

# 2 Types of Noncrystalline Polymers

1. Glassy polymer
  2. Rubbery polymers
- } highly interpenetrated/entangled  
random Gaussian coils

## Glass Transition Temperature

$$\epsilon_{ij}$$

Two viewpoints:

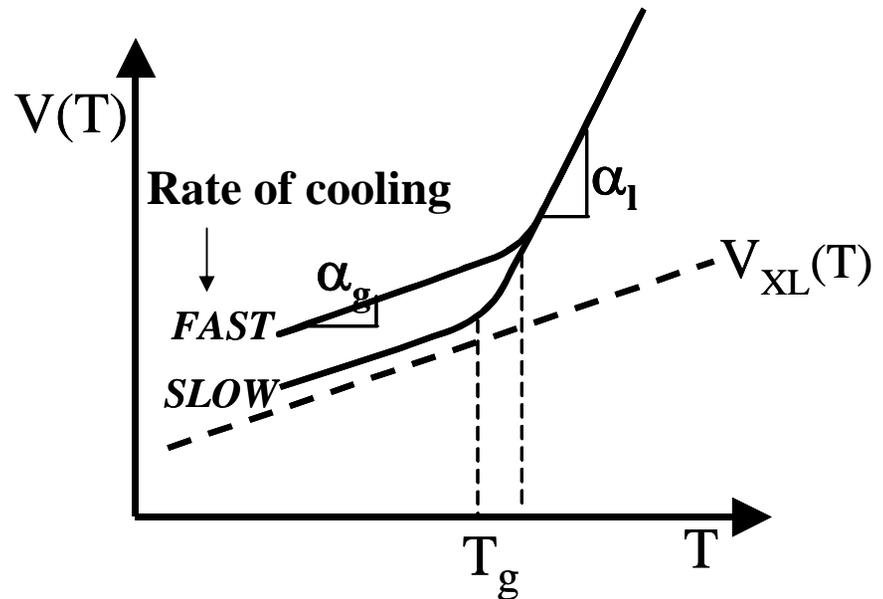
- Increasing T: When  $kT >$  magnitude of  $\epsilon_{ij}$ , the thermal fluctuations can overcome local intermolecular bonds and the frozen (“glassy”) structure becomes “fluid-like”.
- Decreasing T: As the temperature is lowered and T approaches  $T_g$ , the viscosity increases to  $\infty$  and the material becomes “solid”

# Free Volume Theory of $T_g$

Free volume,  $V_F$  – extra space beyond what is present in an ordered crystalline packing (beyond the interstitial volume).

$$V_F(T) \equiv V(T) - V_0(T)$$

- $V_0$  is occupied specific volume of atoms or molecules in the xline state *and* the spaces between them:  $\sim V_{XL}$ .
- $V_F$  increases as  $T$  increases due to the difference in the thermal expansion coefficients ( $\alpha_g$  vs  $\alpha_l$ ).
- $V_0(T) \approx V_{XL}(T) \leftrightarrow$  can take  $\alpha_g \approx \alpha_{XL}$
- $V_F(T) = V_F(T_g) + (T-T_g)\frac{dV_F}{dT} \quad T > T_g$
- define fractional free volume,  $f_F$ :  $V_f/V$



$$f_F(T) = f_F(T_g) + (T-T_g)\alpha_f \quad \alpha_f = \alpha_l - \alpha_g$$

**Viewpoint:**  $T_g$  occurs when available free volume drops below critical threshold for structural rearrangement [VITRIFICATION POINT], *structure “jams up”*.

# T<sub>g</sub> Values of Amorphous Materials

Table of representative amorphous solids, their bonding types, and their glass transition temperatures removed due to copyright restrictions.

See Table 2.2 in Allen, S. M., and E.L. Thomas. *The Structure of Materials*. New York, NY: J. Wiley & Sons, 1999.

# T<sub>g</sub> for Selected Polymers

## Glass Transition Temperature for Selected Polymers

Organic Polymer	T <sub>g</sub> (°C)
Polyacenaphthalene	264
Polyvinyl pyrrolidone	175
Poly- <i>o</i> -vinyl benzyl alcohol	160
Poly- <i>p</i> -vinyl benzyl alcohol	140
Polymethacrylonitrile	120
Polyacrylic acid	106
Polymethyl methacrylate	105
Polyvinyl formal	105
Polystyrene	100
Polyacrylonitrile	96
Polyvinyl chloride	87
Polyvinyl alcohol	85
Polyvinyl acetal	82
Polyvinyl proprional	72
Polyethylene terephthalate	69
Polyvinyl isobutyral	56
Polycaprolactam (nylon 6)	50
Polyhexamethylene adipamide (nylon 6,6)	50
Polyvinyl butyral	49
Polychlorotrifluoroethylene	45
Ethyl cellulose	43

Organic Polymer	T <sub>g</sub> (°C)
Polyhexamethylene sebacamide (nylon 6,10)	40
Polyvinyl acetate	29
Polyperfluoropropylene	11
Polymethyl acrylate	9
Polyvinylidene chloride	-17
Polyvinyl fluoride	-20
Poly-1-butene	-25
Polyvinylidene fluoride	-39
Poly-1-hexene	-50
Polychloroprene	-50
Polyvinyl- <i>n</i> -butyl ether	-52
Polytetramethylene sebacate	-57
Polybutylene oxide	-60
Polypropylene oxide	-60
Poly-1-octene	-65
Polyethylene adipate	-70
Polyisobutylene	-70
Natural rubber	-72
Polyisoprene	-73
Polydimethyl siloxane	-85
Polydimethyl siloxane	-123

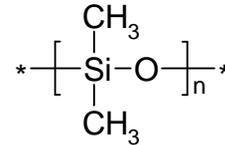
# Effects of Chemical Structure on $T_g$

- Onset of molecular mobility at  $T_g$  involves rotation of chain segments (comprised of  $\sim 10$ -30 repeat units) about the main chain. Such cooperative motion requires
  - #1 sufficient thermal energy ( $kT$  relative to  $\Delta V(\phi)$ ) for ease of rotation about main chain bonds and to overcome local bonding
  - #2 sufficient  $V_F$  for the segments to move into.
- Requirements for a low  $T_g$  polymer:
  1. weak interaction between chains  $\epsilon_{ij}$
  2. easy rotation about main chain bonds  $V(\phi)$
  3. abundant free volume  $V_F$

# Chemical Structure vs. $T_g$

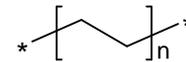
## Low $T_g$ Polymers

Polydimethyl siloxane



-123 °C

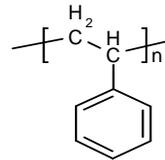
Polyethylene (noncrystalline regions,  
 $T_m = 140\text{C}$ )



-133 °C

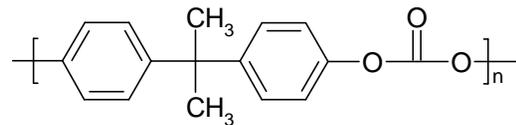
## High $T_g$ Polymers

Polystyrene



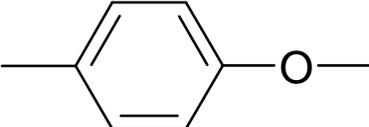
105 °C

Polycarbonate

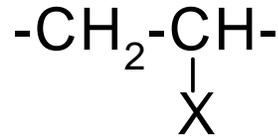


150 °C

# Influence of Backbone Flexibility on $T_g$

<u>Repeat Unit</u>	<u><math>T_g, ^\circ\text{C}</math></u>
$-\text{CH}_2-\text{CH}_2-$	-133
$-\text{CH}_2-\text{CH}_2-\text{O}-$	-67
	80
	84

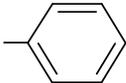
# $T_g$ : Influence of Steric Hinderance and $\epsilon_{ij}$



- steric hinderance

x = side group

$T_g$ , °C

<hr/>	<hr/>
-CH <sub>3</sub>	-23
-CH <sub>2</sub> -CH <sub>3</sub>	-24
-CH <sub>2</sub> -CH(CH <sub>3</sub> ) <sub>2</sub>	-16
-CH <sub>2</sub> -CH <sub>2</sub> -CH <sub>3</sub>	50
	100

- polar chain-chain interactions

-Cl	81
-OH	85
-CN	97

# Molecular Weight Dependence of $T_g$

Data show increase  $T_g$  with MW

Key Concept:

- *Chain ends provide extra space and freedom for motion*

Image removed due to copyright restrictions.

Please see Fig. 3 in Fox, Thomas G., and Flory, Paul J. "Second-Order Transition Temperatures and Related Properties of Polystyrene. I. Influence of Molecular Weight." *Journal of Applied Physics* 21 (June 1950): 581-591.

$$T_g(\overline{M}_n) = T_g^\infty - \frac{C}{\overline{M}_n}$$

# $T_g$ for Random Copolymers and Miscible Blends

- Random copolymers and miscible 2 component blend systems are homogeneous single phased materials and one can assume the rule of mixtures for fractional free volume of each component. This leads to simple relationships for  $T_g$

$$T_{g,co} = T_{g,A} w_A + T_{g,B} w_B$$

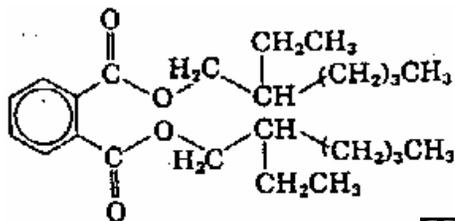
where  $w_i$  is the weight fraction of component  $i$

# Controlling $T_g$ with small molecule additives

## Plasticizers

- low molar mass additives
- act to spread chains apart
- act as lubricant
  - 1) mediate chain-chain polar interactions
  - 2) add free volume
- Example: adding bis(2-ethylhexyl)phthalate (DOP) to neat PVC can significantly lower the  $T_g$  of the blend

$$(T_{g, \text{PVC}} = 70^\circ\text{C}, \quad T_{g, \text{DOP}} = -80^\circ\text{C})$$



DOP

$$T_{g, \text{blend}} = T_{g, \text{DOP}} w_{\text{DOP}} + T_{g, \text{PVC}} w_{\text{PVC}}$$

# Controlling $T_g$ cont'd

## Antiplasticizers

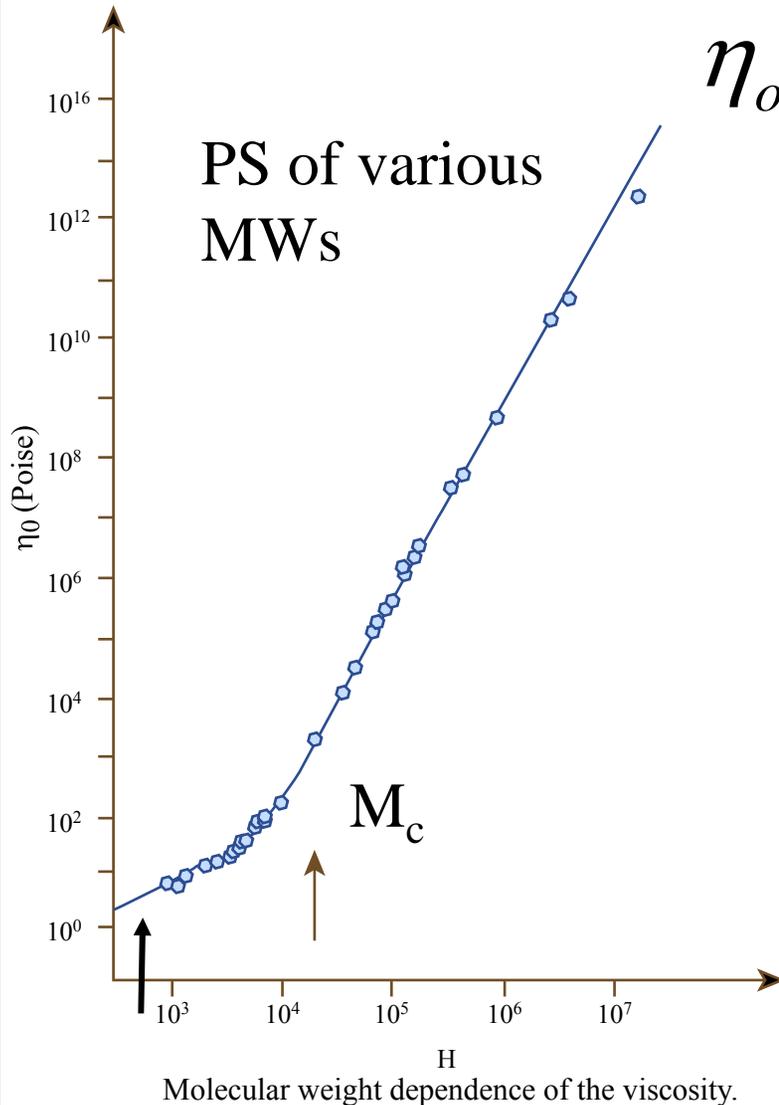
- Small molecule additives which “fit” into regions between chains, lowering the overall  $V_F$  and *raising* the  $T_g$  of the 2 component system

# $\eta_o(M)$ behavior for polymer melts

$\eta_o$  = zero shear rate melt viscosity

$M_c$   
 $M_e$

**Viscosity units: Poise**  
= dynes • s/cm<sup>2</sup>  
= g/cm • s

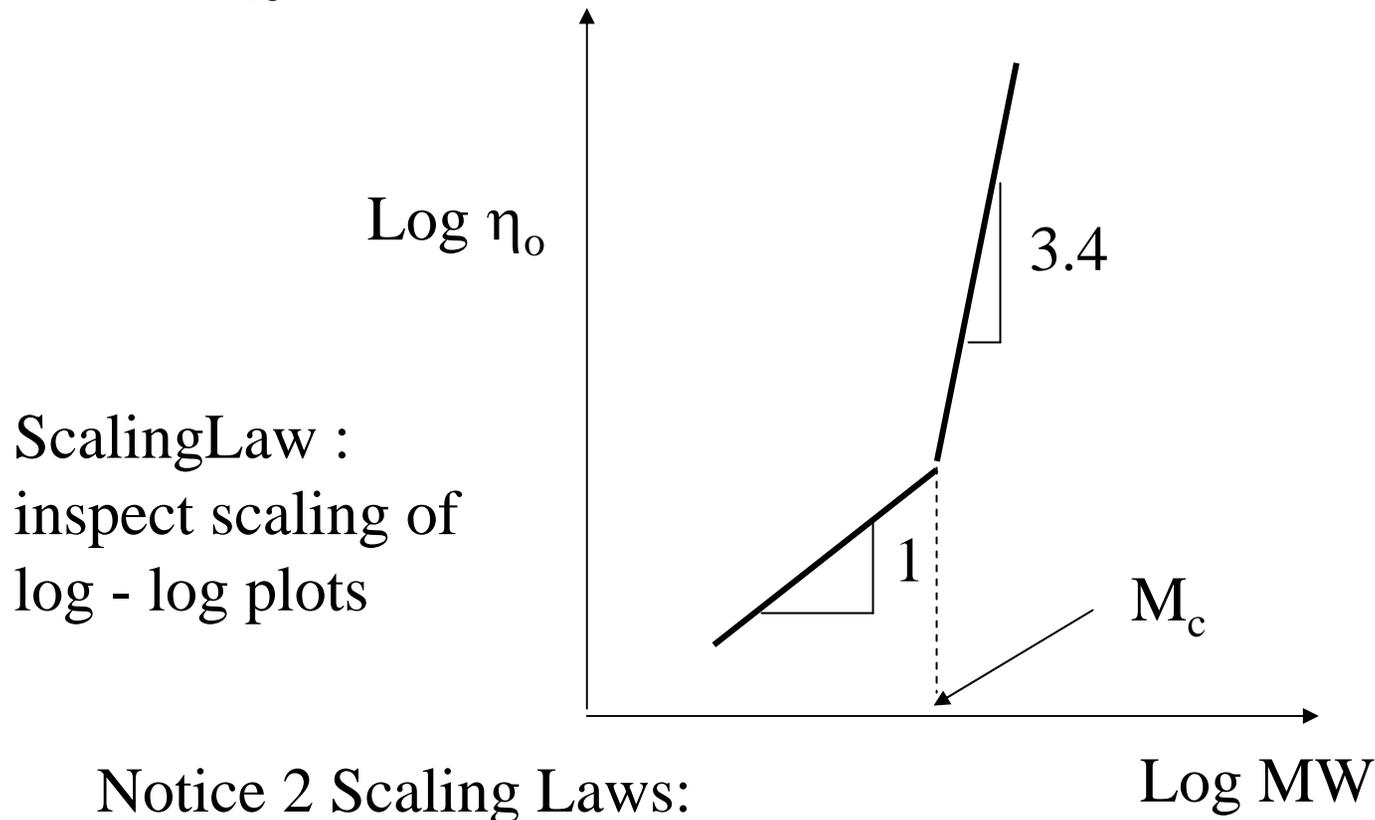


Characteristic molecular weights for selected Polymers.

Polymer species	$M_e$	$M_c$
Polystyrene	19 500	35 000
Poly( $\alpha$ -methyl styrene)	13 500	28 000
1,4-Polybutadiene	1 900	5 000
Polyvinyl acetate	12 000	24 500
Poly(dimethyl siloxane)	10 000	24 400
Polyethylene	1 300	3 800
1,4-Polyisoprene	6 300	10 000
Polyisobutylene	8 900	15 200

# Viscosity of Polymeric Melts

$\eta_0$  is the zero shear rate limit of the melt viscosity



Notice 2 Scaling Laws:

- For low MW
- For high MW

$\eta_0 \sim M$  called **Rouse** Regime

$\eta_0 \sim M^{3.4}$  called **Reptation** Regime

# Diffusivity (and Viscosity) of Polymer Melts

- Center of mass motion is important in determining the diffusivity (and viscosity) of a polymer melt.
- **Small Molecule Liquids** –move by random jumps into adjacent “holes” (free volume concept)  
typical  $D(20\text{ C}) \sim 10^{-5}\text{ cm}^2/\text{sec}$
- **Polymeric Liquids**
  - $D(20\text{ C}) \sim 10^{-14}$  to  $10^{-18}\text{ cm}^2/\text{sec}$
  - 2 regimes of diffusivity vs molecular weight are observed.  
 $D \sim M^{-1}$  and  $D \sim M^{-2}$  (more scaling laws!)

# More Scaling Laws – $D(N)$ behavior

Example: PS melts  
of various MWs

Image removed due to copyright restrictions.

Please see Fig. 2 in Watanabe, Hiroshi, and Kotaka, Tadao. "Viscoelastic and Diffusion Properties of Binary Blends of Monodisperse Polystyrenes." *Macromolecules* 20 (1987): 530-535.

Slope at lower MWs  
is about 1.0

Slope increases  
to  $\sim 2$   
for higher MWs

# Rouse Chain Model

- Rouse chain = a flexible connected string of Brownian particles that interact with a featureless background viscous medium.
- The number of repeat units is less than the entanglement limit, the chain has small  $N$ , where  $N < N_e$

- Viscosity depends on monomeric friction factor  $\xi_M$  and chain length  $x_2$

$$\eta \sim \xi_M \cdot x_2 \quad \text{so} \quad \eta \sim M^1$$

- Diffusivity depends on monomeric friction factor via Einstein

$$D = \frac{kT}{\xi} \quad \text{Hence, } D \sim M^{-1}$$