

# Lecture 4 Honeycombs Notes, 3.054

## Honeycombs-In-plane behavior

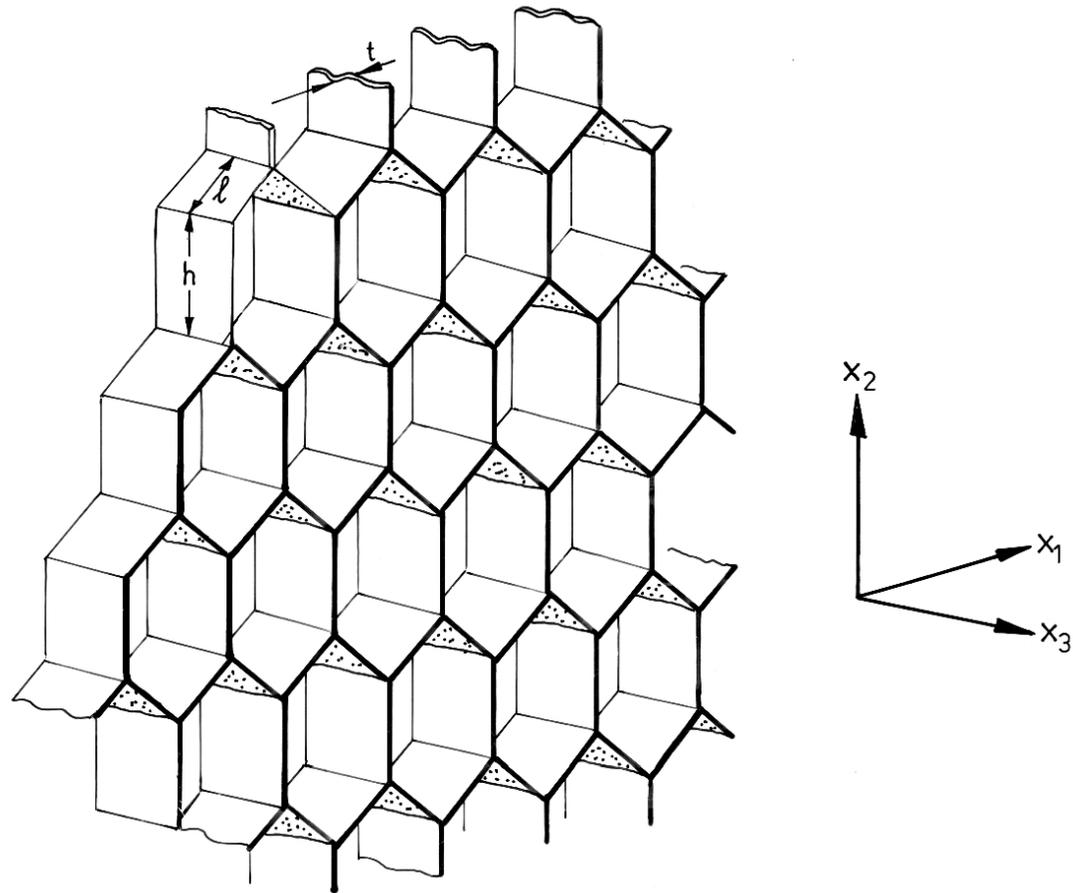
- Prismatic cells
- Polymer, metal, ceramic honeycombs widely available
- Used for sandwich structure cores, energy absorption, carriers for catalysts
- Some natural materials (e.g. wood, cork) can be idealized as honeycombs
- Mechanisms of deformation and failure in hexagonal honeycombs parallel those in foams
  - simpler geometry    unit cell    easier to analyze
- Mechanisms of deformation in triangular honeycombs parallel those in 3D trusses (lattice materials)

## Stress-strain curves and Deformation behavior: In-Plane

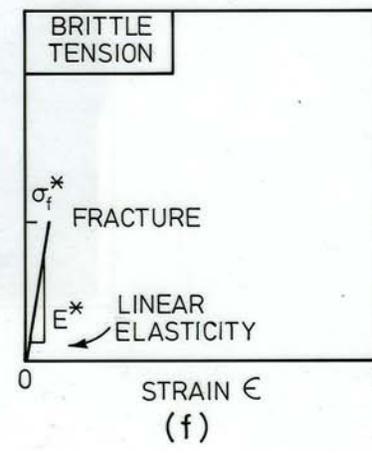
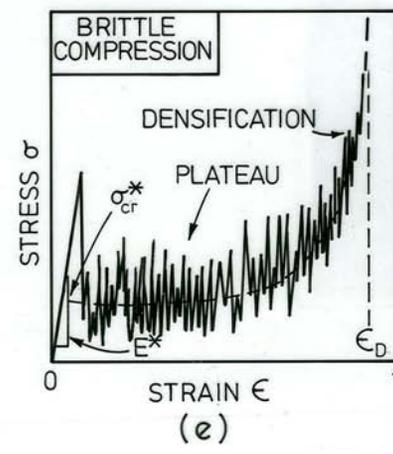
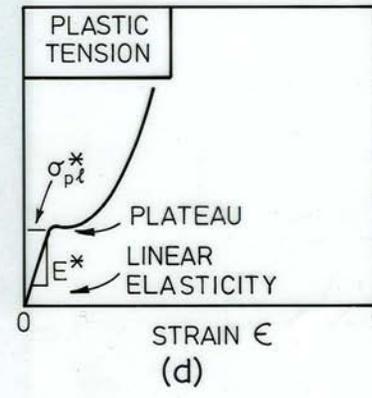
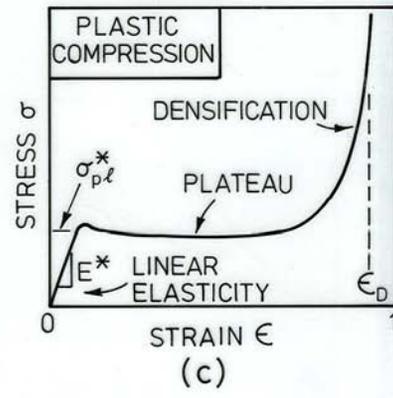
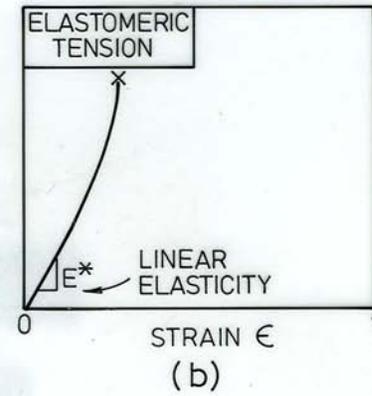
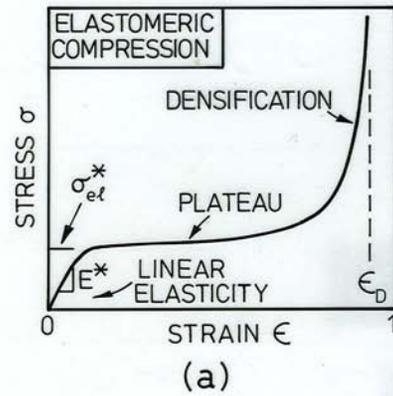
### Compression

- 3 regimes
  - linear elastic
  - stress plateau
  - bending
  - buckling
  - yielding
  - brittle crushing
  - densification
  - cell walls touch
- Increasing  $t/l$      $E^*$      $\sigma^*$      $\epsilon_D$

# Honeycomb Geometry

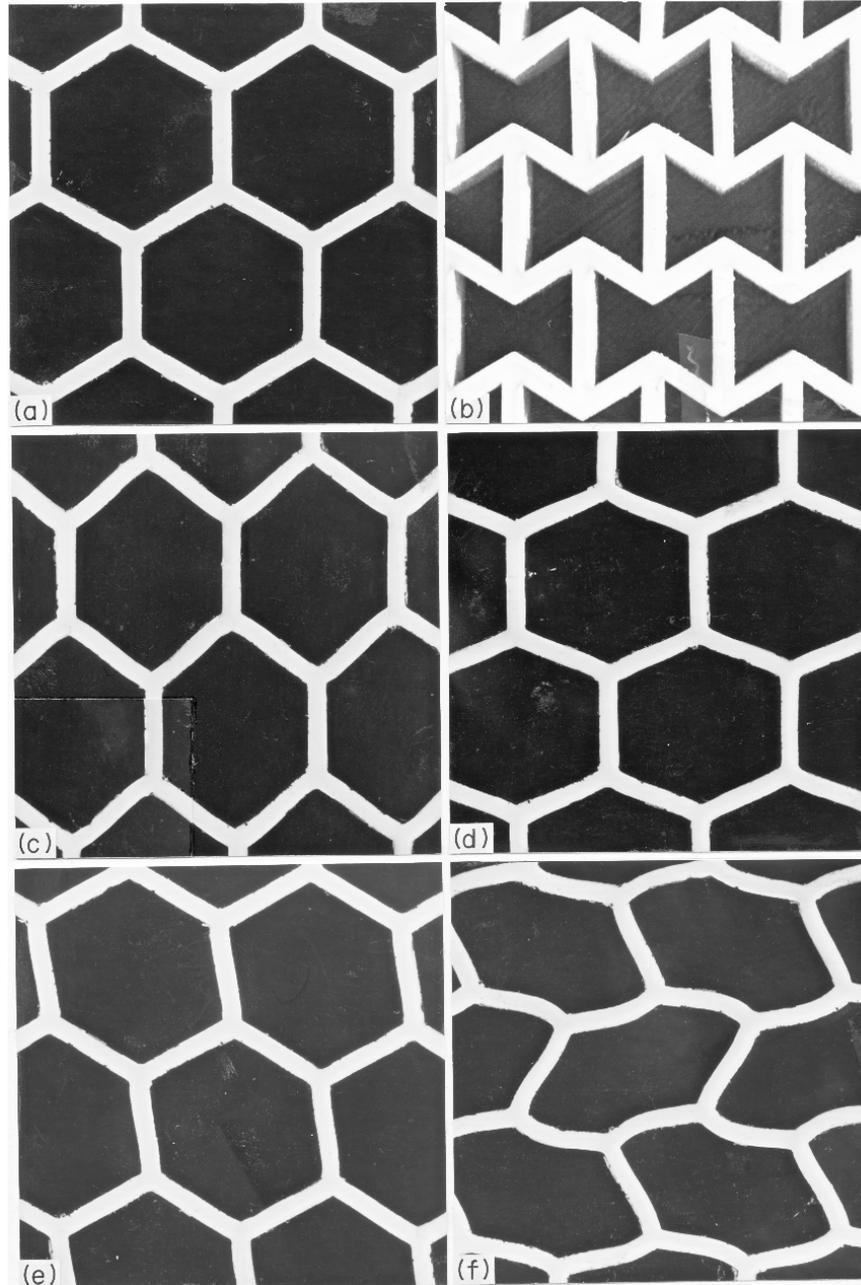


Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, © 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

# Deformation mechanisms

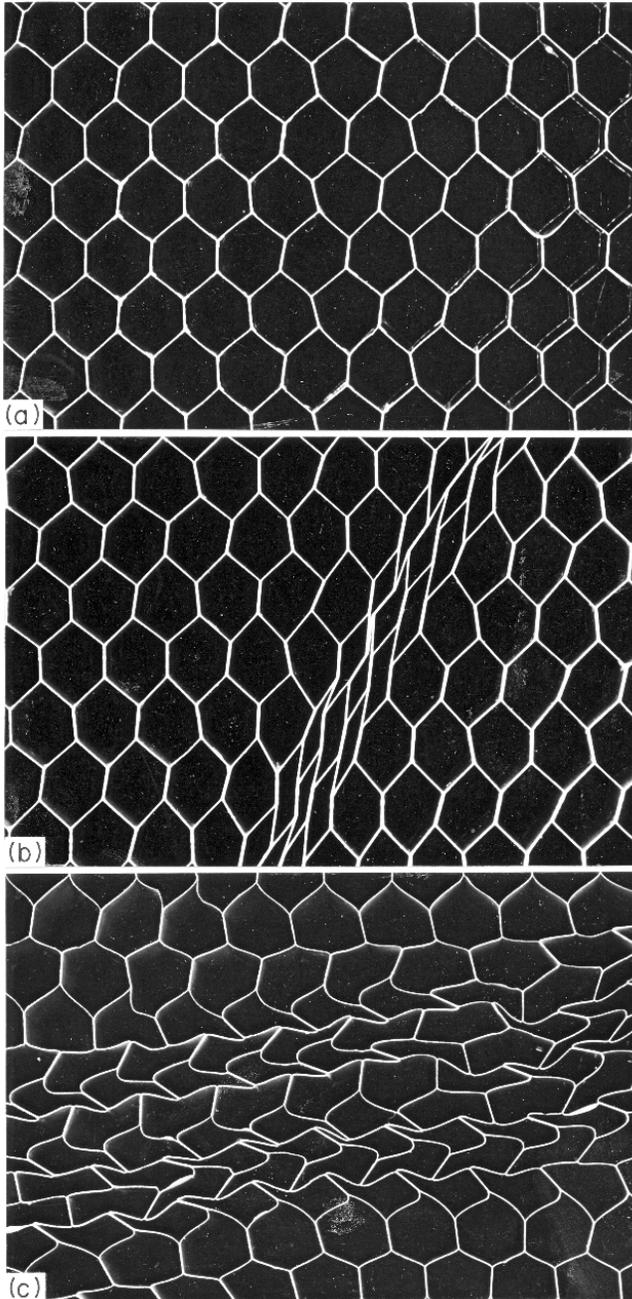


Bending  
 $X_1$  Loading

Bending  
 $X_2$  Loading

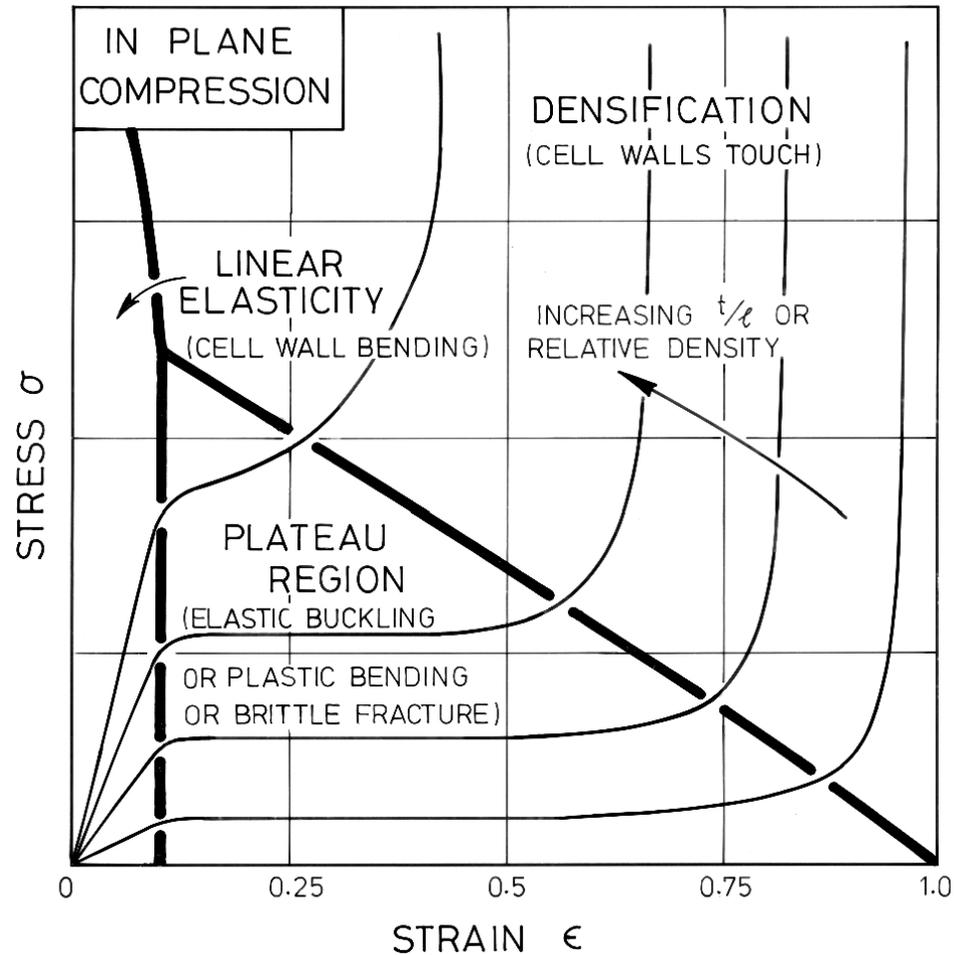
Bending  
Shear

Buckling



## Plastic collapse in an aluminum honeycomb

# Stress-Strain Curve



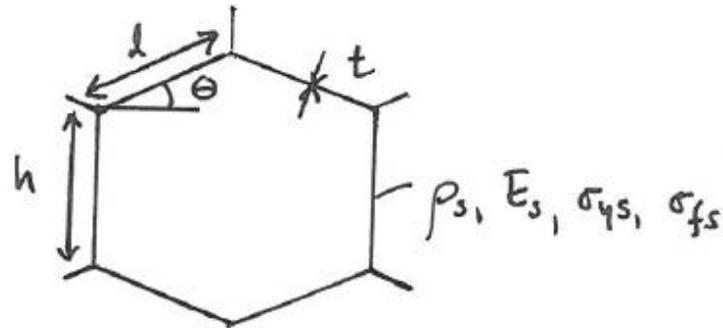
Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, © 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

## Tension

- Linear elastic – bending
- Stress plateau – exists only if cell walls yield
  - no buckling in tension
  - brittle honeycombs fracture in tension

## Variables affecting honeycomb properties

- Relative density  $\frac{\rho^*}{\rho_s} = \frac{\left(\frac{t}{l}\right) \left(\frac{h}{l} + 2\right)}{2 \cos \theta \left(\frac{h}{l} \sin \theta\right)} = \frac{2}{3} \frac{t}{l}$  regular hexagons
- Solid cell wall properties:  $\rho_s, E_s, \sigma_{ys}, \sigma_{fs}$
- Cell geometry:  $h/l, \theta$



## In-plane properties

Assumptions:

- $t/l$  small ( $(\rho_c^*/\rho_s)$  small)    neglect axial and shear contribution to deformation
- Deformations small    neglect changes in geometry
- Cell wall    linear elastic, isotropic

Symmetry

- Honeycombs are orthotropic    rotate  $180^\circ$  about each of three mutually perpendicular axes and structure is the same

## Linear elastic deformation

$$\begin{bmatrix} \epsilon_1 \\ \epsilon_2 \\ \epsilon_3 \\ \epsilon_4 \\ \epsilon_5 \\ \epsilon_6 \end{bmatrix} = \begin{bmatrix} 1/E_1 & -\nu_{21}/E_2 & -\nu_{31}/E_2 & 0 & 0 & 0 \\ -\nu_{12}/E_1 & 1/E_2 & -\nu_{32}/E_3 & 0 & 0 & 0 \\ -\nu_{13}/E_1 & -\nu_{23}/E_2 & -1/E_3 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1/G_{23} & 0 & 0 \\ 0 & 0 & 0 & 0 & 1/G_{13} & 0 \\ 0 & 0 & 0 & 0 & 0 & 1/G_{12} \end{bmatrix} \begin{bmatrix} \sigma_1 \\ \sigma_2 \\ \sigma_3 \\ \sigma_4 \\ \sigma_5 \\ \sigma_6 \end{bmatrix}$$

- Matrix notation:
 
$$\begin{aligned} \epsilon_1 &= \epsilon_{11} & \epsilon_4 &= \gamma_{23} & \sigma_1 &= \sigma_{11} & \sigma_4 &= \sigma_{23} \\ \epsilon_2 &= \epsilon_{22} & \epsilon_5 &= \gamma_{13} & \sigma_2 &= \sigma_{22} & \sigma_5 &= \sigma_{13} \\ \epsilon_3 &= \epsilon_{33} & \epsilon_6 &= \gamma_{12} & \sigma_3 &= \sigma_{33} & \sigma_6 &= \sigma_{12} \end{aligned}$$

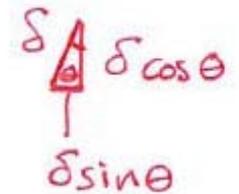
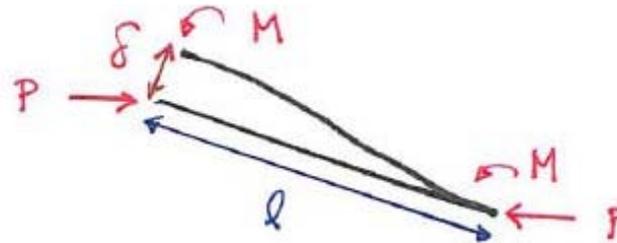
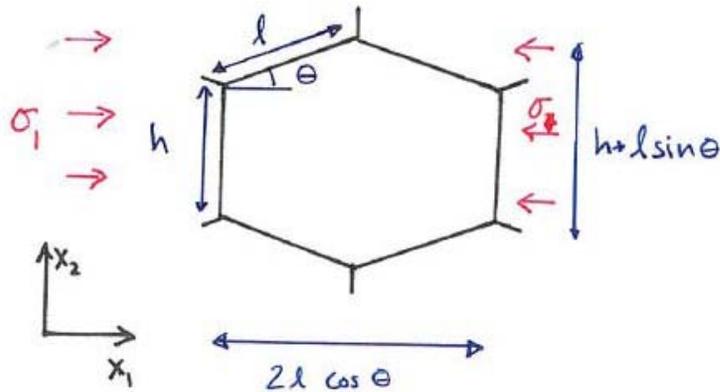
- In-plane ( $x_1 - x_2$ ): 4 independent elastic constants:

$$E_1 \quad E_2 \quad \nu_{12} \quad G_{12}$$

and compliance matrix symmetric  $\frac{-\nu_{12}}{E_1} = \frac{-\nu_{21}}{E_2}$  (reciprocal relation)

$$\left[ \text{notation for Poisson's ratio: } \nu_{ij} = \frac{-\epsilon_j}{\epsilon_i} \right]$$

### Young's modulus in $x_1$ direction



$$\sigma_1 = \frac{P}{(n + l \sin \theta) b}$$

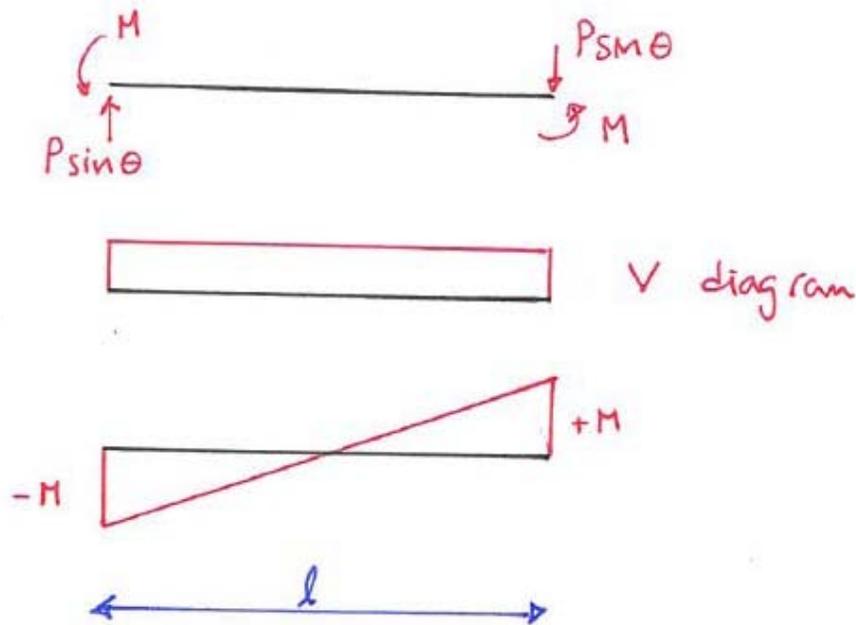
$$\epsilon_1 = \frac{\delta \sin \theta}{l \cos \theta}$$

Unit cell in  $x_1$  direction:  $2l \cos \theta$

Unit cell in  $x_2$  direction:  $h + 2l \sin \theta$

# In-Plane Deformation: Linear Elasticity

Figure removed due to copyright restrictions. See Figure 5: L. J. Gibson,  
M. F. Ashby, et al. "[The Mechanics of Two-Dimensional Cellular Materials.](#)"



M diagram: 2 cantilevers of length  $l/2$

$$\begin{aligned} \delta &= 2 \cdot \frac{P \sin \theta (l/2)^3}{3E_s I} \\ &= \frac{2 P l^3 \sin \theta}{24 E_s I} \\ \delta &= \frac{P l^3 \sin \theta}{12 E_s I} \quad I = \frac{b t^3}{12} \end{aligned}$$

Combining:

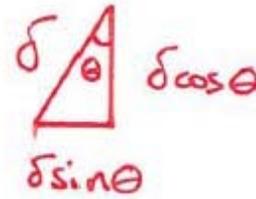
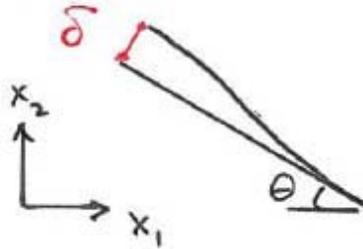
$$\begin{aligned} E_1^* &= \frac{\sigma_1}{\epsilon_1} = \frac{P}{(h + l \sin \theta) b} \frac{l \cos \theta}{\delta \sin \theta} \\ &= \frac{P}{(h + l \sin \theta) b} \frac{l \cos \theta}{P l^3 \sin^2 \theta} 12 E_s \frac{b t^3}{12} \\ E_1^* &= E_s \left( \frac{t}{l} \right)^3 \frac{\cos \theta}{(h/l + \sin \theta) \sin^2 \theta} = \frac{4}{3} \left( \frac{t}{l} \right)^3 E_s \end{aligned}$$

regular  
hexagons  
 $h/l=1 \theta = 30^\circ$

solid property    relative density    cell geometry

## Poisson's ratio for loading in $x_1$ direction

$$\nu_{12}^* = -\frac{\epsilon_2}{\epsilon_1}$$



$$\epsilon_1 = \frac{\delta \sin \theta}{l \cos \theta} \quad \epsilon_2 = \frac{\delta \cos \theta}{h + l \sin \theta} \quad (\text{lengthens})$$

$$\nu_{12}^* = \frac{\delta \cos \theta}{h + l \sin \theta} \left( \frac{l \cos \theta}{\delta \sin \theta} \right) = \frac{\cos^2 \theta}{(h/l + \sin \theta) \sin \theta}$$

- $\nu_{12}^*$  depends ONLY on cell geometry ( $h/l$ ,  $\theta$ ), not on  $E_s$ ,  $t/l$
- Regular hexagonal cells:  $\nu_{12}^* = 1$
- $\nu$  can be negative for  $\theta < 0$

e.g.  $h/l=2 \quad \theta = -30^\circ \quad \nu_{12}^* = \frac{3/4}{(3/2)(-1/2)} = -1$

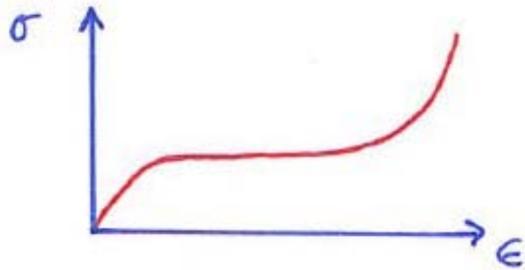
$\mathbf{E}_2^* \quad \nu_{12}^* \quad \mathbf{G}_{12}^*$

- Can be found in similar way; results in book

## Compressive strength (plateau stress)

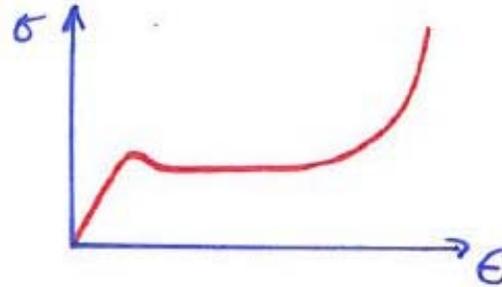
- Cell collapse by:

(1) elastic buckling



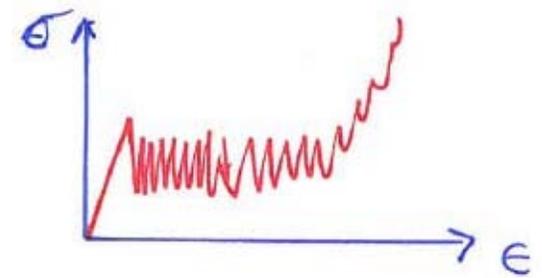
- buckling of vertical struts throughout honeycomb

(2) plastic yielding



- localization of yield
- as deformation progresses, propagation of failure band

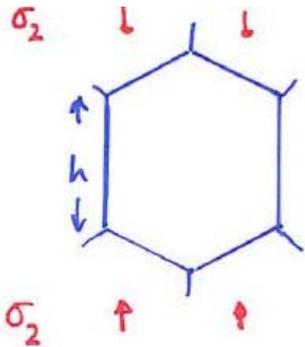
(3) brittle crushing



- peaks and valleys correspond to fracture of individual cell walls

### Plateau stress: elastic buckling, $\sigma_{el}^*$

- Elastomeric honeycombs cell collapse by elastic buckling of walls of length  $h$  when loaded in  $x_2$  direction
- No buckling for  $\sigma_1$ ; bending of inclined walls goes to densification



Euler buckling load

$$P_{cr} = \frac{n^2 \pi^2 E_s I}{h^2}$$

$n$ =end constraint factor



pin-pin  
 $n=1$



fixed-fixed  
 $n=2$

# Elastic Buckling

Figure removed due to copyright restrictions. See Figure 7: L. J. Gibson, M. F. Ashby, et al. "[The Mechanics of Two-Dimensional Cellular Materials](#)."

- Here, constraint  $n$  depends on stiffness of adjacent inclined members
- Can find elastic line analysis (see appendix if interested)
- Rotational stiffness at ends of column,  $h$ , matched to rotational stiffness of inclined members
- Find
 

$n/l=1$	1.5	2
$n=0.686$	0.760	0.860

$$\text{and } (\sigma_{el}^*)_2 = \frac{P_{cr}}{2l \cos \theta b} = \frac{n^2 \pi^2 E_s}{h^2 2l \cos \theta b} \frac{bt^3}{12}$$

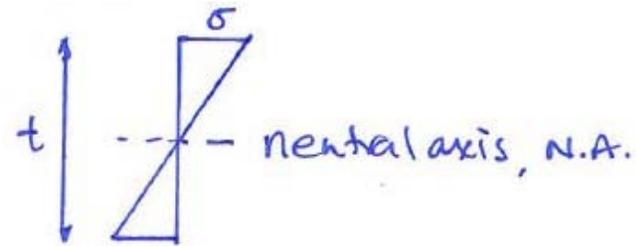
$$\boxed{(\sigma_{el}^*)_2 = \frac{n^2 \pi^2}{24} E_s \frac{(t/l)^3}{(h/l)^2 \cos \theta}}$$

regular hexagons:	$(\sigma_{el}^*)_2 = 0.22 E_s (t/l)^3$
and since	$E_2^* = 4/\sqrt{3} E_s (t/l)^3 = 2.31 E_s (t/l)^3$
strain at buckling	$(\epsilon_{el}^*)_2 = 0.10$ , for regular hexagons, independent of $E_s$ , $t/l$

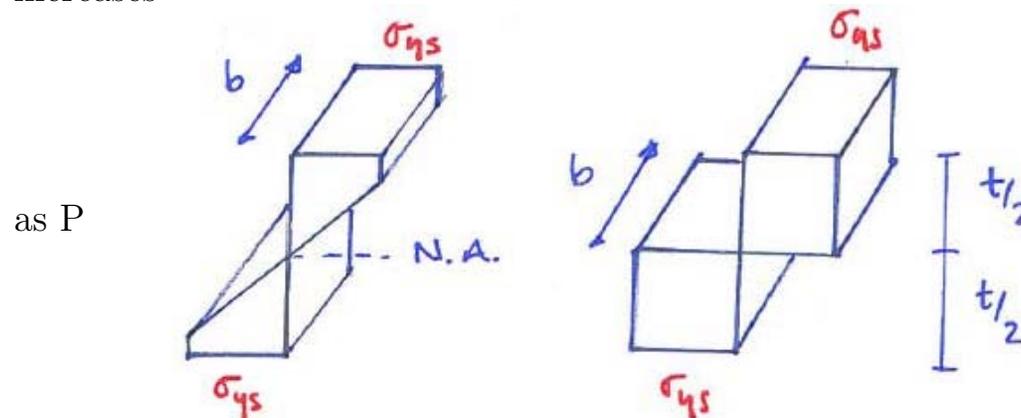
## Plateau stress: plastic yielding, $\sigma_{pl}^*$

- Failure by yielding in cell walls
- Yield strength of cell walls =  $\sigma_{ys}$
- Plastic hinge forms when cross-section fully yields

- Beam theory linear elastic  $\sigma = \frac{My}{I}$



- Once stress outer fiber =  $\sigma_{ys}$ , yielding begins and then progresses through the section, as the load increases



- When section fully yielded (right figure), form plastic hinge
- Section rotates like a pin

# Plastic Collapse

Figure removed due to copyright restrictions. See Figure 8: L. J. Gibson, M. F. Ashby, et al. "[The Mechanics of Two-Dimensional Cellular Materials](#)."

- Moment at formation of plastic hinge (plastic moment,  $M_p$ ):

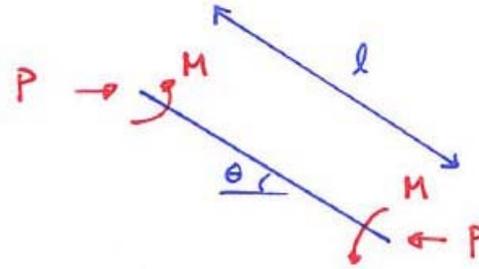
$$M_p = \left( \sigma_{ys} \frac{bt}{2} \right) \left( \frac{t}{2} \right) = \frac{\sigma_{ys} b t^2}{4}$$

- Applied moment, from applied stress

$$2M_{app} - PL \sin \theta = 0$$

$$M_{app} = \frac{Pl \sin \theta}{2}$$

$$\sigma_1 = \frac{P}{(h + l \sin \theta) b}$$



$$M_{app} = \sigma_1 (h + l \sin \theta) b \frac{l \sin \theta}{2}$$

- Plastic collapse of honeycomb at  $(\sigma_{pl}^*)_1$ , when  $M_{app} = M_p$

$$(\sigma_{pl}^*)_1 (h + l \sin \theta) b \frac{l \sin \theta}{2} = \sigma_{ys} \frac{bt^2}{4}$$

$$\boxed{(\sigma_{pl}^*)_1 = \sigma_{ys} \left( \frac{t}{l} \right)^2 \frac{1}{2(h/l + \sin \theta) \sin \theta}}$$

$$\text{regular hexagons: } (\sigma_{pl}^*)_1 = \frac{2}{3} \sigma_{ys} \left( \frac{t}{l} \right)^2$$

$$\text{similarly, } (\sigma_{pl}^*)_2 = \sigma_{ys} \left( \frac{t}{l} \right)^2 \frac{1}{2 \cos^2 \theta}$$

- For thin-walled honeycombs, elastic buckling can precede plastic collapse ( for  $\sigma_2$ )
- Elastic buckling stress = plastic collapse stress  $(\sigma_{el}^*)_2 = (\sigma_{pl}^*)_2$

$$\frac{n^2 \pi^2}{24} E_s \frac{(t/l)^3}{(h/l)^2 \cos \theta} = \frac{\sigma_{ys} (t/l)^2}{2 \cos^2 \theta}$$

$$(t/l)_{\text{critical}} = \frac{12 (h/l)^2}{n^2 \pi^2 \cos \theta} \left( \frac{\sigma_{ys}}{E_s} \right)$$

regular hexagons:  $(t/l)_{\text{critical}} = 3 \frac{\sigma_{ys}}{E_s}$

- E.g. metals  $\sigma_{ys}/E_s \sim .002$   $(t/l)_{\text{critical}} \sim 0.6\%$ 
  - most metal honeycomb denser than this
  - polymer  $\sigma_{ys}/E_s \sim 3 - 5\%$   $(t/l)_{\text{critical}} \sim 10-15\%$
  - low density polymers with yield point may buckle before yield

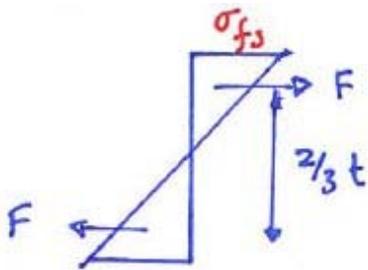
## Plastic stress: brittle crushing, $(\sigma_{cr}^*)_1$

- Ceramic honeycombs — fail in brittle manner
- Cell wall bending — stress reaches modulus of rupture — wall fracture loading in  $x_1$  direction:

$$P = \sigma_1 (h + l \sin \theta) b \quad \sigma_{fs} = \text{modulus of rupture of cell wall}$$

$$M_{\text{max. applied}} = \frac{P l \sin \theta}{2} = \frac{\sigma_1 (h + l \sin \theta) b l \sin \theta}{2}$$

## Moment at fracture, $M_f$

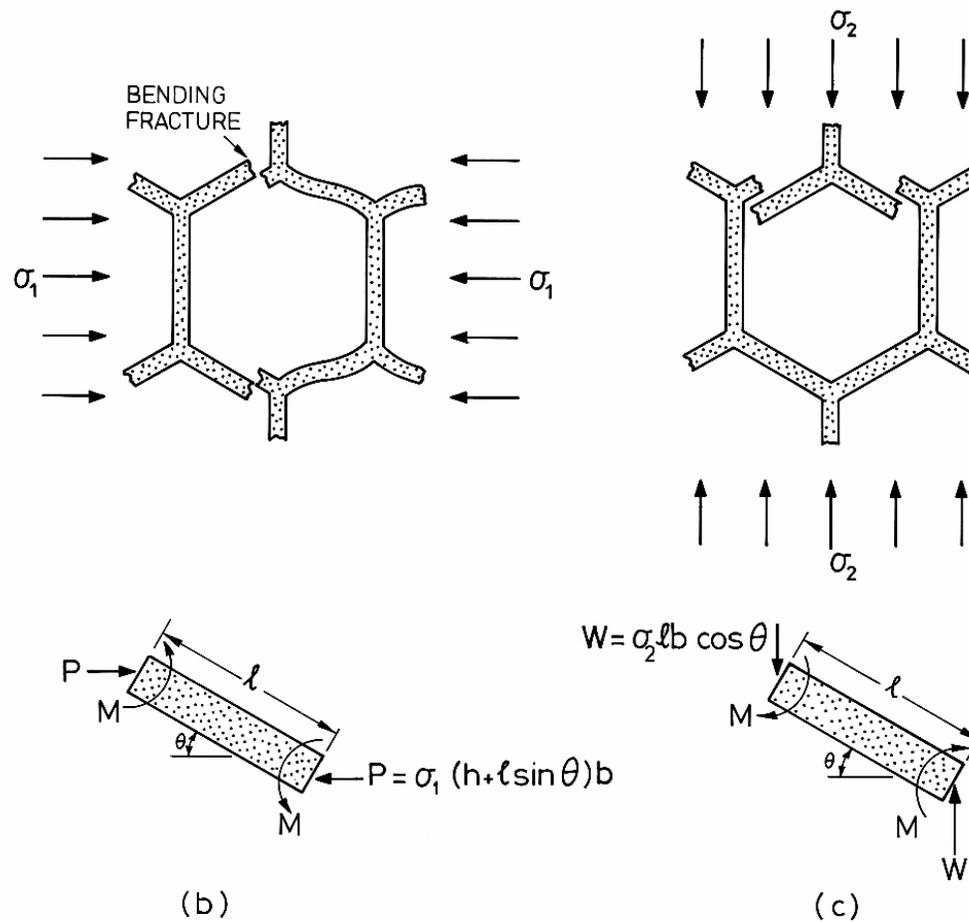


$$M_f = \left( \frac{1}{2} \sigma_{fs} b \frac{t}{2} \right) \left( \frac{2}{3} t \right) = \frac{\sigma_{fs} b t^2}{6}$$

$$(\sigma_{cr}^*)_1 = \sigma_{fs} \left( \frac{t}{l} \right)^2 \frac{1}{3 (h/l + \sin \theta) \sin \theta}$$

$$\text{regular hexagons: } (\sigma_{cr}^*)_1 = \frac{4}{9} \sigma_{fs} \left( \frac{t}{l} \right)^2$$

# Brittle Crushing



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, © 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

## Tension

- No elastic buckling
- Plastic plateau stress approx. same in tension and compression (small geometric difference due to deformation)
- Brittle honeycombs: fast fracture

## Fracture toughness

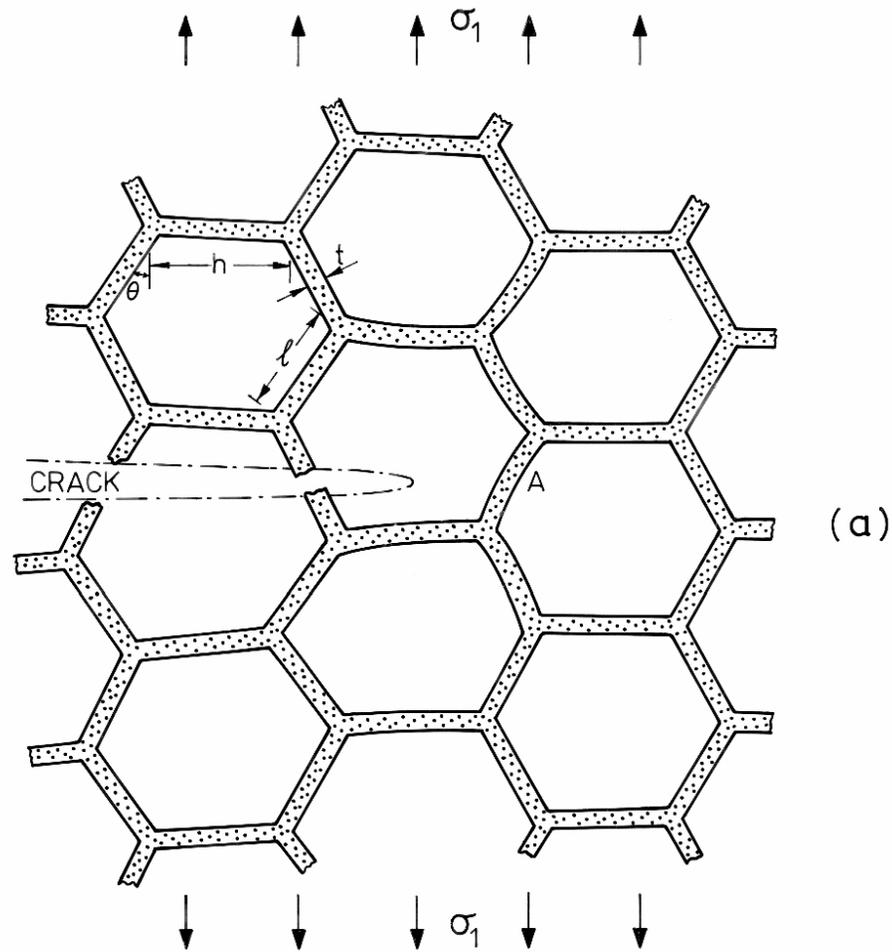
- Assume:
- crack length large relative to cell size (continuum assumption)
  - axial forces can be neglected
  - cell wall material has constant modulus of rupture,  $\sigma_{fs}$

Continuum: crack of length  $2c$  in a linear elastic solid material normal to a remote tension stress  $\sigma_1$  creates a local stress field at the crack tip



$$\sigma_{\text{local}} = \sigma_l = \frac{\sigma_1 \sqrt{\pi c}}{2\pi r}$$

# Fracture Toughness



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, © 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

Honeycomb: cell walls bent fail when applied moment = fracture moment

$M_{app} = P l$  on wall A

$$M_{app} = P l = \sigma_l l^2 b = \frac{\sigma_1 \bar{c} l^2 b}{\bar{l}} = \sigma_{fs} b t^2$$

$$(\sigma_f^*)_1 = \sigma_{fs} \left(\frac{t}{\bar{l}}\right)^2 \sqrt{\frac{\bar{l}}{c}}$$

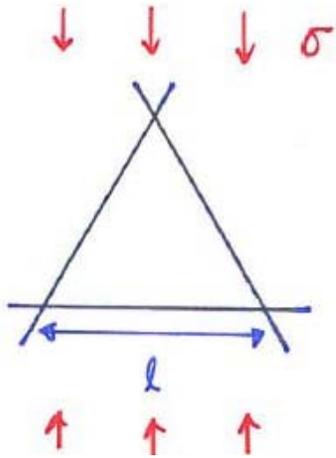
$$\boxed{K_{IC}^* = \sigma_f^* \bar{\pi} c = c \sigma_{fs} \left(\frac{t}{\bar{l}}\right)^2 \bar{l}} \quad \text{depends on cell size, } l!$$

c=constant

Summary: hexagonal honeycombs, in-plane properties

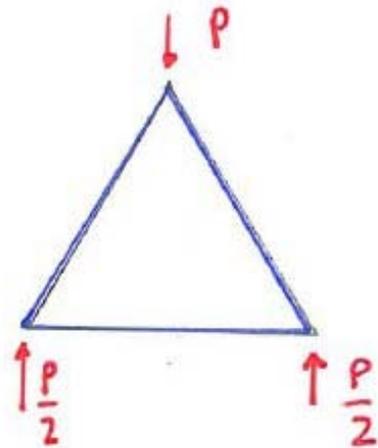
- Linear elastic moduli:  $E_1^* \quad E_2^* \quad \nu_{12}^* \quad G_{12}^*$
- Plateau stresses (compression)  $(\sigma_{el}^*)_2$  elastic buckling  
 $\sigma_{pl}^*$  plastic yield  
 $\sigma_{cr}^*$  brittle crushing
- Fracture toughness (tension)  $K_{IC}^*$  brittle fracture

## Honeycombs: In-plane behavior — triangular cells



depth  $b$  into page

- Triangulated structures - trusses
- Can analyze as pin-jointed (no moment at joints)
- Forces in members all axial (no bending)
- If joints fixed and include bending, difference  $\sim 2\%$
- Force in each member proportional to  $P$



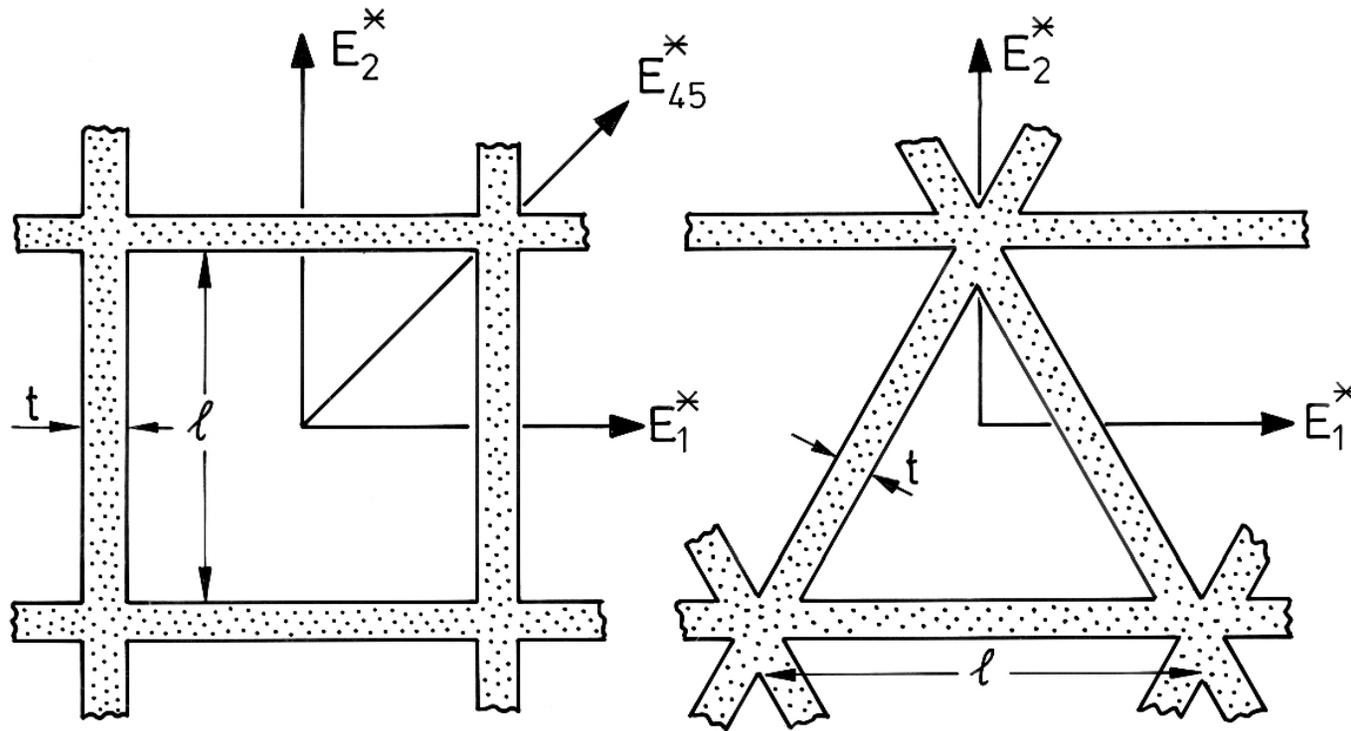
$$\sigma \propto \frac{P}{lb} \quad \epsilon \propto \frac{\delta}{l} \quad \delta \propto \frac{Pl}{AE_s} \text{---axial shortening: Hooke's law}$$

$$E^* \propto \frac{\sigma}{\epsilon} \propto \frac{P}{lb} \frac{l}{\delta} \propto \frac{P}{b} \frac{btE_s}{Pl} \propto E_s \left( \frac{t}{l} \right)$$

$$E^* = c E_s (t/l)$$

exact calculation:  $E^* = 1.15 E_s (t/l)$  for equilateral triangles

# Square and Triangular Honeycombs



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, © 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

MIT OpenCourseWare  
<http://ocw.mit.edu>

3.054 / 3.36 Cellular Solids: Structure, Properties and Applications  
Spring 2015

For information about citing these materials or our Terms of Use, visit: <http://ocw.mit.edu/terms>.