

3.024

**“Electronic, Optical, and Magnetic
Properties of Materials”**

3.024

The Final

- **The final exam will be focused on the second half of the class, but knowledge of basic quantum concepts and potentials is a must.**
- **There will be no complicated equations, so you won't need an equation sheet. Concepts come first.**
- **Bring your brain, pens, pencils and a calculator (strongly recommended).**

3.024 Topics Discussed

- Hamiltonian mechanics with application to normal vibrations in crystals
Phonons: dispersion relations, normal modes.
- Introduction to Quantum Mechanics: Schrodinger's Equation.
Applications to quantum dots, tunneling devices.
- Localized vs. delocalized states: from a free electron to an atom.
- Electronic states in crystals: DOS, bandgaps, interpretation of band diagrams.
- Fermions, symmetrization and Pauli's exclusion principle:
Electrons in bands and the classification of solids.
- "Free electron gas" description of carriers
- The chemical potential: Fermi level, statistics of electron distribution.
- Electronic structure of semiconductors: intrinsic and extrinsic.
- Semiconductor devices: p-n junctions under illumination and applied voltage.
- Maxwell's equations: electromagnetic waves in materials.
- Indices of refraction: reflection and transmission.
- Periodic optical materials: photonic bands and bandgaps.
- Magnetization in materials: para-, ferro-, anti-ferro and ferrimagnets.
- Magnetic domains.

Quantum Mechanical Potentials

Particle in Free space



Gravitational lensing in the Abel 370 galaxy cluster. (Photo courtesy of NASA, ESA, the Hubble SM4 ERO Team and ST-ECF <http://www.spacetelescope.org/images/heic0910b/>.)

Hamiltonian:

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2}$$

Energy eigenvalues:

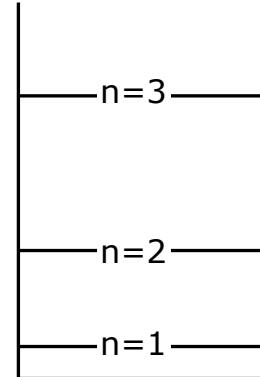
$$E_n = \frac{\hbar^2 k^2}{2m}$$

Energy eigenfunctions:

$$u_k(x) = e^{\pm ikx}$$

$$k = \frac{\sqrt{2mE}}{\hbar}$$

Particle in 1D box



Hamiltonian:

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x)$$

$$V(x) = \begin{cases} 0, & 0 \leq x \leq d \\ \infty, & x > d \text{ & } x < 0 \end{cases}$$

Energy eigenvalues:

$$E_n = \frac{\hbar^2 \pi^2 n^2}{2md^2}$$

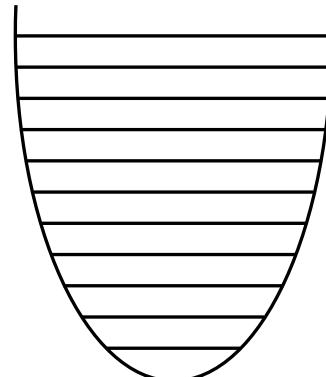
$$\Delta E_n = E_{n+1} - E_n = \frac{\hbar^2 \pi^2}{2md^2} (2n+1)$$

Energy eigenfunctions:

$$u_n(x) = A_n \sin \frac{n\pi x}{d}$$

$$A_n = \sqrt{\frac{2}{d}}$$

1D SHO



Hamiltonian:

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} m\omega^2 x^2$$

Energy eigenvalues:

$$E_n = \hbar\omega \left(n + \frac{1}{2} \right)$$

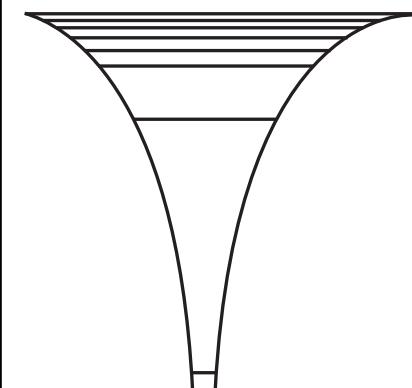
$$\Delta E_n = E_{n+1} - E_n = \hbar\omega$$

Energy eigenfunctions:

$$u_n(x) = e^{-\frac{m\omega}{2\hbar}x^2} h_n(x)$$

$h_n(x)$ – Hermite polynomial

Hydrogen atom



Hamiltonian:

$$\hat{H} = -\frac{\hbar^2}{2m_r} \nabla^2 - \frac{e^2}{r}$$

Energy eigenvalues:

$$E_n = -\frac{E_I}{n^2}$$

$$\Delta E_n = E_{n+1} - E_n = \frac{(2n+1)E_I}{n^2(n+1)^2}$$

Energy eigenfunctions:

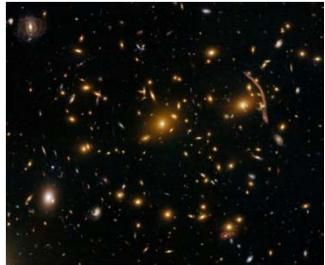
$$u_n(r) = R_{nl}(r) Y_l^m(\theta, \phi)$$

$$0 \leq l \leq n-1$$

$$-l+1 \leq m \leq l-1$$

Quantum Mechanical Potentials

Particle in Free space



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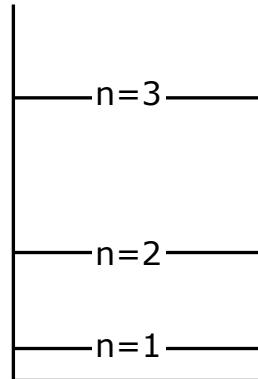
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Particle in 1D box



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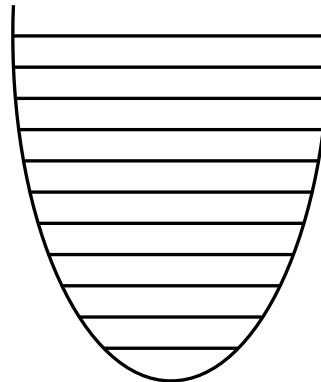
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Energy eigenfunctions:

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$$A_n = \sqrt{\frac{2}{d}}$$

1D SHO



Hamiltonian:

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} m\omega^2 x^2$$

Energy eigenvalues:

$$E_n = \hbar\omega \left(n + \frac{1}{2} \right)$$

$$\Delta E_n = E_{n+1} - E_n = \hbar\omega$$

Energy eigenfunctions:

$$u_n(x) = e^{-\frac{m\omega}{2\hbar}x^2} h_n(x)$$

$h_n(x)$ – Hermite polynomial

Periodic Potential

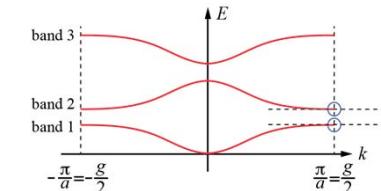
Figure removed due to copyright restrictions. See Fig. 3:Kittel, Charles. *Introduction to Solid State Physics*. 8th ed. Wiley 2004, p. 178.

Hamiltonian:

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \sum_G V_G e^{iGx}$$

$$G = ng, \quad g = \frac{2\pi}{a}$$

Energy eigenvalues:



Energy eigenfunctions:

$$u_{n,k}(x) = e^{ikx} \sum_G C_{k-G} e^{-iGx}$$

$$-\frac{g}{2} \leq x \leq \frac{g}{2} \Rightarrow BZ$$

Carriers in Intrinsic Semiconductors

D.O.S.:

$$g_c(E) = \frac{m_c^{3/2}}{\pi^2 \hbar^3} \sqrt{2(E - E_c)}$$

$$g_v(E) = \frac{m_v^{3/2}}{\pi^2 \hbar^3} \sqrt{2(E_v - E)}$$

Figure removed due to copyright restrictions. Fig. 2.16: Pierret, Robert F. *Semiconductor Fundamentals*. 2nd ed. Prentice Hall, 1988.

Carrier concentrations:

$$n = \int_{E_c}^{\infty} g_c(E) f(E) dE = N_c \exp\left(-\frac{E_c - E_F}{k_B T}\right)$$

$$p = \int_{-\infty}^E g_v(E) (1 - f(E)) dE = N_v \exp\left(-\frac{E_F - E_v}{k_B T}\right)$$

Fermi distribution function:

$$f_e(E) = \frac{1}{1 + e^{(E - E_F)/k_B T}}$$

$$f_h(E) = 1 - f_e(E) = \frac{1}{1 + e^{(E_F - E)/k_B T}}$$

Law of Mass Action:

$$np = N_c N_v \exp\left(-\frac{E_g}{k_B T}\right) = n_i^2, \quad n_i = \sqrt{N_c N_v} \exp\left(-\frac{E_g}{2k_B T}\right)$$

Figure removed due to copyright restrictions. Fig. 3: Kittel, Charles. *Introduction to Solid State Physics*. 8th ed. Wiley, 2004, p. 147.

Intrinsic carrier concentration:

$$n_i(\text{Si}) = 10^{10} \text{ cm}^{-3}$$

$$n_i(\text{GaAs}) = 2 \cdot 10^6 \text{ cm}^{-3}$$

Fermi level in Intrinsic SC:

$$n = p = n_i \Rightarrow \mu \equiv E_F = E_v + \frac{1}{2} E_g + \frac{3}{4} k_B T \ln\left(\frac{m_v^*}{m_c^*}\right)$$

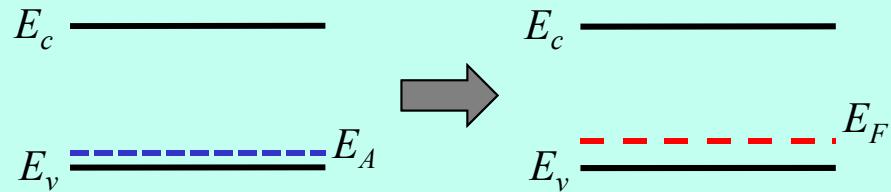
Extrinsic (Doped) Semiconductors

p-type

Figure removed due to copyright restrictions. Fig. 2.16: Pierret, Robert F. *Semiconductor Fundamentals*. 2nd ed. Prentice Hall, 1988.

$$p \sim N_A, \quad n \sim \frac{n_i^2}{N_A}$$

$$E_F = E_{Fi} - k_B T \ln \left(\frac{N_A}{n_i} \right)$$

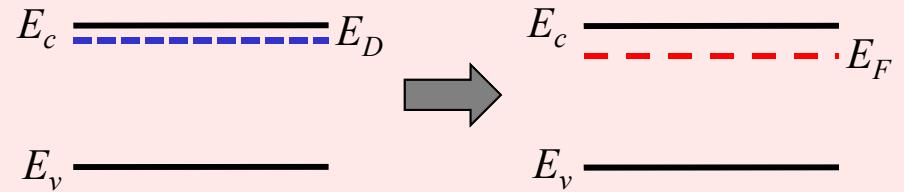


n-type

Figure removed due to copyright restrictions. Fig. 2.16: Pierret, Robert F. *Semiconductor Fundamentals*. 2nd ed. Prentice Hall, 1988.

$$n \sim N_D, \quad p \sim \frac{n_i^2}{N_D}$$

$$E_F = E_{Fi} + k_B T \ln \left(\frac{N_D}{n_i} \right)$$



Conductivity of SC: $\sigma = e n_c \mu_n + e p_v \mu_p$

Temperature dependence for *n*-type:

Figure removed due to copyright restrictions. Fig. 2.22: Pierret, Robert F. *Semiconductor Fundamentals*. 2nd ed. Prentice Hall, 1988.

PN Junction

W – depletion width
 V_{bi} – built-in voltage

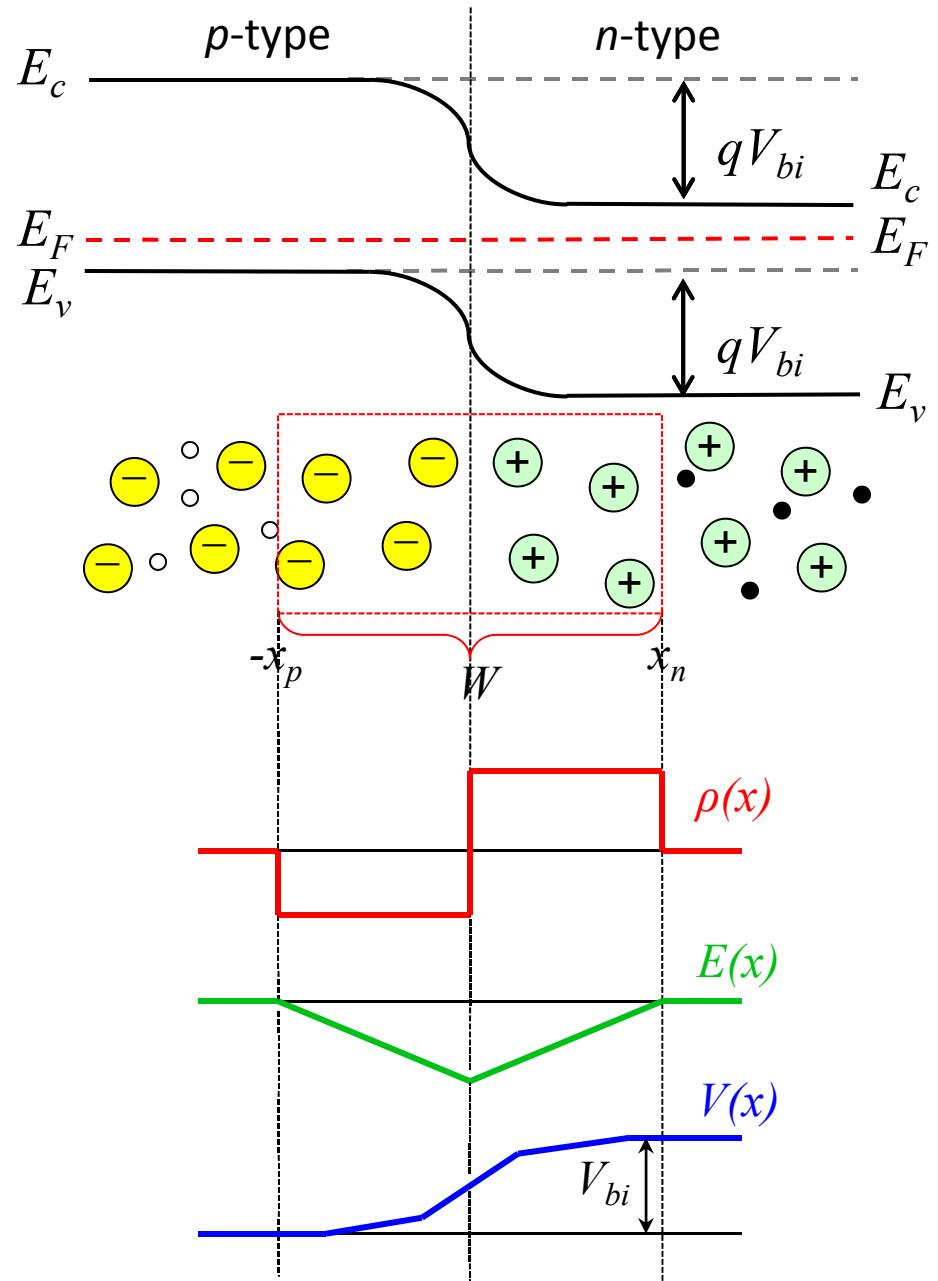
$$V_{bi} = E_F^n - E_F^p = \frac{k_B T}{e} \ln \left(\frac{N_D N_A}{n_i^2} \right)$$

$$N_D x_n = N_A x_p$$

$$x_p = \sqrt{\frac{2\epsilon_r \epsilon_0 V_{bi}}{e} \frac{N_D}{N_A(N_D + N_A)}}$$

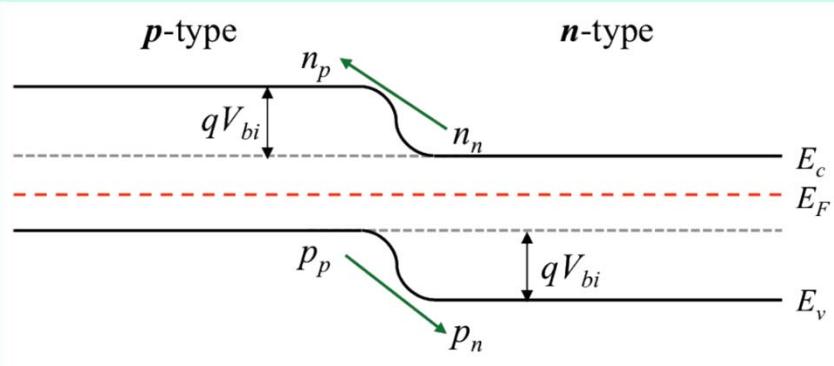
$$x_n = \sqrt{\frac{2\epsilon_r \epsilon_0 V_{bi}}{e} \frac{N_A}{N_D(N_D + N_A)}}$$

$$W = x_p + x_n = \sqrt{\frac{2\epsilon_r \epsilon_0 V_{bi}}{e} \frac{(N_D + N_A)}{N_D N_A}}$$



PN Junction Diodes: PVs and LEDs

Diode IV-characteristic



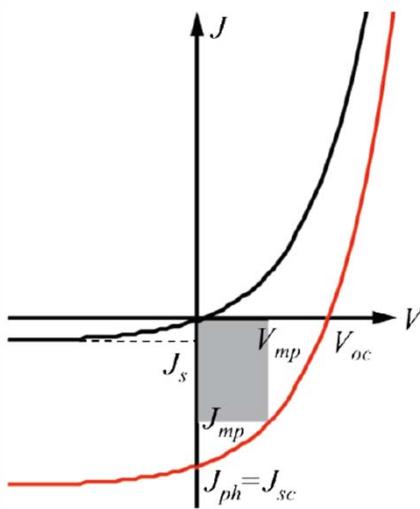
$$J = J_{diff} + J_{drift}$$

Forward bias: diffusion current
Reverse bias: drift current (J_s)

$$J = q \left(\frac{D_n}{L_n} \frac{n_i^2}{N_A} + \frac{D_p}{L_p} \frac{n_i^2}{N_D} \right) \left(e^{\frac{qV_a}{k_B T}} - 1 \right) = J_s \left(e^{\frac{qV_a}{k_B T}} - 1 \right)$$

Figure removed due to copyright restrictions. Fig. 15.1-18(c): Saleh, Bahaa E. A., and Malvin Carl Teich. *Fundamentals of Photonics*. 2nd ed. Wiley, 2007.

Photovoltaic cells:



$$J_{light} = J_s \left(e^{\frac{qV_a}{k_B T}} - 1 \right) - J_{ph}$$

$$V_{oc} = \frac{k_B T}{q} \ln \left(\frac{J_{ph}}{J_s} + 1 \right)$$

$$FF = \frac{V_{mp} J_{mp}}{V_{oc} J_{sc}}$$

Light Emitting Devices:

$$k_{sp}(\omega) = D \sqrt{\hbar\omega - E_g} e^{-(\hbar\omega - E_g)/k_B T}, \text{ where } D = \frac{(2m_r)^{\frac{3}{2}}}{2\pi^2 \hbar^2 \tau_R} e^{(E_{Fc} - E_{Fv} - E_g)/k_B T}$$

Figure removed due to copyright restrictions. Fig. 16.1-4: Saleh, Bahaa E. A., and Malvin Carl Teich. *Fundamentals of Photonics*. 2nd ed. Wiley, 2007.

$$\hbar\omega_{peak} = E_g + \frac{1}{2} k_B T$$

Electromagnetic Waves in Materials

Maxwell's Equations:

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\nabla \cdot \vec{D} = \rho_f$$

$$\nabla \times \vec{H} = \vec{J}_f + \frac{\partial \vec{D}}{\partial t}$$

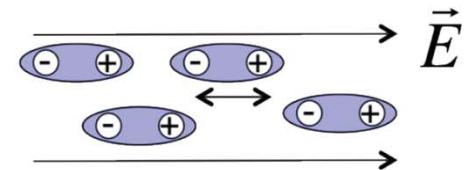
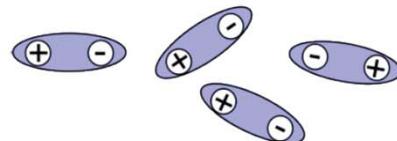
$$\nabla \cdot \vec{B} = 0$$

Constitutive Relations in Linear Media:

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P} = \epsilon_0 (1 + \chi) \vec{E} = \epsilon_r \epsilon_0 \vec{E} = \epsilon \cdot \vec{E}$$

$$\vec{B} = \mu_0 \vec{H} + \vec{M} = \mu_0 (1 + \chi_m) \vec{H} = \mu_r \mu_0 \vec{H} = \mu \cdot \vec{H}$$

Damped oscillator model for polarization:



$$\frac{d^2 \vec{P}}{dt^2} + \sigma \frac{d\vec{P}}{dt} + \omega_0^2 \vec{P} = \omega_0^2 \epsilon_0 \chi_0 \vec{E}$$

$$\vec{P} = \frac{\omega_0^2 \epsilon_0 \chi_0}{\omega_0^2 - \omega^2 - i\sigma\omega} \vec{E} = \epsilon_0 \chi(\omega) \vec{E}$$

$$\chi(\omega) = \chi'(\omega) + i\chi''(\omega)$$

$$\chi'(\omega) = \chi_0 \frac{\omega_0^2 (\omega_0^2 - \omega^2)}{(\omega_0^2 - \omega^2)^2 + (\sigma\omega)^2}$$

$$\chi''(\omega) = \chi_0 \frac{\omega_0^2 \sigma \omega}{(\omega_0^2 - \omega^2)^2 + (\sigma\omega)^2}$$

Wave Equations:

$$\left(\nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) \vec{E} = 0$$

$$\left(\nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) \vec{B} = 0$$

$$\vec{E}(\vec{r}, t) = \vec{E}_0 e^{i\vec{k}\vec{r} - i\omega t} + c.c.$$

$$\vec{H}(\vec{r}, t) = \vec{H}_0 e^{i\vec{k}\vec{r} - i\omega t} + c.c.$$

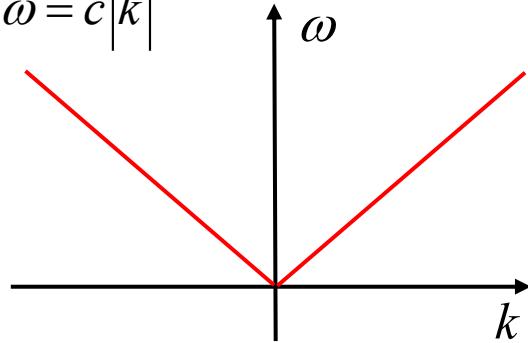
$$c = \frac{c_0}{n} = \frac{1}{\sqrt{\mu\epsilon}}, \quad c_0 = \frac{1}{\sqrt{\mu_0\epsilon_0}} \approx 3 \times 10^8 \frac{m}{s}$$

Figure removed due to copyright restrictions. See Fig. 5.5-5: Saleh, Bahaa E. A., and Malvin Carl Teich. *Fundamentals of Photonics*. 2nd ed. Wiley, 2007.

Optical Interfaces: Continuity of Phase

Dispersion relation for photons:

$$\omega = c|\vec{k}|$$



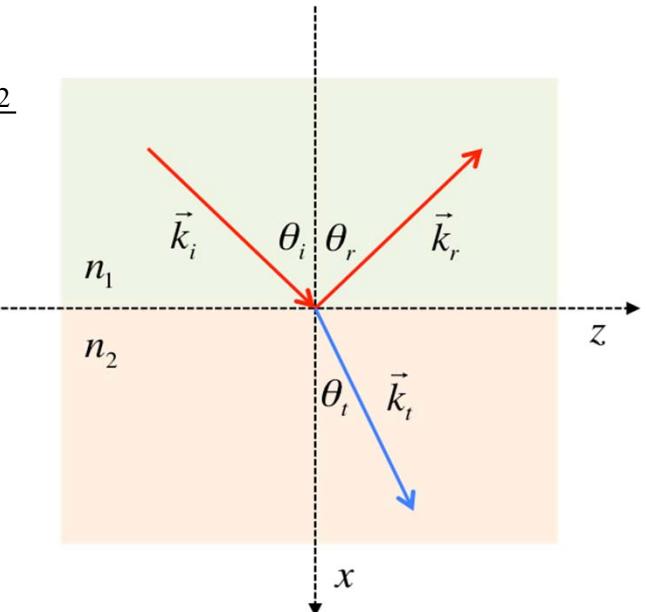
$$|\vec{k}_i| = |\vec{k}_r| = \frac{\omega n_1}{c_0}, \quad |\vec{k}_t| = \frac{\omega n_2}{c_0}$$

$$k_{iz} = k_{rz} = k_{tz} \equiv \beta$$

$$\theta_i = \theta_r$$

Snell's law:

$$n_1 \sin \theta_i = n_2 \sin \theta_t$$



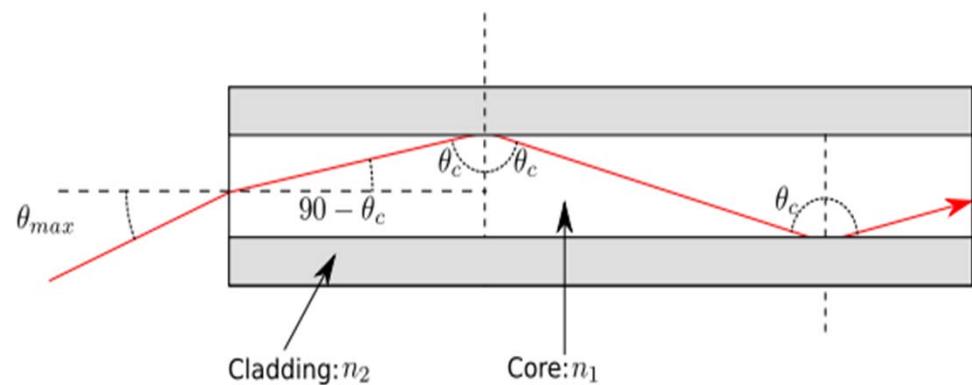
Total Internal Reflection:

$$\sin \theta_c = \frac{n_2}{n_1} \sin \theta_t = \frac{n_2}{n_1}$$

Numerical Aperture for a Waveguide:

$$NA = n \cdot \sin \theta_{\max} = \sqrt{n_1^2 - n_2^2}$$

For air: $n = 1$



Optical Interfaces: Boundary Conditions

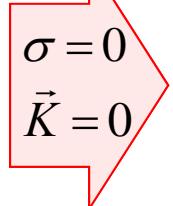
Boundary Conditions:

$$\vec{E}_{1\parallel} = \vec{E}_{2\parallel}$$

$$D_{2\perp} - D_{1\perp} = \sigma$$

$$\vec{H}_{2\parallel} - \vec{H}_{1\parallel} = \vec{K}$$

$$B_{1\perp} = B_{2\perp}$$



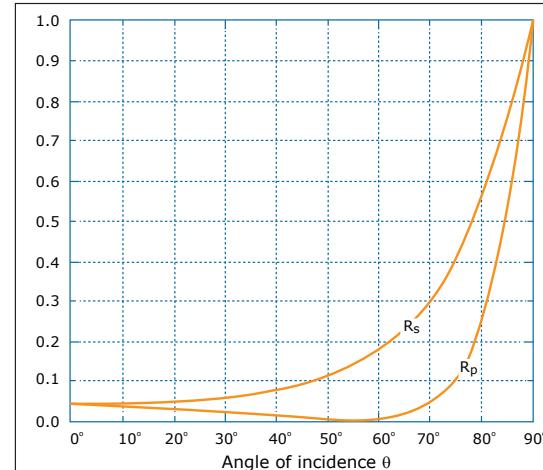
$$\vec{E}_{1\parallel} = \vec{E}_{2\parallel}$$

$$\epsilon_2 E_{2\perp} = \epsilon_1 E_{1\perp}$$

$$\vec{H}_{2\parallel} = \vec{H}_{1\parallel}$$

$$\mu_1 H_{1\perp} = \mu_2 H_{2\perp}$$

Reflection Coefficients:

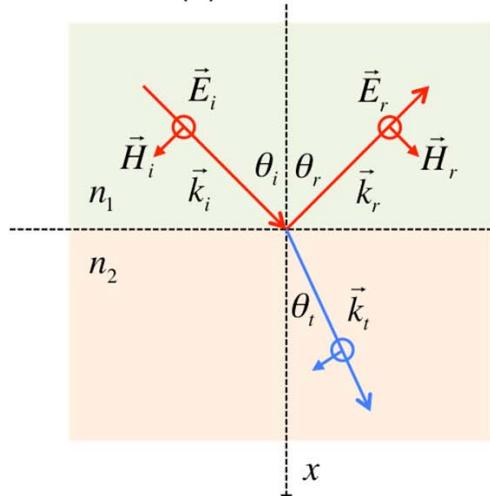


Brewster angle:

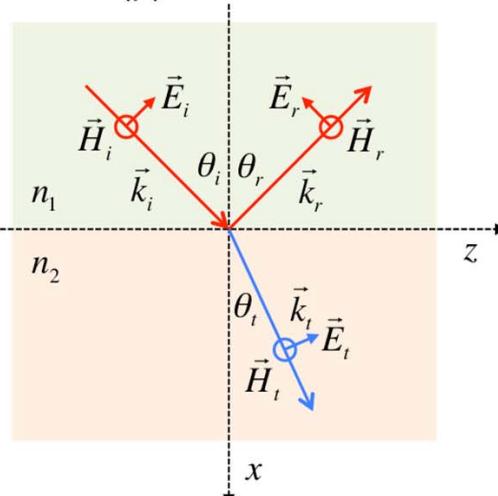
$$\sin^2 \theta_B = \frac{n_2^2}{n_1^2 + n_2^2}$$

Image by MIT OpenCourseWare.

(s) Polarization



(p) Polarization



Antireflective coatings:

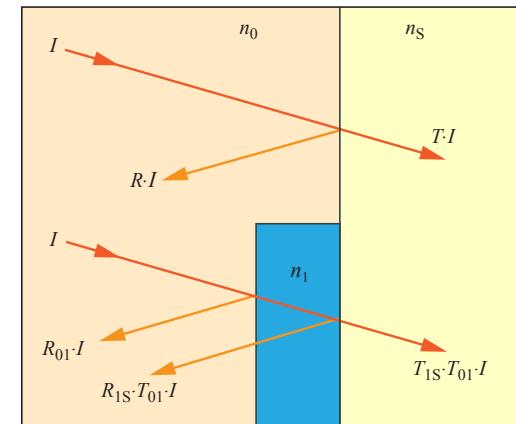
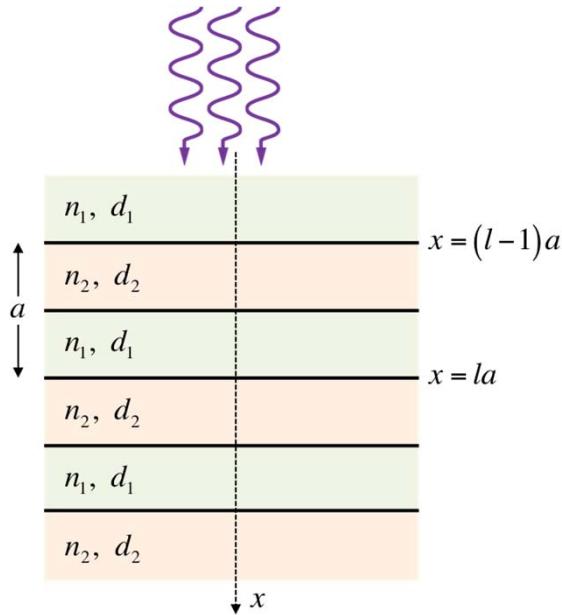


Image by MIT OpenCourseWare.

$$n_I = \sqrt{n_0 n_S} \quad d = \frac{\lambda_0}{4n_I}$$

Periodic Optical Materials: Photonic Crystals



$$\begin{pmatrix} E_{1(l-1)} \\ E'_{1(l-1)} \end{pmatrix} = \underbrace{D_{21}P_2D_{12}P_1}_M \begin{pmatrix} E_{1l} \\ E'_{1l} \end{pmatrix}$$

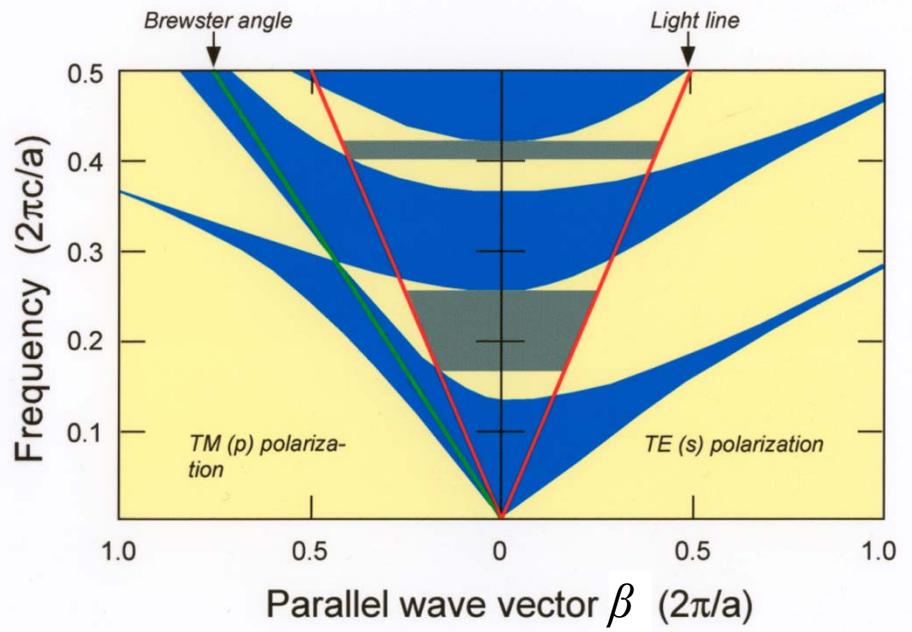
$$\begin{pmatrix} E_{1(l-1)} \\ E'_{1(l-1)} \end{pmatrix} = e^{-iKa} \begin{pmatrix} E_{1l} \\ E'_{1l} \end{pmatrix}$$

$$M \begin{pmatrix} E_{1l} \\ E'_{1l} \end{pmatrix} = e^{-iKa} \begin{pmatrix} E_{1l} \\ E'_{1l} \end{pmatrix}$$

$$k_z = \beta = 0$$

Figure removed due to copyright restrictions. Bloch waves corresponding to the A and B solutions for frequency at the edge of the Brillouin zone: Unknown source.

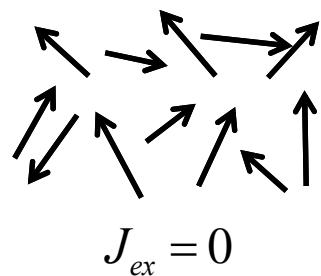
$$g = \frac{2\pi}{a}, \quad a = d_1 + d_2$$



Classification of Magnetic Materials

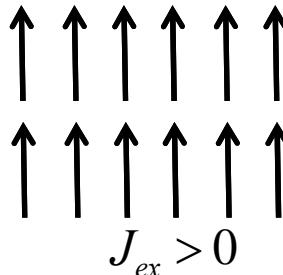
$$\text{Exchange interaction: } E_{ex} = -2J_{ex}\vec{S}_1\vec{S}_2$$

Paramagnetic



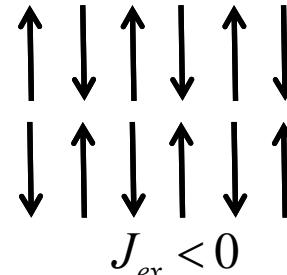
$$J_{ex} = 0$$

Ferromagnetic



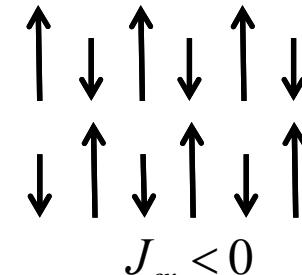
$$J_{ex} > 0$$

Anti-Ferromagnetic



$$J_{ex} < 0$$

Ferrimagnetic



$$J_{ex} < 0$$

Magneto crystalline anisotropy energy:

$$E_a = \sum_n K_{un} \sin^{2n} \theta \approx K_u \sin^2 \theta$$

Easy axis: low magnetic field needed to magnetize to saturation

Hard axis: high magnetic field needed to magnetize to saturation

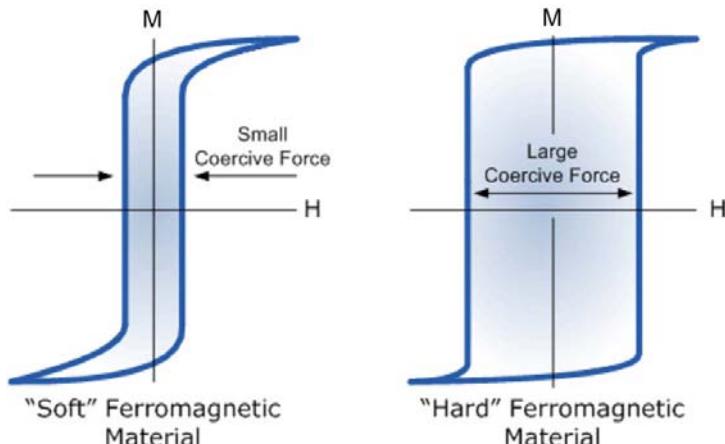
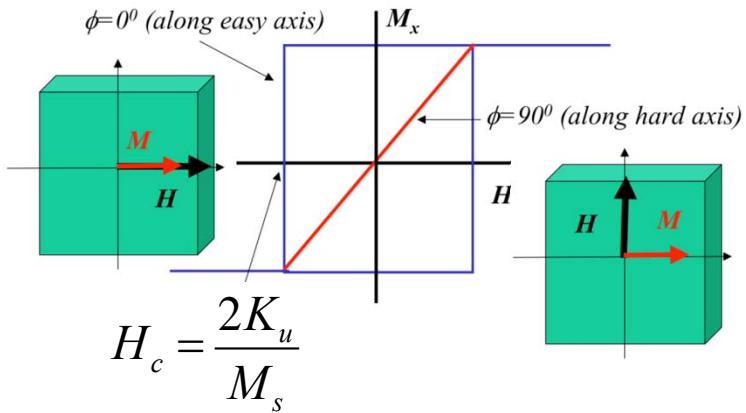
Mn, Fe, Ni – cubic, low anisotropy

Co – hexagonal, high anisotropy

Figure removed due to copyright restrictions. See Fig. 6.1: O' Handley, Robert C. *Modern Magnetic Materials*. Wiley, 1999.

Hysteresis in Ferromagnetic Materials

Single domain ferromagnet:



Courtesy of [Wayne Storr](#). Used with permission.

Soft: low anisotropy, easy to magnetize
(transformers, generators)

Hard: high anisotropy, hard to magnetize
(hard drives, permanent magnets)

Domains and domain walls

Domain wall width: $\delta = Na = \pi S \sqrt{\frac{2J}{K_u a^3}} \Rightarrow \begin{cases} J \uparrow \Rightarrow \delta \uparrow \\ K_u \uparrow \Rightarrow \delta \downarrow \end{cases}$

Domain wall energy: $\sigma_{BW} = \pi S \sqrt{\frac{2JK_u}{a}}$

Types of walls:

Domain wall motion:

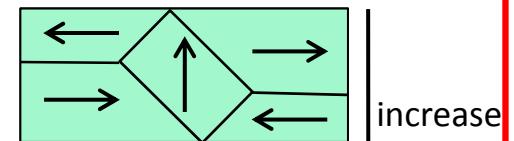
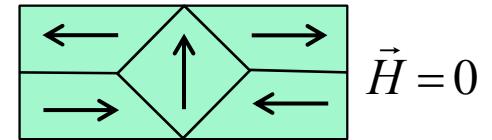
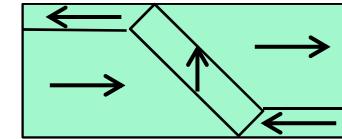


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Bloch wall and Neél wall: O'Handley, Robert C. *Modern Magnetic Materials*. Wiley, 1999.



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3.024 Electronic, Optical and Magnetic Properties of Materials

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