Lecture 14: 10.31.05 Batteries continued; Thermodynamic stability

Today:

THERMODYNAMICS OF A BATTERY, CONTINUED	
IMPLICATIONS OF STABILITY REQUIREMENTS FOR THE PROPERTIES OF MATERIALS	
Le Chatelier's principle	
Constraints on thermodynamic properties ensure stability: example of internal energy constraints	
REQUIREMENTS FOR THE SHAPE OF FREE ENERGY CURVES	
MAXWELL RELATIONS	.(
The Euler Relationship and Maxwell Relations	(
References	

Reading: C.H.P. Lupis, Chemical Thermodynamics of Materials, 'Stability of one-

component systems,' pp. 34-35.

H.B. Callen, Thermodynamics and an Introduction to Thermostatistics, 'Stability

of Thermodynamic Systems,' pp. 202-212

Engel and Reid 6.2 (Maxwell relations)

Supplementary Reading:

Thermodynamics of a battery, continued

- Last time, we began the analysis of how a battery converts chemical work into electrical work (transport of
 electrons). This system introduces a new form of internal energy, arising from the transport of charged
 species, which is incorporated into the fundamental equation for a system as the electrochemical potential:
 - The script F in this expression is the Faraday constant: $F = 96,485 \frac{C}{mole\ e^-}$
- We applied this new term for electrochemical potentials in our analysis of the Daniell cell.

Half-cell reactions:

$$Zn^{I} \Leftrightarrow (Zn^{2+})^{II} + 2(e^{-})^{I}$$

 $(Cu^{2+})^{III} + 2(e^{-})^{IV} \Leftrightarrow Cu^{IV}$

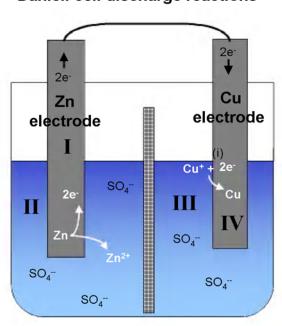
Total reaction for system:

$$Zn^{I} + \left(Cu^{2+}\right)^{III} \Leftrightarrow \left(Zn^{2+}\right)^{II} + Cu^{IV}$$

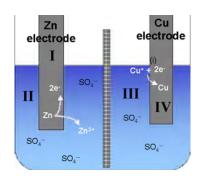
$$\Delta \overline{G}_{rxn} = \mu_{Cu}^{IV} + \widetilde{\mu}_{Zn^{++}}^{II} - \mu_{Zn}^{I} - \widetilde{\mu}_{Cu^{++}}^{III}$$

 We showed last time that the electrical potential difference obtained across the electrodes I and IV is related to the free energy of reaction by the Nernst Equation:

Daniell cell discharge reactions



 How can we further simplify this expression? Recall that we have sulfate ion equilibration through the membrane between the aqueous solutions II and III:



$$\begin{pmatrix} \tilde{\mu}_{SO_4^{2-}} \end{pmatrix}^{II} = \left(\tilde{\mu}_{SO_4^{2-}} \right)^{III} \\
\mu_{SO_4^{2-}}^{II} - 2\phi^{II} F = \mu_{SO_4^{2-}}^{III} - 2\phi^{III} F$$
(c) $\phi^{II} = \phi^{III}$

- For a reversible reaction, we call the potential difference the electromotive force, or EMF:
- The U.S. convention is that EMF is positive when reactions are written to proceed spontaneously from left to right (i.e. Zn is the left electrode).

An example calculation: Thermodynamics of a fuel cell

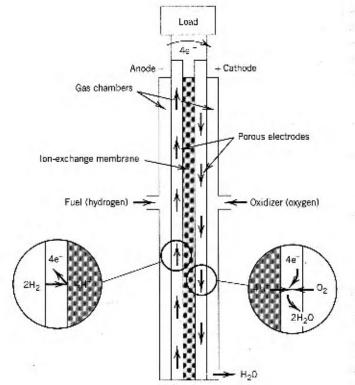
 Consider again a fuel cell similar to the one shown schematically in the last lecture notes. A more simplified schematic is shown below:

(Silbey et al.)

The Half-cell reactions and overall reaction are:

$$\begin{split} \frac{1}{2}O_{2(g)} + 2H^{+} + 2e^{-} &\longleftrightarrow H_{2}O_{(I)} \\ \frac{H_{2(g)} &\longleftrightarrow 2H^{+} + 2e^{-}}{\frac{1}{2}O_{2(g)} + H_{2(g)} &\longleftrightarrow H_{2}O_{(I)} \\ & \text{E}^{\circ} = \text{1.2288 V} \end{split}$$

If a certain fuel cell runs with P_{O2} = P_{H2} = 5000 psi (340 atm), what voltage is achieved by the cell? Assume the gases can be treated as ideal.

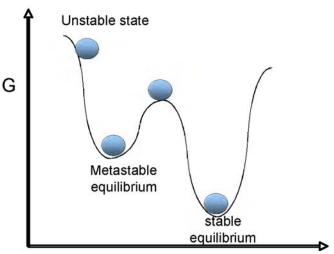


Implications of stability requirements for the properties of materials

Le Chatelier's principle

Materials are most useful in stable equilibrium states- conditions where the material is resistant to change
if the system experiences small fluctuations in its state. For example, if the temperature and pressure are
held constant at (T₀,P₀), then stable equilibrium is identified by finding the minimum value of G in the
landscape of possible U, S, and V values taken by the system. Using a figure from our first lecture:

$$G = H - TS = U - T_0S + P_0V$$

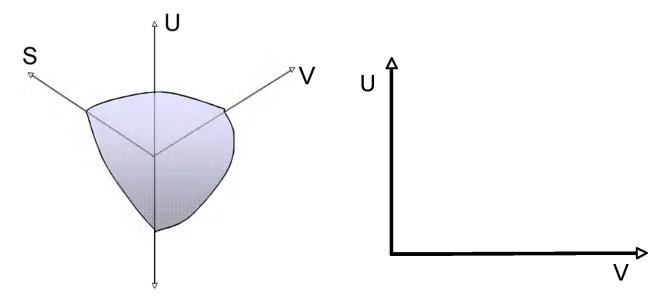


Extensive variable of the system (U, V, É)

• **Le Chatelier's principle** states that a system perturbed by a small fluctuation will feel a thermodynamic driving force to return to the stable equilibrium state.

Constraints on thermodynamic properties ensure stability: example of internal energy constraints

- MAIN IDEA: Many of the thermodynamic properties of materials can only have certain values if the material is in a stable equilibrium state.
- Recall the equilibrium condition in terms of internal energy: For conditions of fixed total entropy, the internal energy of the system is minimized. Graphically, this means:



• The surface drawn in the diagram represents all of the possible equilibrium states of some system. The mathematical requirements for a stable equilibrium are:

For equilibrium to be stable against fluctuations in S only or V only:

For equilibrium to be stable against coupled fluctuations in both S and V:

• These requirements on the shape of the internal energy surface are linked to values of thermodynamic parameters of the system:

- This is the requirement for thermal stability. Since the absolute temperature must be > 0 the
 heat capacity must also be greater than 0. In addition, the entropy must increase if the
 temperature increases, for stability.
- · Next, we have:

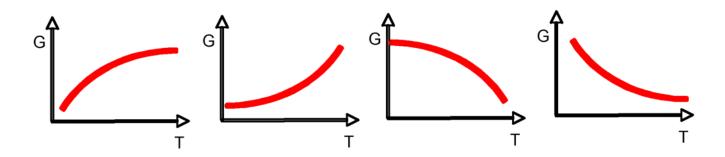
- o This is the condition for mechanical stability.
- From this analysis, we see that many properties of materials may only have certain values if the material is to be stable. There are numerous other examples that can be derived using the other thermodynamic functions. In summary, one can prove:

$$C_P \ge C_V \ge 0$$

$$\kappa_T \geq \kappa_S \geq 0$$

Requirements for the shape of free energy curves

- What does the free energy as a function of typical experimental parameters look like- for example, what
 does a plot of G vs. T look like? What about G vs. P? Do we know anything (qualitatively) about how it
 must look? The answer is yes: the relationships between the Gibbs free energy and other
 thermodynamic parameters tell us numerous things about how plots of G vs. various thermodynamic
 variables must behave in stable thermodynamic systems:
 - o Consider a plot of G vs. temperature:



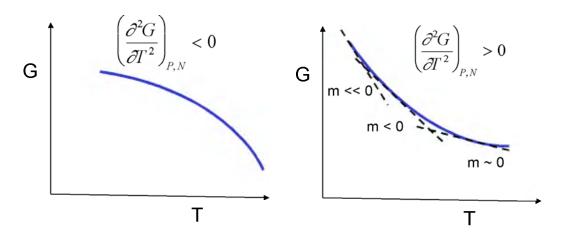
- o What should the slope of this curve be?
 - Using our differential expression for G:

$$dG = VdP - SdT + \sum_{j=1}^{P} \sum_{i=1}^{C} \mu_i^j dn_i^j$$

... and combining this with the algebraic definition of the differential for G(T,P,N):

$$dG = \left(\frac{\partial G}{\partial T}\right)_{P,N} dT + \left(\frac{\partial G}{\partial P}\right)_{T,N} dP + \sum_{j=1}^{P} \sum_{i=1}^{C} \left(\frac{\partial G}{\partial n_{i}^{j}}\right)_{T,N_{j\neq i}} dn_{i}^{j}$$

- ...thus, the slope of the curve is -S.
- o what is the curvature of G vs. T?
 - The curvature is given by the second derivative of G vs. T- it is the rate of change of the slope:



o The curvature is also related to known thermodynamic quantities:

- The heat capacity and the absolute temperature must always be ≥ 0, thus the curvature must be < 0. Thus plots of the free energy vs. temperature at constant pressure must always have the general shape shown at left above.</p>
- Similar analyses may be applied to predict the shape of other free energy curves of interest.

Maxwell Relations

The Euler Relationship and Maxwell Relations

• We've already discussed some of the useful properties of state functions like U, S, and H- they are path independent and integrable. An additional useful characteristic is that state functions must obey the Euler reciprocal relationship, which is expressed mathematically for a multivariable function f(x,y) as:

$$\frac{\partial^2 f(x,y)}{\partial x \partial y} = \frac{\partial^2 f(x,y)}{\partial y dx}$$

EULER RELATIONSHIP

- o The Euler relationship can be used to identify identities between thermodynamic variables that are not obvious: these are called the *Maxwell relations*.
- A set of Maxwell relations can be derived for each thermodynamic state function. For example, starting with the internal energy *U*(*S*, *V*, *N*):

- Maxwell's relations can help us further derive thermodynamic quantities from measurable parameters of our materials.
- Another example is the Maxwell relation obtained by taking second derivatives of the enthalpy *H*(*S*,*P*,*N*):

• Summarizing these relationships for each of our main thermodynamic functions:

$$\begin{array}{ll}
\mathsf{U} & -\left(\frac{\partial P}{\partial S}\right)_{V,n} = \left(\frac{\partial T}{\partial V}\right)_{S,n} \\
H \equiv U + PV & \left(\frac{\partial V}{\partial S}\right)_{P,n} = \left(\frac{\partial I}{\partial P}\right)_{S,n} \\
G \equiv H - TS & -\left(\frac{\partial S}{\partial P}\right)_{T,n} = \left(\frac{\partial V}{\partial T}\right)_{P,n} \\
F \equiv U - TS & \left(\frac{\partial S}{\partial V}\right)_{T,n} = \left(\frac{\partial P}{\partial T}\right)_{V,n}
\end{array}$$

References

Reference Type: Book Record Number: 34

Author: Silbey, R.J.; Alberty, R.A.; Bawendi, M.G.

Year: 2005

Title: Physical Chemistry

City: New York

Publisher: John Wiley Number of Pages: 944

Edition: 4th

Reference Type: Electronic Source

Record Number: 1 Author: Carter, W.C.

Year: 2002

Title: 3.00 Thermodynamics of Materials Lecture Notes

URL: http://pruffle.mit.edu/3.00/