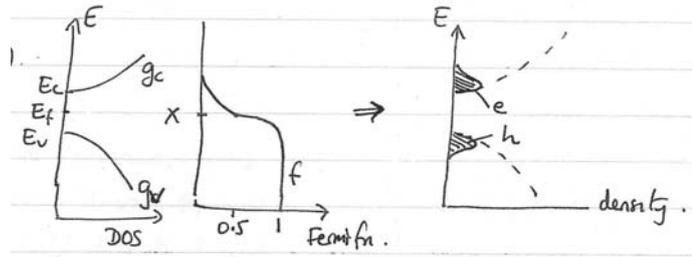
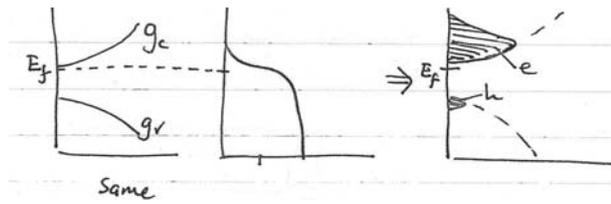


3.15 - Problem Set 1 Solutions

2.
a. (i)



- (ii)



- b.

$$n_i = \sqrt{N_c N_v} \exp \frac{-E_g}{2kT}$$

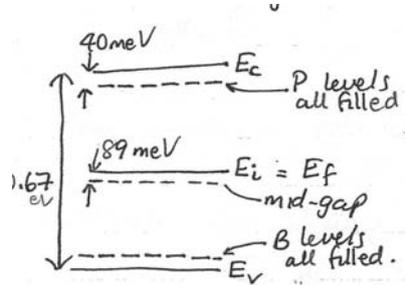
For Si: $n_i = 10^{10} \text{cm}^{-3}$, $E_g = 1.12 \text{eV}$

$$\begin{aligned} \text{Therefore: } \sqrt{N_c N_v} &= 10^{10} \exp(1.12/2 \times 0.0258) \\ &= 2.167 \times 10^{19} \text{cm}^{-3} \end{aligned}$$

$$\begin{aligned} \text{For Ge: } n_i &= 2.167 \times 10^{19} \exp -(0.67/2 \times 0.0258) \\ &= 6.1 \times 10^{13} \text{cm}^{-3} \end{aligned}$$

This is 6000 times larger because band gap is smaller.

c.



This is *compensated* (donors and acceptors cancel). Since $N_A = N_D$, $E_f = E_i$, and $n = p = n_i$. Difference between E_i and midgap:

$$\begin{aligned}
 E_i &= \text{midgap} + \frac{3}{4}kT \ln \frac{m_p^*}{m_n^*} \\
 &= \frac{3}{4} \times 0.0256 \ln 100 \\
 &= 89 \text{ meV}
 \end{aligned}$$

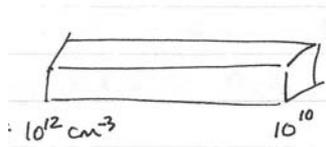
d. Lower because compensated material's dopants scatter carriers, lowering mobility.

Problem 3

- a. As T increases, μ decreases
 As doping level increases, μ decreases
 As lattice defects increase, μ decreases

All since greater scattering $\mu = \frac{e\tau}{m^*}$. Also, μ is lower for heavier carriers.

b.



p type: $p = N_A = 10^{18} \text{cm}^{-3}$ for both temperatures. Fully ionized.

Hot End: $n = 10^5 \text{cm}^{-3} = n_i^2 / N_A$

Cold End: $n = 10^2 \text{cm}^{-3}$

So we have diffusion of electrons until the electric field balances the concentration gradient. Neglect hole diff.

Thermal R & G occurs everywhere, but more carriers at hot end.

$$J_n = eD_n \frac{dn}{dx} = en\mu_n \epsilon \text{ at steady state. Diffusion} = \text{Drift.}$$

$$\text{Concentration Gradient} = \frac{10^6 - 10^2}{1 \text{cm}} \text{cm}^{-3} \approx 10^6 \text{cm}^{-4}$$

Substitute:

$$D_n = kT\mu_n/e$$

$$\text{Therefore, } kT\mu_n \frac{dn}{dx} = en\mu_n \epsilon$$

$$\epsilon = \frac{kT}{en} \frac{dn}{dx}$$

At the hot end:

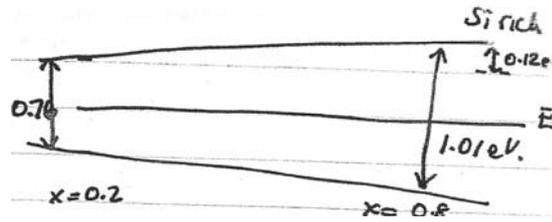
$$n = 10^6 \text{cm}^{-3}, \epsilon = \frac{kT}{e} \cdot \frac{10^6 \text{cm}^{-4}}{10^6 \text{cm}^{-3}} = \frac{kT}{e} = 0.026 \text{V/cm}$$

At the cold end:

$$n = 10^2 \text{cm}^{-3}, \epsilon = 260 \text{V/cm}.$$

Problem 4

$x=0, 0.67 \text{ eV}$
 $x=0.2, 0.76 \text{ eV}$
 $x=0.8, 1.01 \text{ eV}$
 $x=1.0, 1.1 \text{ eV}$



Intrinsic - assume E_f is very close to midgap.

Electrons \leftarrow drift due to band slope.

Electrons \rightarrow diffuse since more e^- in narrow gap.

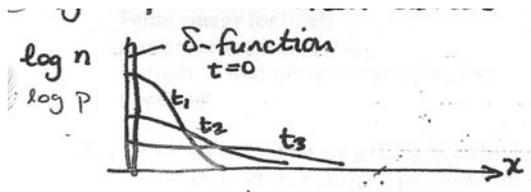
Holes \rightarrow diffuse.

Holes \leftarrow drift.

The only R & G is thermal.

Problem 5

Light produces excess carriers at surface.



Intrinsic so $n = p$. Area under curve decreases due to recombination. Curve spreads due to diffusion.

$$\frac{\partial n}{\partial t} = \frac{\partial n}{\partial t}_{\text{drift}} + \text{diffusion} + \text{R \& G.}$$

Drift $\rightarrow 0$ because no fields.

Diffusion $\rightarrow \frac{1}{e} \nabla J = \frac{1}{e} \frac{d^2 n}{dx^2}$.

R & G: $\frac{dn}{dt} = -\frac{n_e}{\tau}$ where $n_e =$ excess carrier concentration.

$$\frac{dn}{dt} = D_n \frac{d^2 n}{dx^2} + -\frac{n_e}{\tau}$$

Same type of expression for holes.

τ is due to recombination at traps, since it's intrinsic material. These are more effective than band to band recombination but there are relatively few of them.

$\tau = 1/r_2 N_\tau$. $r_2 \gg r$ but $N_\tau <$ typical.