

Thermal properties of foams

- closed cell foams widely used for thermal insulation
- only materials with lower conductivity are aerogels (tend to be brittle + weak)
 - + vacuum insulation panels
- low thermal conductivity of foams arises from:
 - low volume fraction of solid
 - high volume fraction of gas with low λ
 - small cell size suppresses convection + radiation (through repeated absorption + reflection)
- applications: buildings, refrigerated vehicles, LNG tankers
- foams also have good thermal shock resistance since coeff. of thermal expansion of foam equal to that of the solid + modulus much lower ($\epsilon = \alpha \Delta T$ $\sigma = E \alpha \Delta T = \sigma_f$)
- ⇒ used as heat shields
- ceramic foams used as firebrick
 - ceramic has high T_m
 - foam - low λ - low heat loss
 - low heat capacity - lowers energy to heat furnace to temperature
 - good thermal shock resistance - resists spalling

Thermal conductivity, λ

- steady state conduction (T constant with time)

Fourier's law: $q = -\lambda \nabla T$

q = heat flux [$J/(m^2 s)$]

λ = thermal conductivity [W/mK]

∇T = temperature gradient

$$= i \frac{\partial T}{\partial x} + j \frac{\partial T}{\partial y} + k \frac{\partial T}{\partial z}$$

- non-steady heat conduction (T varies with time, t)

$$\frac{\partial T}{\partial t} = a \frac{\partial^2 T}{\partial x^2}$$

$$a = \text{thermal diffusivity} = \frac{\lambda}{\rho c_p}$$

ρ = density

c_p = specific heat = heat req'd to raise temp. of unit mass by 1°K

ρc_p = volumetric heat capacity [$J/m^3 K$]

- values for λ, a Table 7.1

Table 7.1 Thermal conductivities and diffusivities

Material	Thermal conductivity λ (W/m K)	Thermal diffusivity a (m ² /s)
Copper (solid)	384 ^a	8.8×10^{-5} a
Aluminium (solid)	230 ^a	8.9×10^{-5} a
Alumina (solid)	25.6 ^a	8.2×10^{-6} a
Glass (solid)	1.1 ^a	4.5×10^{-7} a
Polyethylene (solid)	0.35 ^a	1.7×10^{-7} a
Polyurethane (solid)	0.25 ^c	
Polystyrene (solid)	0.15 ^a	1.0×10^{-7} a
Air	0.025 ^a	—
Carbon dioxide	0.016 ^a	—
Trichlorofluoromethane (CCl ₃ F)	0.008 ^a	—
Oak ($\rho^*/\rho_s = 0.40$)	0.150 ^a	—
White pine ($\rho^*/\rho_s = 0.34$)	0.112 ^a	—
Balsa ($\rho^*/\rho_s = 0.09$)	0.055 ^a	—
Cork ($\rho^*/\rho_s = 0.14$)	0.045 ^a	—
Polystyrene foam ($\rho^*/\rho_s = 0.025$)	0.040 ^b	1.1×10^{-6} b
Polyurethane foam ($\rho^*/\rho_s = 0.02$)	0.025 ^b	9.0×10^{-7} b
Polystyrene foam ($\rho^*/\rho_s = 0.029\text{--}0.057$)	0.029–0.035 ^d	
Polyisocyanurate foam, (CFC-11) ($\rho^* = 32$ kg/m ³)	0.020 ^d	
Phenolic foam, (CFC-11, CFC-113) ($\rho^* = 48$ kg/m ³)	0.017 ^d	
Glass foam ($\rho^*/\rho_s = 0.05$)	0.050 ^d	
Glass wool ($\rho^*/\rho_s = 0.01$)	0.042 ^d	
Mineral fibre ($\rho^*/\rho_s = 4.8\text{--}32$ kg/m ³)	0.046 ^d	

All values for room temperature.

References

^aHandbook of Chemistry and Physics, 66th edn (1985–6) Chemical Rubber Co. ed. R. C. Weast.

^bPatten, G. A. and Skochdopole, R. E. (1962) *Mod. Plast.*, **39**, 149.

^cSchuetz, M. A. and Glicksman, L. R. (1983) *Proc. SPI 6th International Technical/Marketing Conference*, pp. 332–40.

^dGlicksman, L. R. (1994) Heat transfer in foams, in *Low Density Cellular Plastics* ed. Hilyard, N. C. and Cunningham, A. Chapman and Hall.

Data for thermal conductivity and thermal diffusivity

thermal diffusivity, α

- materials with a high value of α rapidly adjust their temp. to that of surroundings, because they conduct heat rapidly in comparison to their volumetric heat capacity; do not require much energy to reach thermal equilibrium

$$\text{e.g. Cu } \alpha = 112 \times 10^{-6} \text{ m}^2/\text{s}$$

$$\text{nylon } \alpha = 0.09 \times 10^{-6} \text{ m}^2/\text{s}$$

$$\text{wood } \alpha = 0.082 \times 10^{-6} \text{ m}^2/\text{s}$$

Thermal conductivity of a foam, λ^*

λ^* - contributions from - conduction through solid, λ_s^*

- " " " gas, λ_g^*

- convection within cells, λ_c^*

- radiation through cell walls + across voids, λ_r^*

$$\lambda^* = \lambda_s^* + \lambda_g^* + \lambda_c^* + \lambda_r^*$$

• conduction through solid: $\lambda_s^* = \eta \lambda_s (\rho^*/\rho_s)$ η = efficiency factor $\sim 2/3$

• conduction through gas: $\lambda_g^* = \lambda_g (1 - \rho^*/\rho_s)$

For example, 2.5% dense closed-cell polystyrene foam $\lambda^* = 0.040 \text{ W/mK}$

$$\lambda_s = 0.15 \text{ W/mK} \quad \lambda_g = 0.025 \text{ W/mK (air)}$$

$$\begin{aligned}\lambda_s^* + \lambda_g^* &= \frac{1}{3}(0.15)(0.025) + (0.025)(0.975) \\ &= 0.003 + 0.024 \\ &= 0.027 \text{ W/mK}\end{aligned}$$

- Most of conductivity comes from conduction through gas
- foams for insulation blown with low λ_g gases
- problem with aging - low λ_g gases diffuse out of foam over time, air diff. in, $\lambda_g^* \uparrow$

Convection within the cell



- gas rises + falls due to density changes with temperature
- density changes - buoyancy forces
- also have viscous forces from drag of gas as it moves past cell wall
- Convection important if Rayleigh number > 1000

$$Ra = \frac{\rho g \beta \Delta T_c l^3}{\mu a}$$

ρ = density of gas

g = grav. accn

β = volume expansion

$a_{\text{gas}} = \frac{1}{\mu} \tau / (\text{membrane})$ a = thermal diffusivity

ΔT_c = temp. diff. across cell

l = cell size

μ = dynamic viscosity of gas

Convection

$$\text{for } Ra = 1000 \quad \text{air} \quad p = p_{\text{atm}} \quad T = \text{room temp} \quad \beta = \gamma_T = 1/300 \text{ } (\text{K}^{-1})$$

$$\Delta T_c = 1^\circ\text{K} \quad \mu_{\text{air}} = 2 \times 10^{-5} \text{ Pa}\cdot\text{s} \quad \rho_{\text{air}} = 1.2 \text{ kg/m}^3$$

$$a_{\text{air}} = 2.0 \times 10^{-5} \text{ m}^2/\text{s}$$

$$\Rightarrow l = 20 \text{ mm}$$

- Convection important if cell size $> 20 \text{ mm}$
- Most foams: cell size $< 1 \text{ mm} \Rightarrow$ convection negligible

Radiation

- heat flux passing by radiation, q_r° , from a surface at temperature T_1 to one at a lower temp. T_0 , with a vacuum between them, is:

$$q_r^\circ = \beta_1 \sigma (T_1^4 - T_0^4) \quad \text{Stefan's law}$$

$$\sigma = \text{Stefan's constant} = 5.67 \times 10^{-8} \text{ W/m}^2\text{K}^4$$

β_1 = constant (<1) describing emissivity of the surface

(emitted radiant flux per unit area of sample relative to black body radiator at same temp. + conditions; black body absorbs all energy; black body emissivity = 1)

radiation

- If put foam between two surfaces, heat flux is reduced, since radiation is absorbed by the solid + reflected by cell walls

- Attenuation $q_r = q_r^* \exp(-k^* t^*)$ Beer's law
 - k^* = extinction coeff. for foam
 - t^* = thickness of foam

- for optically thin walls + struts ($t < 10\mu\text{m}$) (transparent to radiation)

$$k^* = (\rho^* / \rho_s) k_s$$

- heat flux by radiation then

$$q_r = \lambda_r^* \frac{dT}{dx}$$

$$q_r = \beta_1 \sigma (T_1^4 - T_0^4) \exp[-(\rho^*/\rho_s) k_s t^*] = \lambda_r^* \frac{dT}{dx}$$

- obtain λ_r^* using some approximations:

approximations :

$$\frac{dT}{dx} \approx \frac{T_1 - T_0}{t^*} = \frac{\Delta T}{t^*}$$

$$T_1^4 - T_0^4 = 4 \Delta T \bar{T}^3 \quad \bar{T} = \left(\frac{T_1 + T_0}{2} \right)$$

$$q_r = \beta_1 \sigma 4 \Delta T \bar{T}^3 \exp[-(\rho^*/\rho_s) k_s t^*] = \lambda_r^* \frac{\Delta T}{t^*}$$

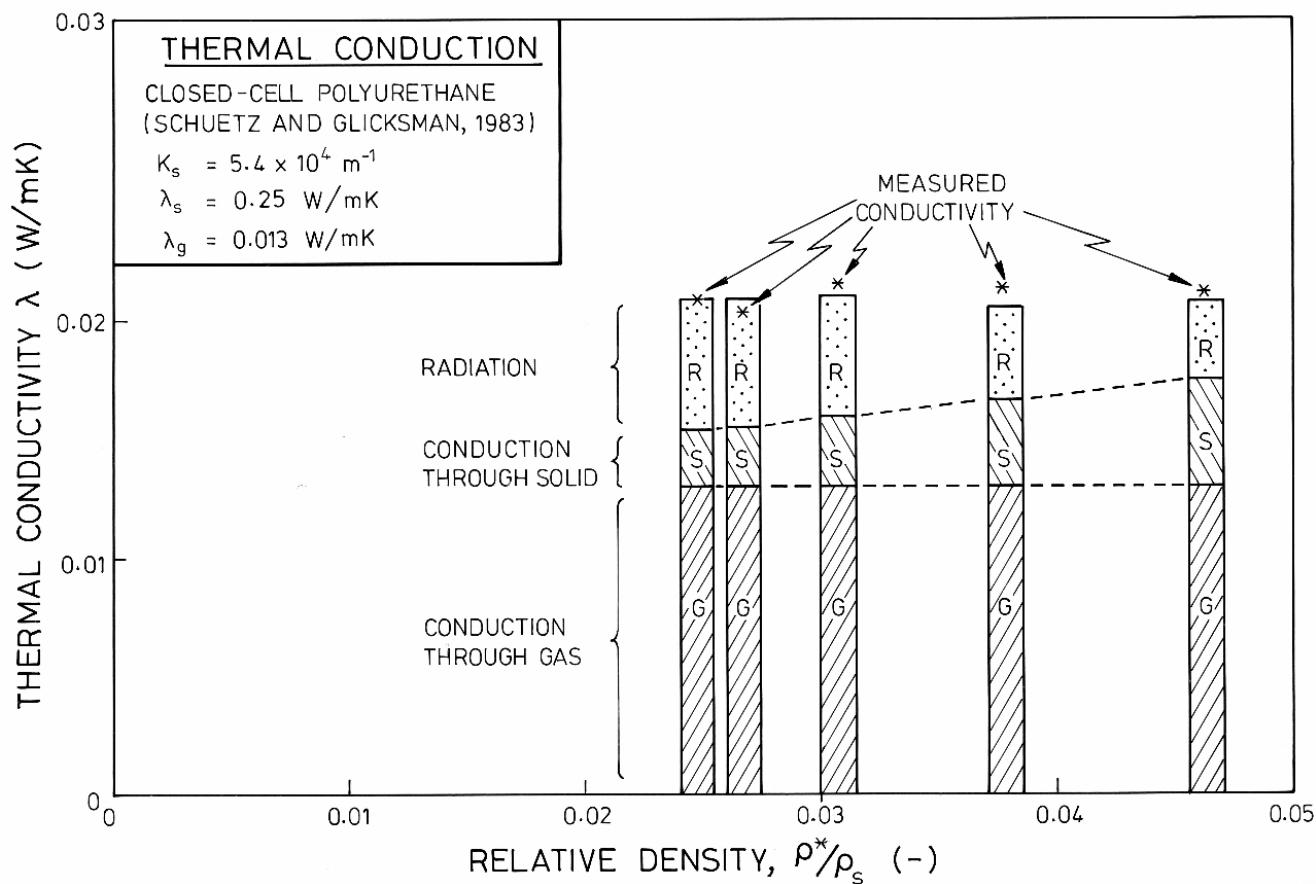
$$\lambda_r^* = 4 \beta_1 \sigma \bar{T}^3 t^* \exp[-(\rho^*/\rho_s) k_s t^*]$$

as $\rho^*/\rho_s \downarrow \lambda_r^* \uparrow$

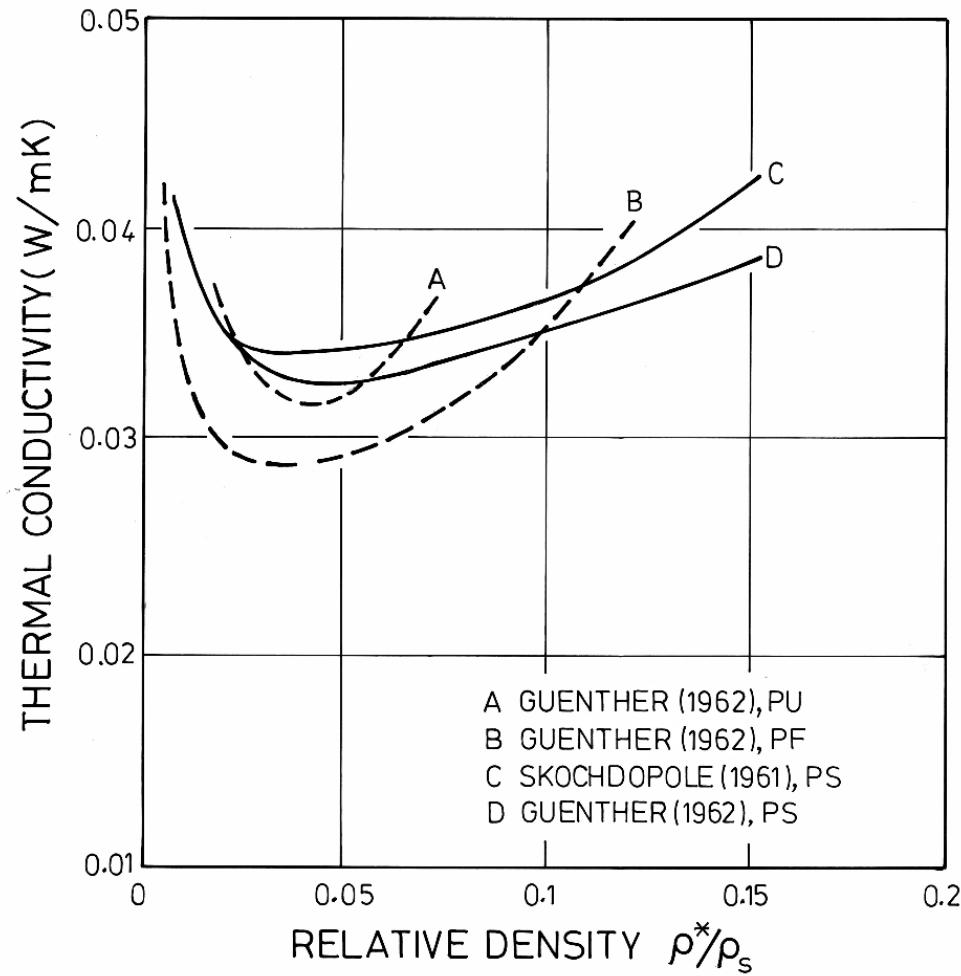
Thermal conductivity

- relative contributions of λ_s^* , λ_g^* , λ_r^* shown in fig 7.1
 - largest contribution λ_g^*
- λ^* plotted against relative density Fig 7.2
 - minimum @ between ρ^*/ρ_s of 0.03 & 0.07
 - at which point λ^* only slightly larger than λ_s^*
 - at low ρ^*/ρ_s , λ^* increases - increasing transparency to radiation
(also, walls may rupture)
 - tradeoff: as $\rho^*/\rho_s \downarrow$, $\lambda_s^* \downarrow$ but $\lambda_r^* \uparrow$

Thermal Conductivity

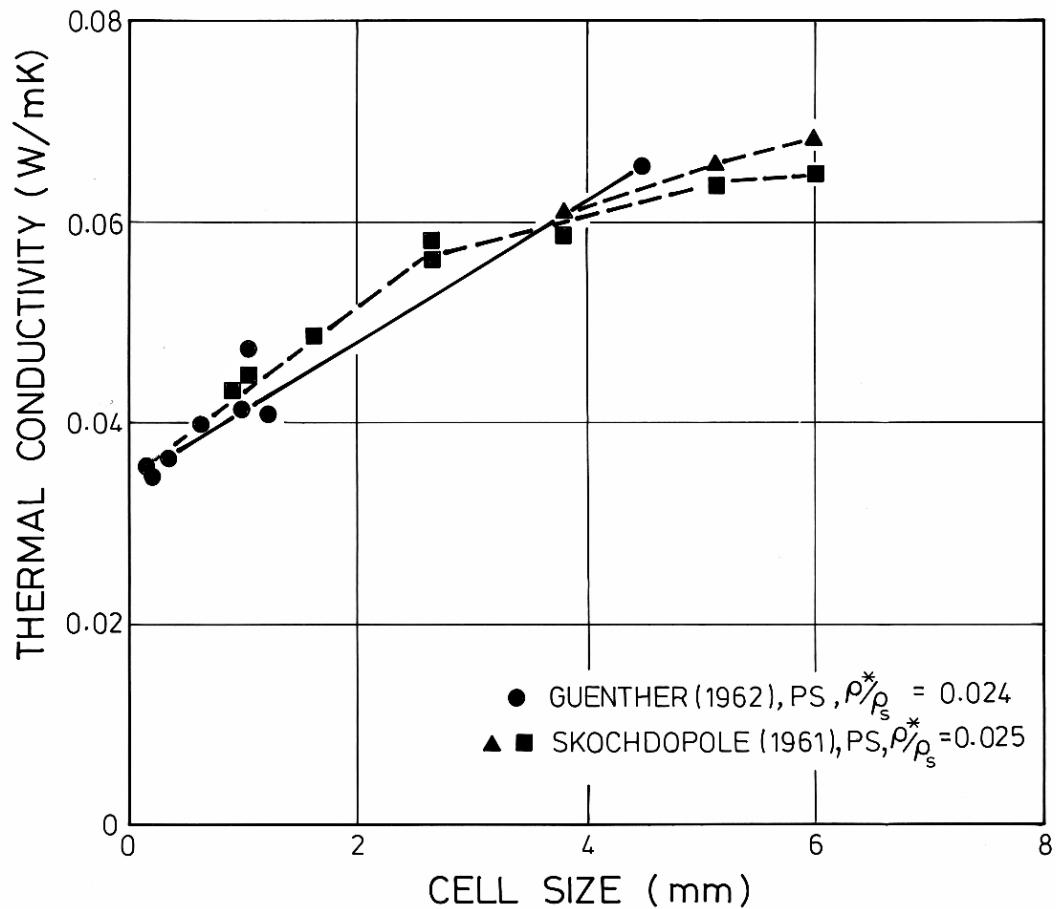


Cond. Vs. Relative Density



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, © 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

Cond. vs. Cell Size



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, © 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

λ^* plotted against cell size Fig 7.3

- λ^* increases with cell size
- radiation reflected less often

Note: aerogels

- pore size < 100 nm
- mean free path of air at ambient pressure = 68 nm
 ↳ avg. distance molecules move before collision with another molecule
- aerogels - pore size < mean free path of air - reduces conduction through gas.

Specific heat C_p

- specific heat = energy req'd to raise temp. of unit mass by unit temp
- $$C_p^* = C_{ps} \quad [J/kg \cdot K]$$

Thermal expansion coefficient

$$\alpha^* = \alpha_s \quad (\text{consider foam as framework})$$

(but if closed-cell foam cooled dramatically - gas can freeze, collapsing the cells; or if heated - gas expands, increasing the internal pressure + strains)

Thermal shock resistance

- If material subjected to sudden change in surface temp. - induces thermal stresses at surface + cracking + spalling.
- Consider material at T_1 dropped into water at T_2 ($T_1 > T_2$)
 - Surface temp. drops to T_2 , contracting surface layers
 - Thermal strain $\epsilon_T = \alpha \Delta T$
- If surface bonded to underlying block of material - constrained to original dimension

$$\sigma = \frac{E \alpha \Delta T}{1-\nu} \quad \text{in the surface}$$

- Cracking / spalling when $\sigma = \sigma_f$

$$\Delta T_c = \sigma_f \frac{(1-\nu)}{E\alpha} = \text{critical } \Delta T \text{ to just cause cracking}$$

- for foam: (open cells)

$$\Delta T_c^* = \frac{0.2 \sigma_{fs} (\rho^*/\rho_s)^{3/2} (1-\nu^*)}{E_s (\rho^*/\rho_s)^2 \alpha_s} = \frac{0.2}{(\rho^*/\rho_s)^{1/2}} \frac{\sigma_{fs} (1-\nu)}{E_s \alpha_s} = \frac{0.2}{(\rho^*/\rho_s)^{1/2}} \Delta T_{cs}$$

- As foam density \downarrow $\Delta T_c^* \uparrow$ firebrick - porous ceramic

Case study: optimization of foam density for thermal insulation

- there is an optimal foam density for a given thermal insulation problem
- already saw λ^* has a minimum as a $f(\rho^*/\rho_s)$
- typically, have a constraint on the foam thickness, t^* $t = \text{constant}$

$$\lambda^* = \frac{2}{3}(\rho^*/\rho_s)\lambda_s + (1-\rho^*/\rho_s)\lambda_g^* + 4\rho_s \beta_1 \sigma \bar{T}^3 t^* \exp[-k_s (\rho^*/\rho_s) t^*]$$

- What is optimum ρ^*/ρ_s for a given t^* ?

$$\frac{d\lambda^*}{d(\rho^*/\rho_s)} = 0 \Rightarrow (\rho^*/\rho_s)_{\text{opt}} = \frac{1}{k_s t^*} \ln \left[\frac{4 k_s \beta_1 \sigma \bar{T}^3 t^{*2}}{\frac{2}{3} \lambda_s - \lambda_g^*} \right]$$

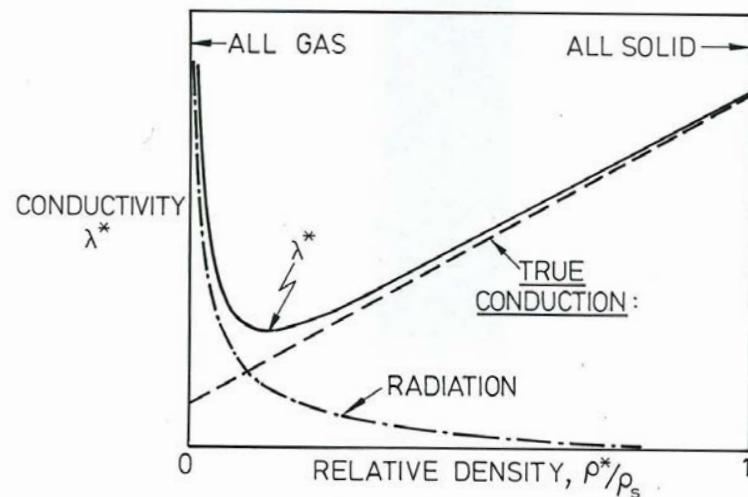
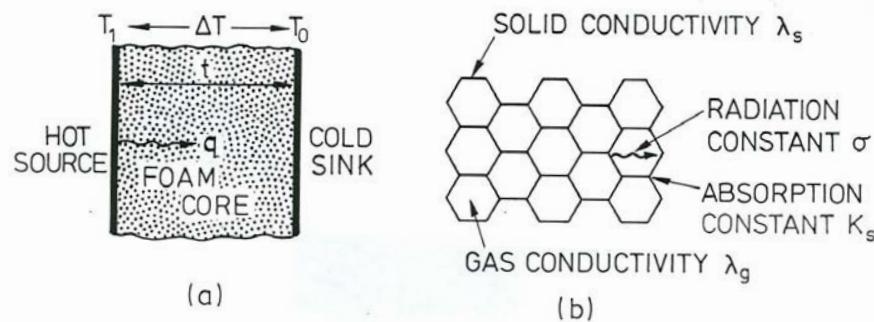
- as given thickness t^* increases, $(\rho^*/\rho_s)_{\text{opt}}$ decreases
- as \bar{T} increases, $(\rho^*/\rho_s)_{\text{opt}}$ increases

e.g. coffee cup $t^* = 3\text{mm}$ $(\rho^*/\rho_s)_{\text{opt}} = 0.08$

refrigerator $t^* = 50\text{mm}$ $(\rho^*/\rho_s)_{\text{opt}} = 0.02$

(see PP slide Table 7.3
for data used in
calculation).

Case Study: Optimization of Relative Density



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Case Study: Optimum Relative Density

Table 7.3 Data for optimization case study

Extinction coefficient of solid polymer, K_s	$5.67 \times 10^4 \text{ m}^{-1}$
Emissivity factor, β_l	0.5
Conductivity of solid polymer, λ_s	0.22 W/m K
Conductivity of gas, λ_g	0.02 W/m K
Mean temperature, \bar{T}	300°K
Stefan's constant, σ	$5.67 \times 10^{-8} \text{ W/m}^2\text{K}^4$

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Case study: insulation for refrigerators

- insulation reduces energy cost, but has a cost itself
- total cost = cost of insulation + cost of energy lost by heat transfer through walls
- objective function: minimize total cost
- given: x = thickness of insulation

C_M = cost of insulation/mass

ΔT = temp. diff. across insulation

C_E = " " energy/joule

t_d = design life of refrigerator

C_T = total cost/area

$$C_T = x \rho^* C_M + \lambda \frac{\Delta T}{x} t_d C_E \quad (\text{heat flux } q = \lambda \frac{\Delta T}{x} \frac{J}{m^2 s})$$

$$\text{define } M_1 = \frac{1}{\rho^* C_M} \quad M_2 = \frac{1}{\lambda}$$

$$\frac{C_T}{x} = \frac{1}{M_1} + \left[\frac{\Delta T}{x^2} t_d C_E \right] \frac{1}{M_2}$$

- two terms are equal when:

$$M_2 = \left[\frac{\Delta T}{x^2} t_e C_E \right] M_1$$

↳ coupling constant

- family of parallel straight lines of constant value $\frac{\Delta T}{x^2} t_e C_E$

- Fig 13.11 $\Delta T = 20^\circ$ $x = 10\text{mm}$ $C_E = 0.01/\text{mJ}$

two lines for $t_e = 10\text{ years}$ & $t_e = 1\text{ month}$

(note error in book $t_e = 10\text{ yrs}$ line should be mixed and)

- also plotted a set of curved contours - plots of C_T/x

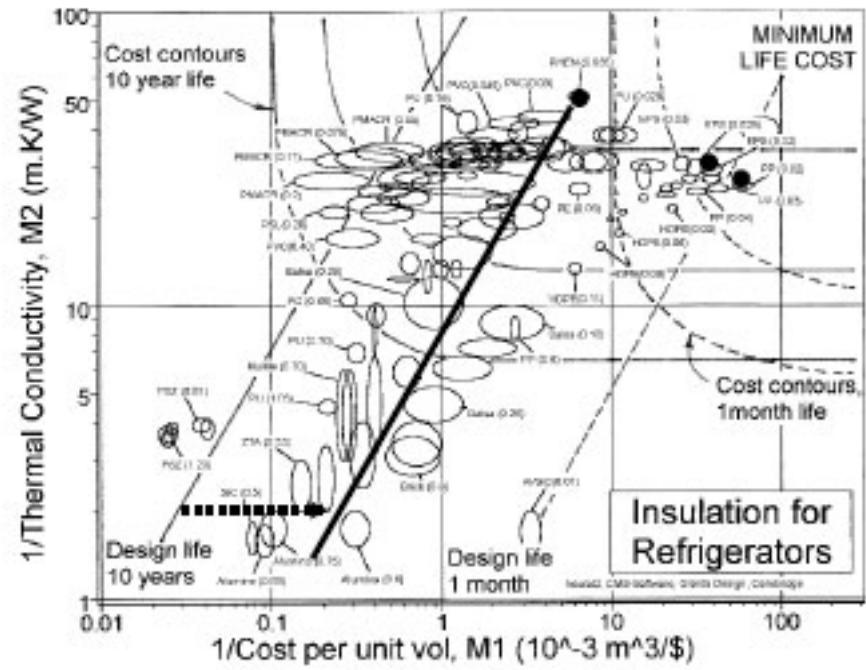
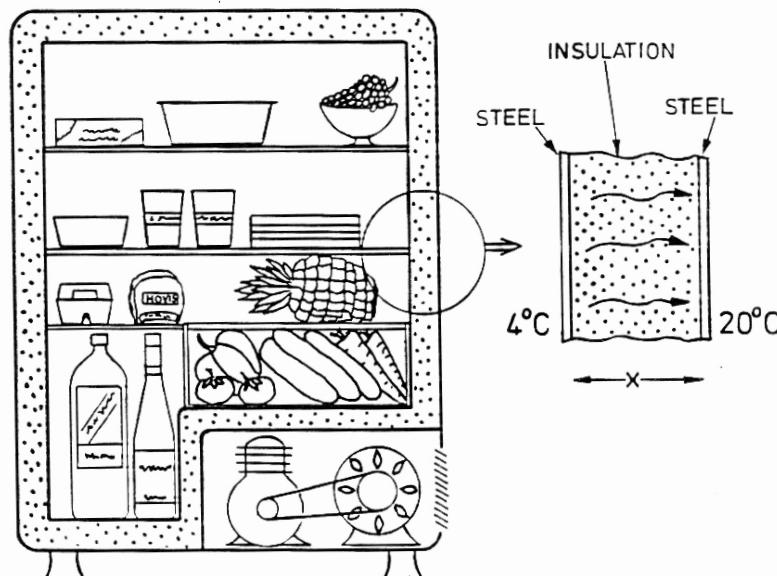
as move up + to right of plot, the value of $C_T/x \downarrow$

- for $t_e = 10\text{ years} \Rightarrow$ phenolic foam $\rho^* = 0.035 \text{ Mg/m}^3$

$t_e = 1\text{ month} \Rightarrow$ EPS $\rho^* = 0.02 \text{ Mg/m}^3$

PP $\rho^* = 0.02 \text{ Mg/m}^3$

Case Study: Insulation for Refrigerators



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