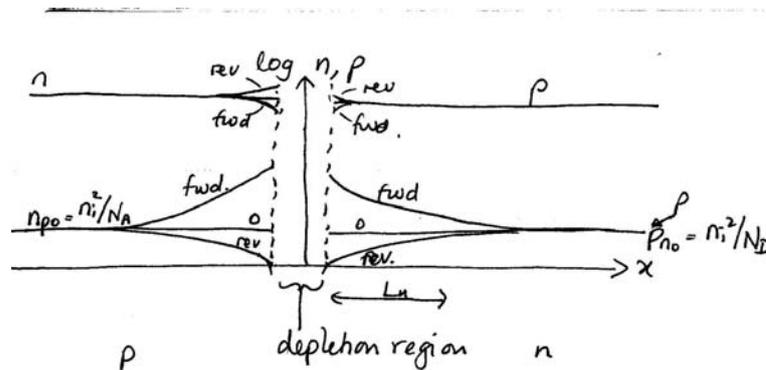
c. This is *compensated* (donors and acceptors cancel). Since $N_A = N_D$,



$E_f = E_i$, and $n = p = n_i$. Difference between E_i and midgap:

$$\begin{aligned} E_i &= \text{midgap} + \frac{3}{4} kT \ln \frac{m_p^*}{m_n^*} \\ &= \frac{3}{4} \times 0.0256 \ln 100 \\ &= 89 \text{ meV} \end{aligned}$$

d. Lower because compensated material's dopants scatter carriers, lowering mobility.



Problem 2

Problem 1

a.

****FIGURE****

Forward Bias: e^- injected into p side, leads to high e concentration near depletion region, decays away with characteristic length L_n (10s of μm) slight reduction of electron concentration near depletion region of n side.

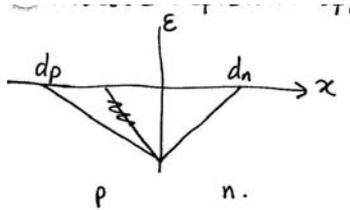
Reverse Bias: e collected from p side, pins $n = 0$. (Very slight increase in n on the n side.)

b.

Breakdown: ϵ in depletion region = 10^5V/cm .

Assume depletion approximation: if $N_D = N_A$, $d_n = d_p$.

$$\text{Max } \epsilon = \frac{N_A e d_p}{E_D E_R} \text{ or } \frac{N_D e d_n}{E_D E_R}$$



$$\text{We know: } d_n = \left(\frac{2E_D E_R V_0}{e} \frac{N_A}{N_D(N_A + N_D)} \right)^{\frac{1}{2}}$$

and

$$V_0 = \frac{kT}{e} \ln \left(\frac{N_A N_D}{n_i^2} \right) = 0.0258 \ln \left(\frac{10^{15} \cdot 10^{15}}{10^{20}} \right)$$

So:

$$\begin{aligned} d_n &= \left(\frac{2 \times 1.05 \cdot 10^{-12} \times 0.59}{1.6 \cdot 10^{-19}} \times \frac{10^{15}}{10^{15} (2.5 \times 10^{15})} \right)^{\frac{1}{2}} \\ &= 6.2 \times 10^{-5} \text{cm or } 0.6 \mu\text{m} \end{aligned}$$

Therefore:

$$\epsilon = \frac{N_D e d_n}{E_D E_r} = \frac{10^{15} \times 1.6 \cdot 10^{-19} \times 6 \cdot 10^{-5}}{1.05 \cdot 10^{-12}} = 0.9 \cdot 10^4 \text{V/cm}$$

For breakdown, applying voltage V_A extends the depletion region, and raises ϵ . d_n increases by a factor of $\sqrt{\frac{V_0 + V_A}{V_0}}$ and so does ϵ .

If avalanche occurs at 10^5 V/cm:

$$\frac{10^5}{0.9 \times 10^4} = \sqrt{\frac{V_0 + V_A}{V_0}}$$

$$V_0 + V_A = 72.8 \text{V}$$

So reverse bias of $\approx 72 \text{V}$ is needed for breakdown.

c.

We collect carriers in reverse bias pn-jn so we might bias both jns in reverse. Then any holes produced in base will go across either EB or BC-jn, depending on which direction they go. A large hole current flows into E and C from B. Electrons flow out of the base contact. $I_B \propto$ light intensity. (In fact it would work even without bias.)

Note: Putting EB in fwd bias means a large current already flows through E \rightarrow C, so the light-generated carriers would not add much to this.

Problem 3

a.

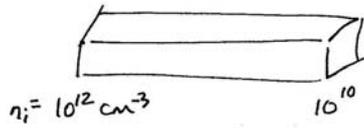
As T increases, μ decreases

As doping level increases, μ decreases

As lattice defects increase, μ decreases

All since greater scattering $\mu = \frac{e\tau}{m^*}$. Also, μ is lower for heavier carriers.

b.



p type: $p = N_A = 10^{18} \text{cm}^{-3}$ for both temperatures. Fully ionized.

Hot End: $n = 10^5 \text{cm}^{-3} = n_i^2 / N_A$

Cold End: $n = 10^2 \text{cm}^{-3}$

So we have diffusion of electrons until the electric field balances the concentration gradient. Neglect hole diff.

Thermal R & G occurs everywhere, but more carriers at hot end.

$$J_n = eD_n \frac{dn}{dx} = en\mu_n\epsilon \text{ at steady state. Diffusion = Drift.}$$

$$\text{Concentration Gradient} = \frac{10^6 - 10^2}{1 \text{cm}} \text{cm}^{-3} \approx 10^6 \text{cm}^{-4}$$

Substitute:

$$D_n = kT\mu_n/e$$

$$\text{Therefore, } kT\mu_n \frac{dn}{dx} = en\mu_n\epsilon$$

$$\epsilon = \frac{kT}{en} \frac{dn}{dx}$$

At the hot end:

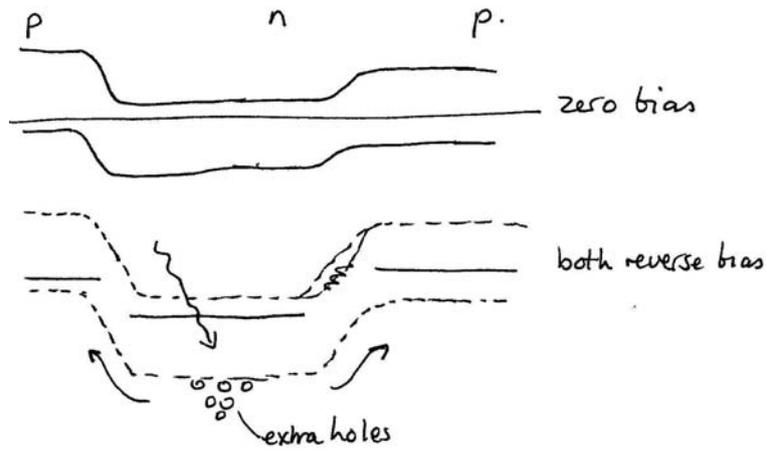
$$n = 10^6 \text{ cm}^{-3}, \epsilon = \frac{kT}{e} \cdot \frac{10^6 \text{ cm}^{-4}}{10^6 \text{ cm}^{-3}} = \frac{kT}{e} = 0.026 \text{ V/cm}$$

At the cold end:

$$n = 10^2 \text{ cm}^{-3}, \epsilon = 260 \text{ V/cm}.$$

d.

We collect carriers in reverse bias pn-jn so we might bias both jns in reverse. Then any holes produced in base will go across either EB or BC jn, depending on which direction they go. A large hole current flows into E and C from B. Electrons flow out of the base contact. $I_B \propto$ light intensity. (In fact it would work even without bias.)



Note: Putting EB in fwd bias means a large current already flows through E \rightarrow C, so the light-generated carriers would not add much to this.