

# Particle Systems and ODE Solvers II, Mass-Spring Modeling

With slides from Jaakko Lehtinen  
and others

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# ODEs and Numerical Integration

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$$\frac{d \mathbf{X}(t)}{dt} = f(\mathbf{X}(t), t)$$

- Given a function  $f(\mathbf{X}, t)$  compute  $\mathbf{X}(t)$
- Typically, *initial value problems*:
  - Given values  $\mathbf{X}(t_0) = \mathbf{X}_0$
  - Find values  $\mathbf{X}(t)$  for  $t > t_0$
- We can use lots of standard tools

# Reduction to 1st Order

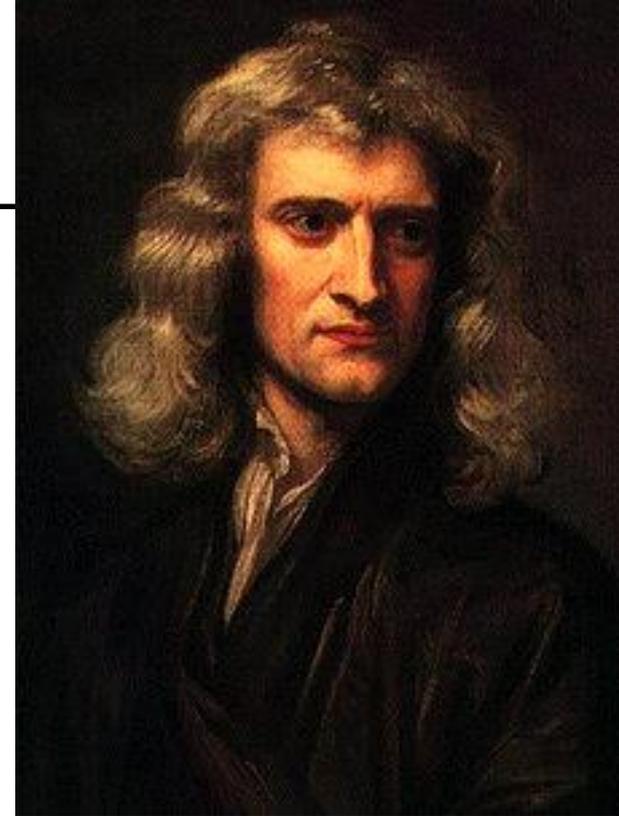
- Point mass: 2<sup>nd</sup> order ODE

$$\vec{F} = m\vec{a} \quad \text{or} \quad \vec{F} = m \frac{d^2 \vec{x}}{dt^2}$$

- Corresponds to system of first order ODEs

$$\begin{cases} \frac{d}{dt} \vec{x} = \vec{v} \\ \frac{d}{dt} \vec{v} = \vec{F} / m \end{cases}$$

2 unknowns ( $\mathbf{x}$ ,  $\mathbf{v}$ )  
instead of just  $\mathbf{x}$



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# ODE: Path Through a Vector Field

- $\mathbf{X}(t)$ : path in multidimensional phase space

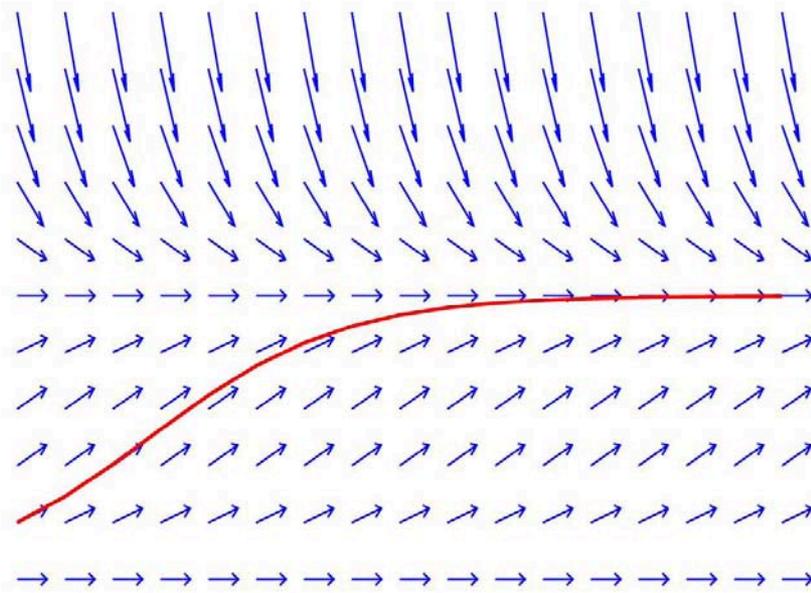


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$$\frac{d}{dt} \mathbf{X} = f(\mathbf{X}, t)$$

“When we are at state  $\mathbf{X}$  at time  $t$ , where will  $\mathbf{X}$  be after an infinitely small time interval  $dt$ ?”

- $f = d/dt \mathbf{X}$  is a vector that sits at each point in phase space, pointing the direction.

# Euler, Visually

$$\frac{d}{dt}\mathbf{X} = f(\mathbf{X}, t)$$

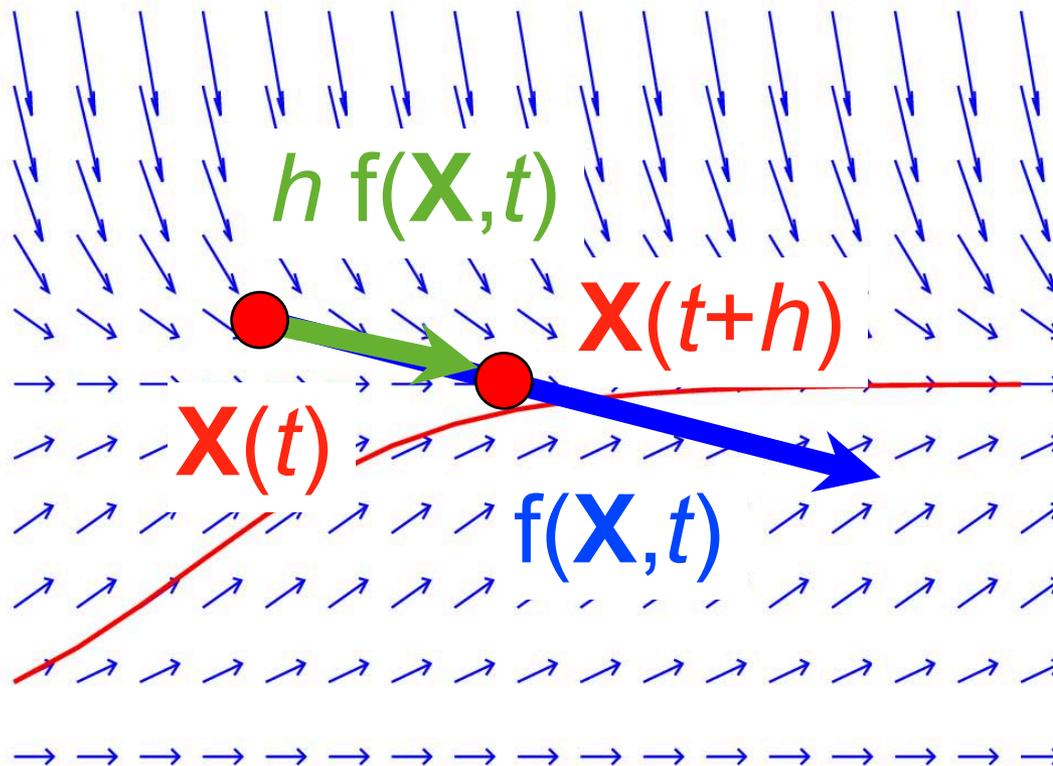


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# Euler's Method: Inaccurate

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- Moves along tangent; can leave solution curve, e.g.:

$$f(\mathbf{X}, t) = \begin{pmatrix} -y \\ x \end{pmatrix}$$

- Exact solution is circle:

$$\mathbf{X}(t) = \begin{pmatrix} r \cos(t+k) \\ r \sin(t+k) \end{pmatrix}$$

- Euler spirals outward  
no matter how small  $h$  is
  - will just diverge more slowly

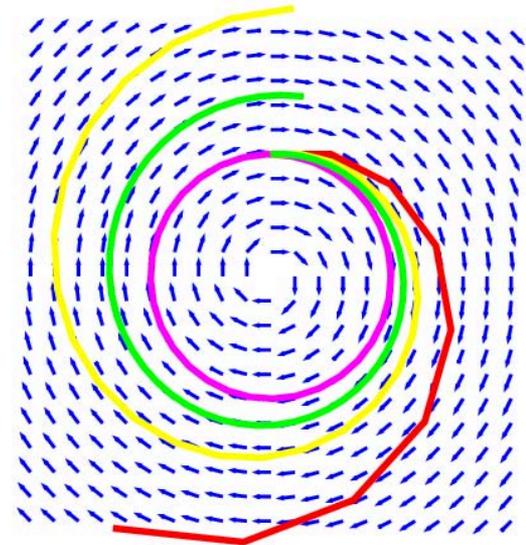


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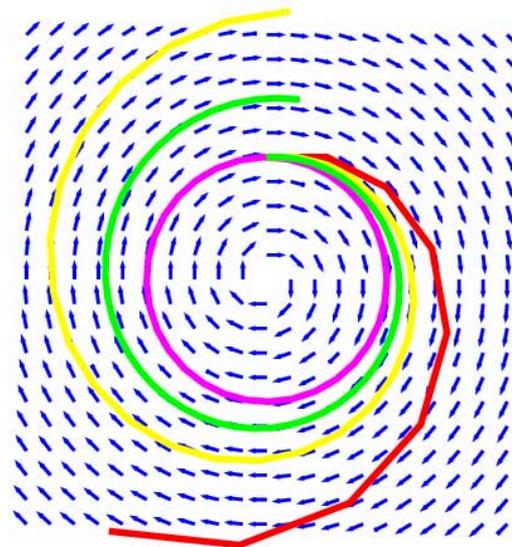


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## Questions?

# Euler's Method: Not Always Stable

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- “Test equation”  $f(x, t) = -kx$

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$$x(t) = x_0 e^{-kt}$$

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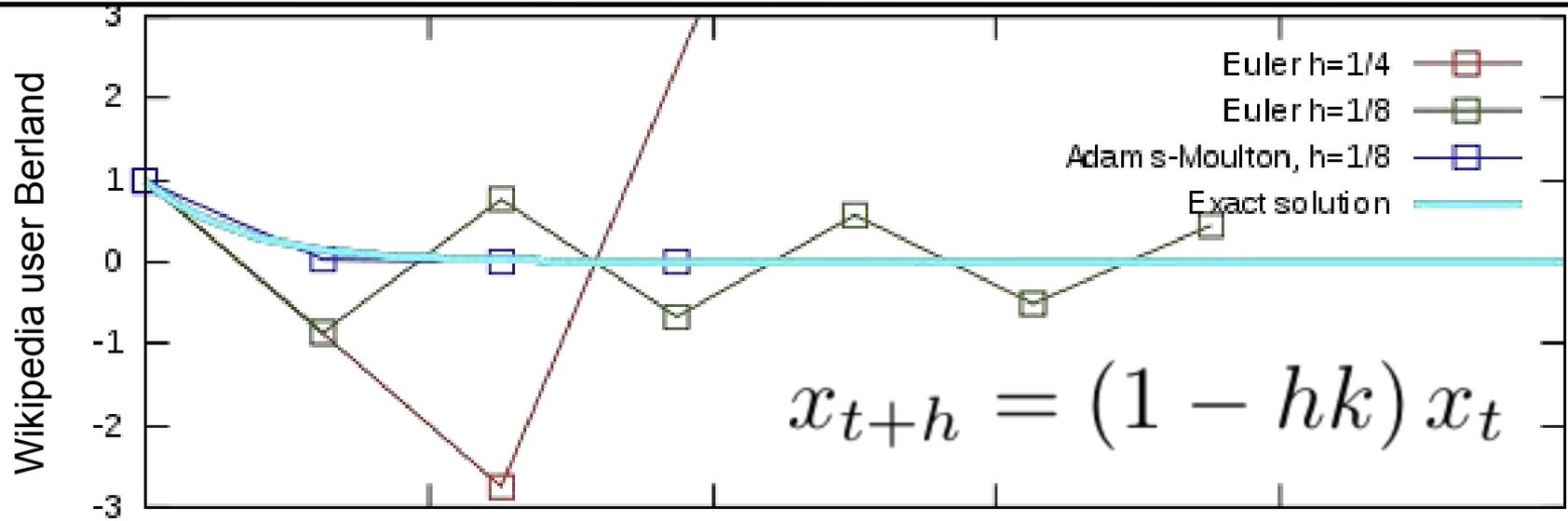
- Let's apply Euler's method:

$$x_{t+h} = x_t + h f(x_t, t)$$

$$= x_t - hkx_t$$

$$= (1 - hk) x_t$$

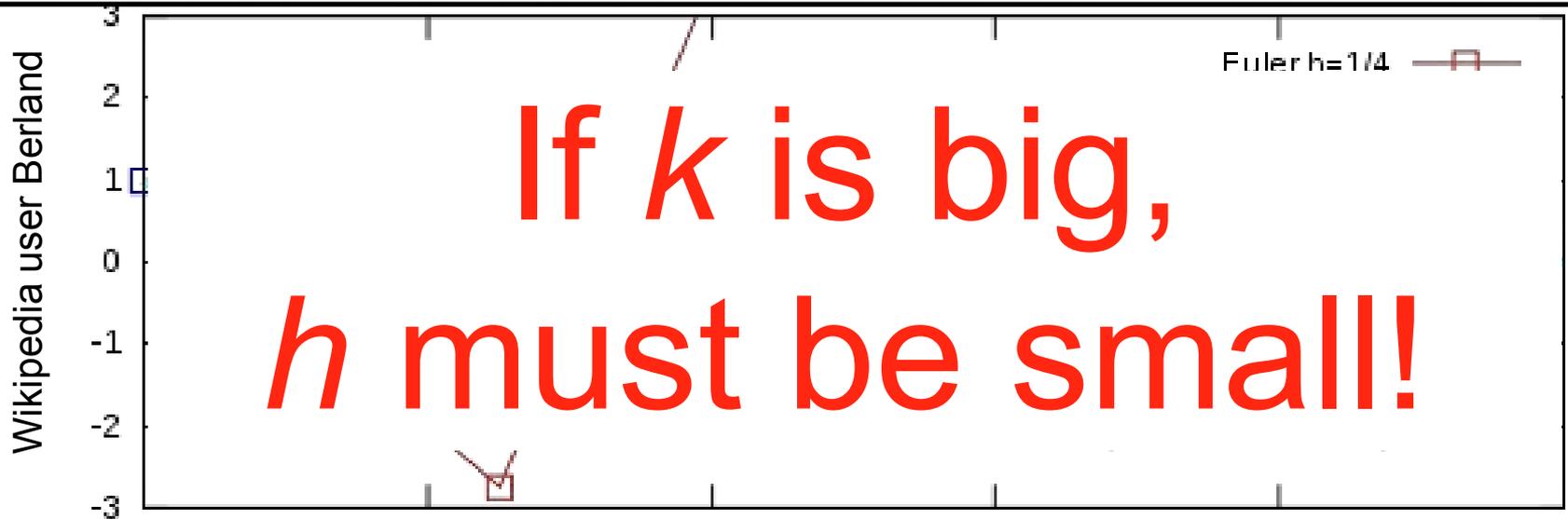
# Euler's Method: Not Always Stable



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- Limited step size!
  - When  $0 \leq (1 - hk) < 1 \Leftrightarrow h < 1/k$   
things are fine, the solution decays
  - When  $-1 \leq (1 - hk) \leq 0 \Leftrightarrow 1/k \leq h \leq 2/k$   
we get oscillation
  - When  $(1 - hk) < -1 \Leftrightarrow h > 2/k$  things explode

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# Analysis: Taylor Series

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- Expand exact solution  $\mathbf{X}(t)$

$$\mathbf{X}(t_0 + h) = \mathbf{X}(t_0) + h \left( \frac{d}{dt} \mathbf{X}(t) \right) \Big|_{t_0} + \frac{h^2}{2!} \left( \frac{d^2}{dt^2} \mathbf{X}(t) \right) \Big|_{t_0} + \frac{h^3}{3!} (\dots) + \dots$$

- Euler's method approximates:

$$\mathbf{X}(t_0 + h) = \mathbf{X}_0 + h f(\mathbf{X}_0, t_0) \quad \dots + O(h^2) \text{ error}$$

$$h \rightarrow h/2 \Rightarrow \text{error} \rightarrow \text{error}/4 \text{ per step} \times \text{twice as many steps} \\ \rightarrow \text{error}/2$$

- First-order method: Accuracy varies with  $h$
- To get 100x better accuracy need 100x more steps

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# Can We Do Better?

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- Problem:  $f$  varies along our Euler step
- Idea 1: look at  $f$  at the arrival of the step and compensate for variation

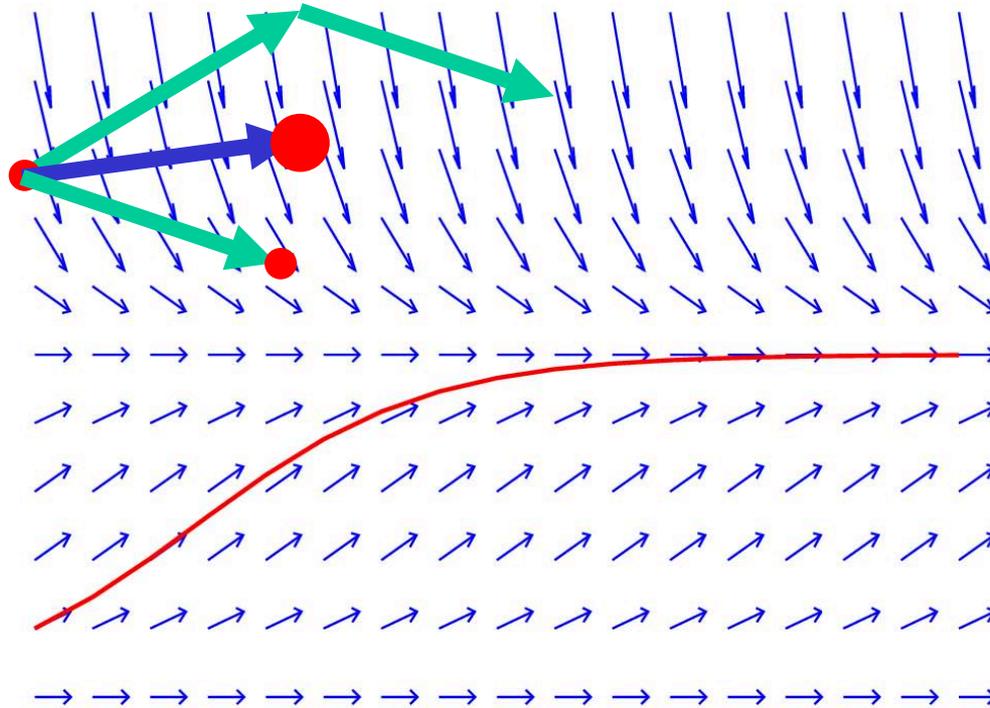


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# 2nd Order Methods

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- This translates to...

$$\begin{aligned} f_0 &= f(\mathbf{X}_0, t_0) \\ f_1 &= f(\mathbf{X}_0 + h f_0, t_0 + h) \end{aligned}$$

- and we get

$$\mathbf{X}(t_0 + h) = \mathbf{X}_0 + \frac{h}{2}(f_0 + f_1) + O(h^3)$$

- This is the *trapezoid method*
  - Analysis omitted (see 6.839)
- Note: What we mean by “2<sup>nd</sup> order” is that the error goes down with  $h^2$ , not  $h$  – the equation is still 1<sup>st</sup> order!

# Can We Do Better?

---

- Problem:  $f$  has varied along our Euler step
- Idea 2: look at  $f$  after a smaller step, use that value for a full step from initial position

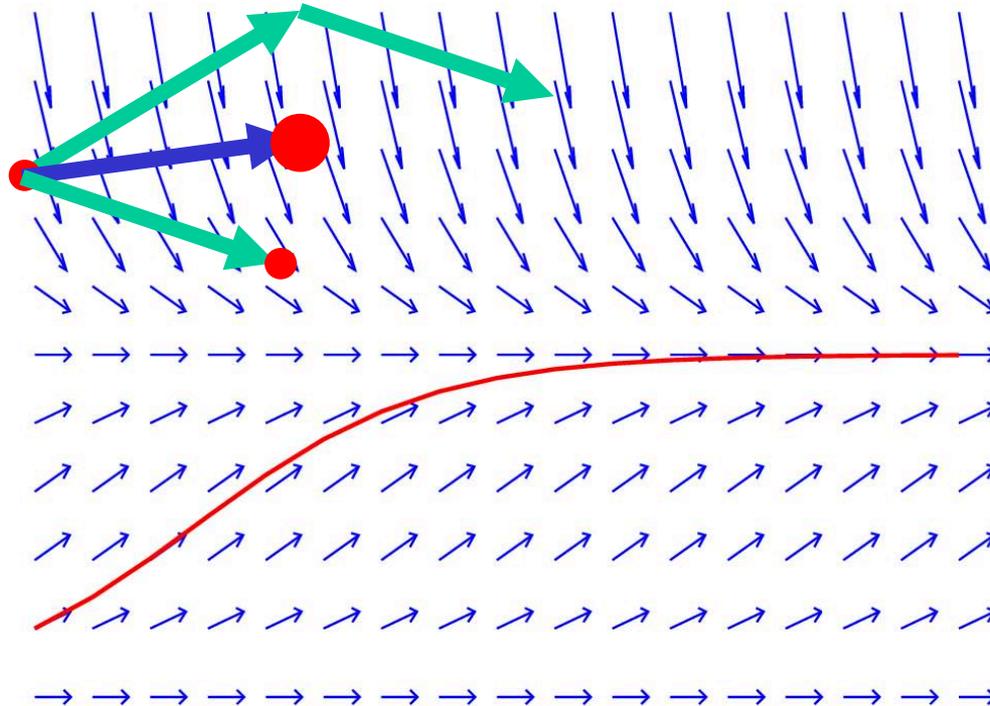


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# 2nd Order Methods Cont'd

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- This translates to...

$$\begin{aligned} f_0 &= f(\mathbf{X}_0, t_0) \\ f_m &= f\left(\mathbf{X}_0 + \frac{h}{2} f_0, t_0 + \frac{h}{2}\right) \end{aligned}$$

- and we get  $\mathbf{X}(t_0 + h) = \mathbf{X}_0 + h f_m + O(h^3)$

- This is the *midpoint method*

- Analysis omitted again,  
but it's not very complicated, see [here](#).

# Comparison

- **Midpoint:**
  - $\frac{1}{2}$  Euler step
  - evaluate  $f_m$
  - full step using  $f_m$
- **Trapezoid:**
  - Euler step (a)
  - evaluate  $f_1$
  - full step using  $f_1$  (b)
  - average (a) and (b)
- Not exactly same result,  
but same order of accuracy

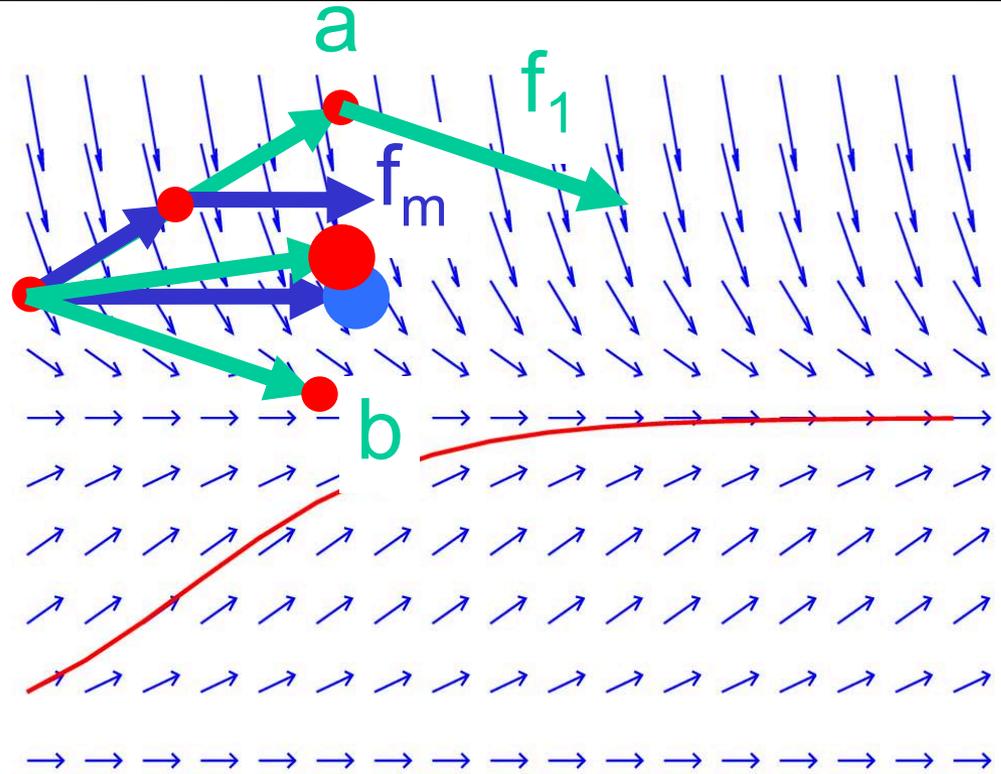


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# Can We Do Even Better?

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- You bet!
- You will implement **Runge-Kutta** for assignment 3
- Again, see **Witkin, Baraff, Kass: Physically-based Modeling Course Notes, SIGGRAPH 2001**
- See eg <http://www.youtube.com/watch?v=HbE3L5CIdQg>

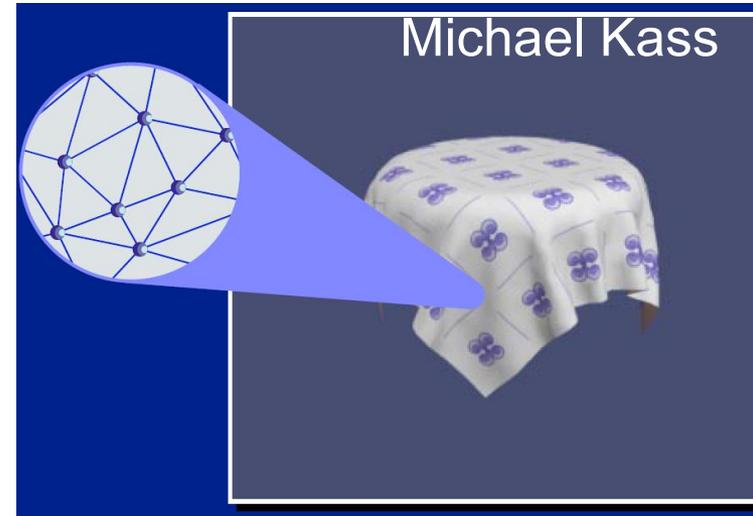
# Can We Do Even Better? Questions?

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# Mass-Spring Modeling

- Beyond pointlike objects: strings, cloth, hair, etc.
- Interaction between particles
  - Create a network of spring forces that link pairs of particles



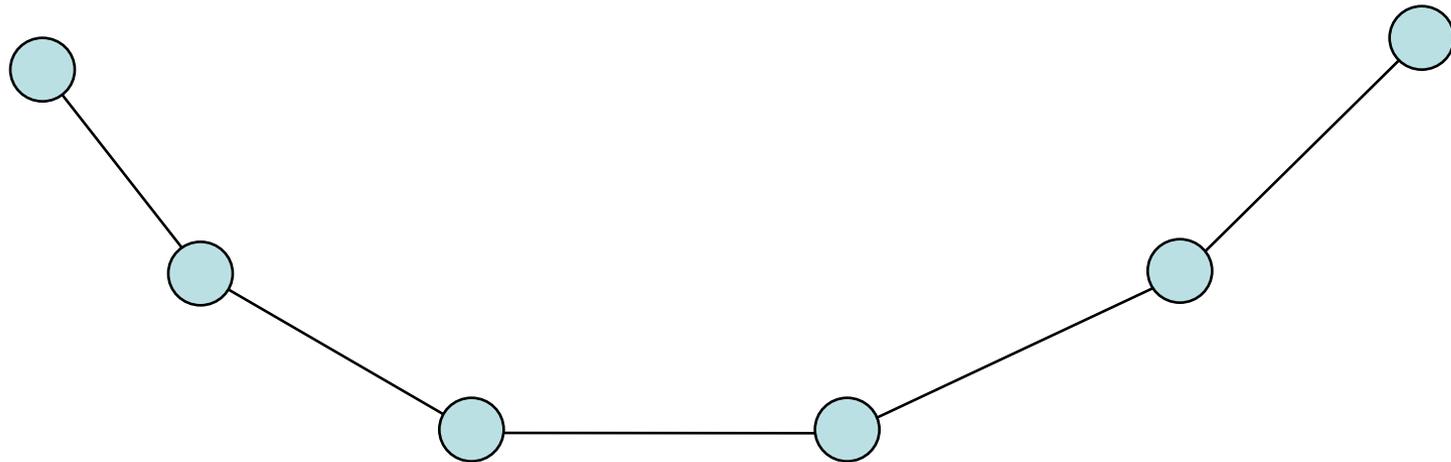
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- First, slightly hacky version of cloth simulation
- Then, some motivation/intuition for *implicit integration* (NEXT LECTURE)

# How Would You Simulate a String?

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- Each particle is linked to two particles (except ends)
- Come up with forces that try to keep the distance between particles constant



# Springs

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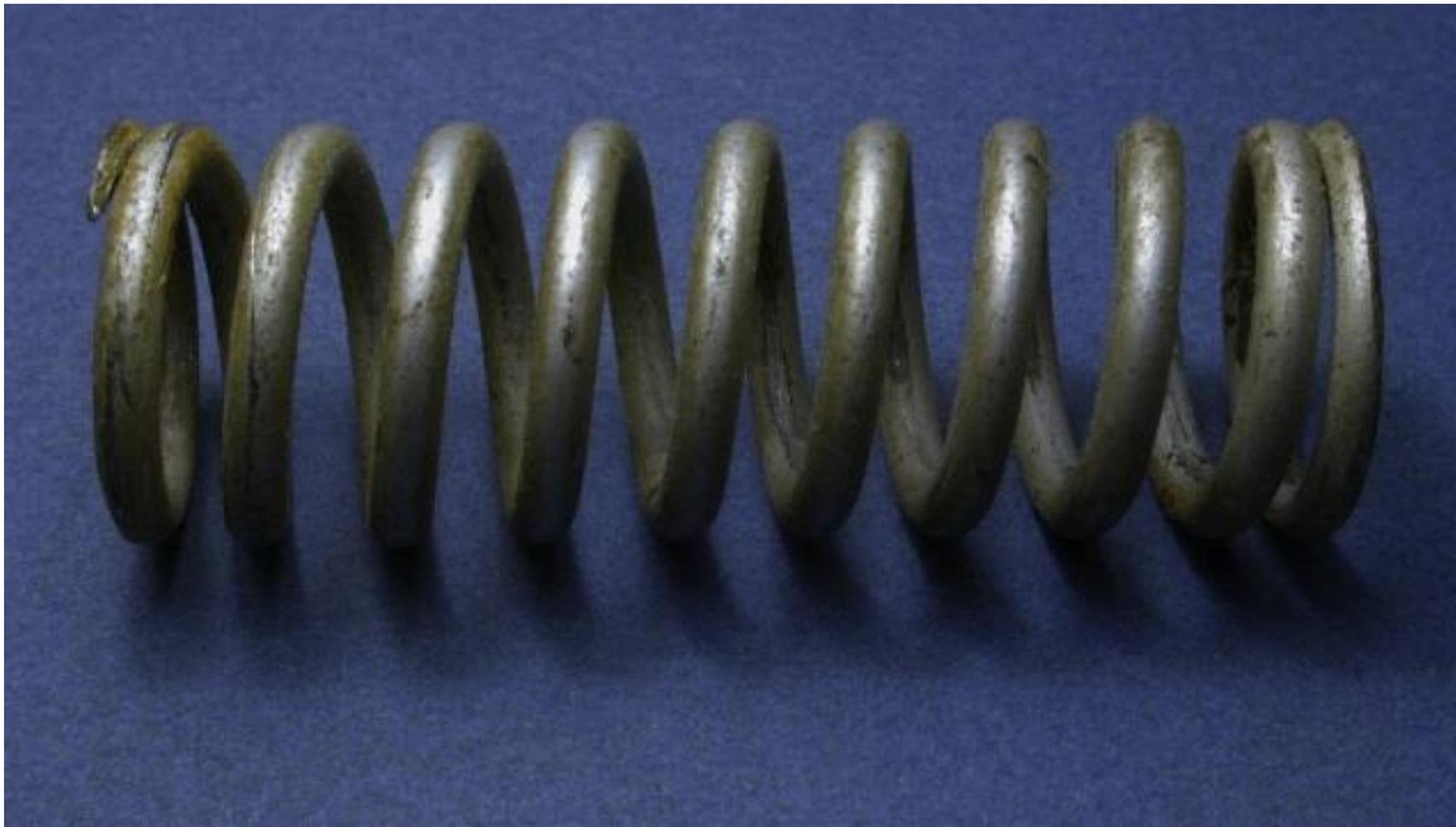
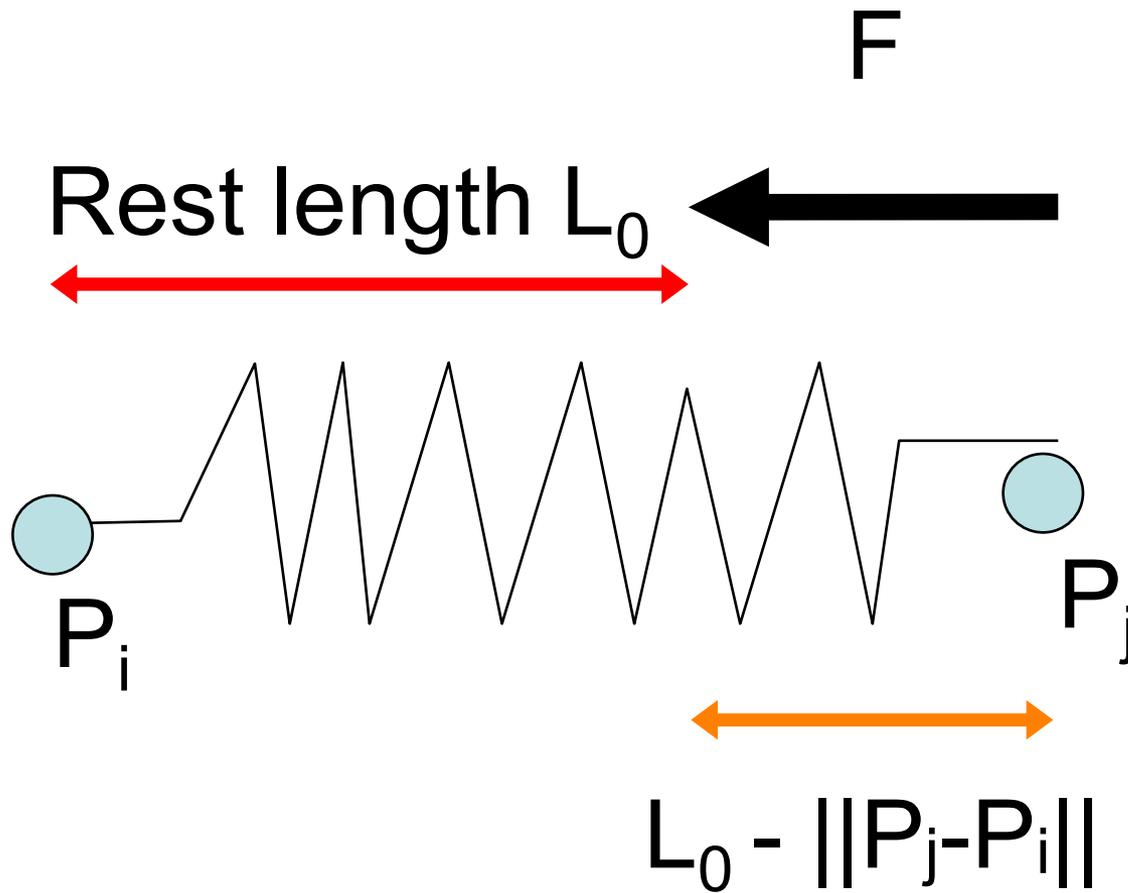


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# Spring Force – Hooke's Law

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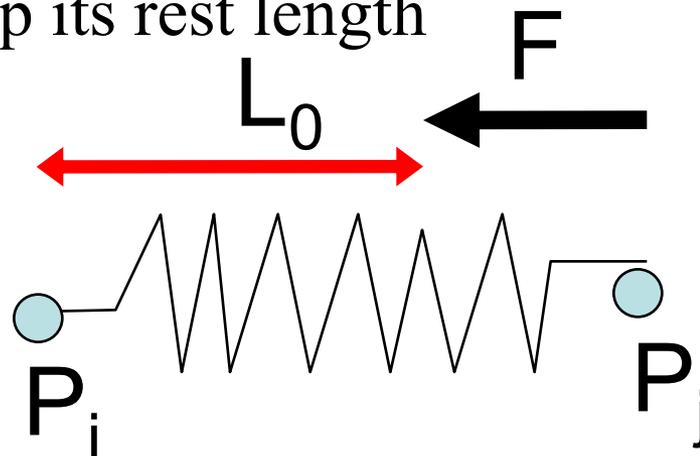


# Spring Force – Hooke's Law

- Force in the direction of the spring and proportional to difference with rest length  $L_0$ .

$$F(P_i, P_j) = K(L_0 - \|P_i\vec{P}_j\|) \frac{P_i\vec{P}_j}{\|P_i\vec{P}_j\|}$$

- K is the stiffness of the spring
  - When K gets bigger, the spring *really* wants to keep its rest length

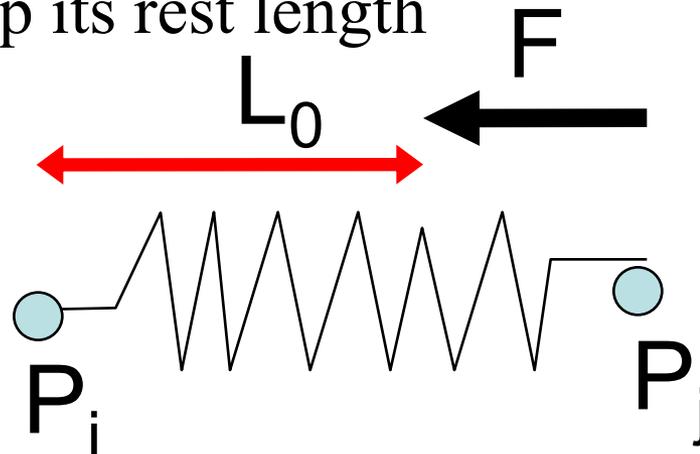


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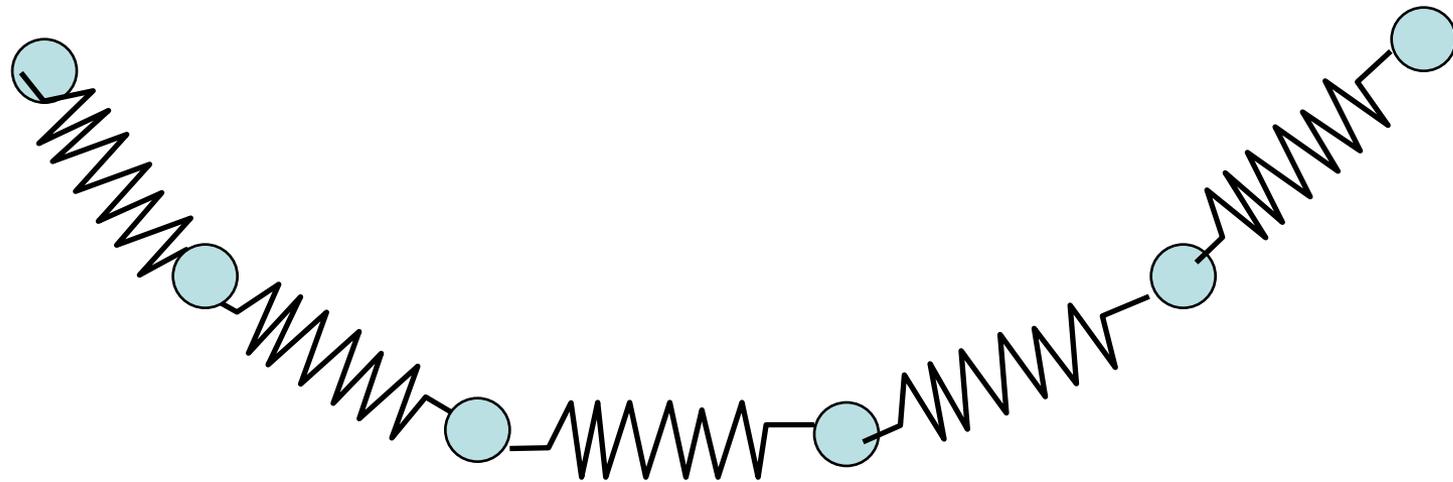


This is the force on  $P_j$ .  
**Remember Newton:**  
 $P_i$  experiences force of equal magnitude but opposite direction.

# How Would You Simulate a String?

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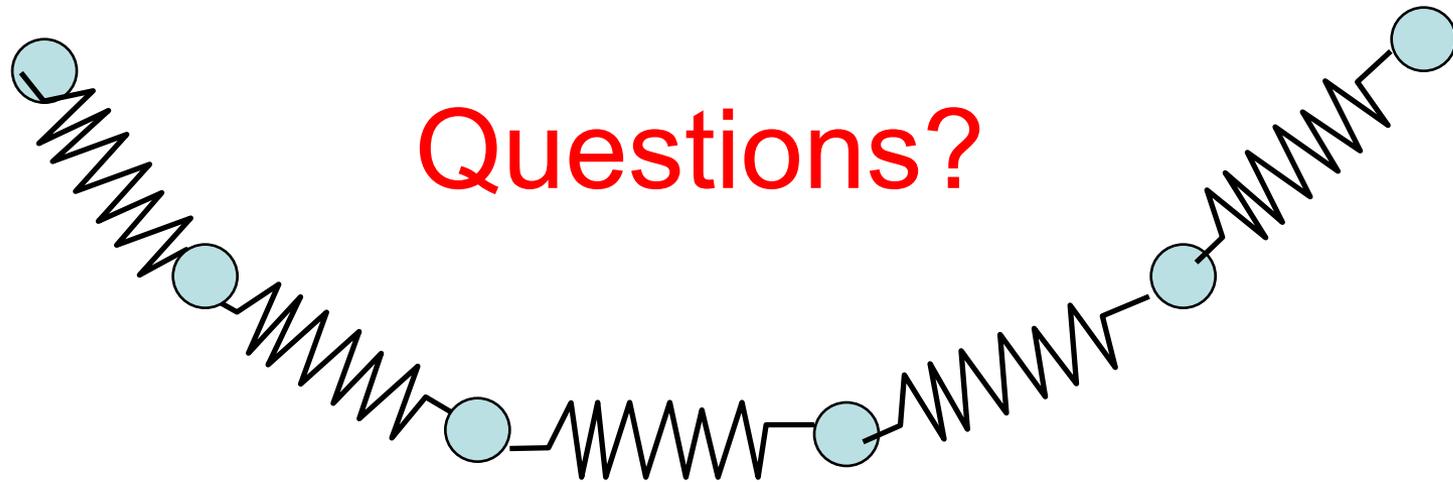
- Springs link the particles
- Springs try to keep their rest lengths and preserve the length of the string
- Not exactly preserved though, and we get oscillation
  - Rubber band approximation



# How Would You Simulate a String?

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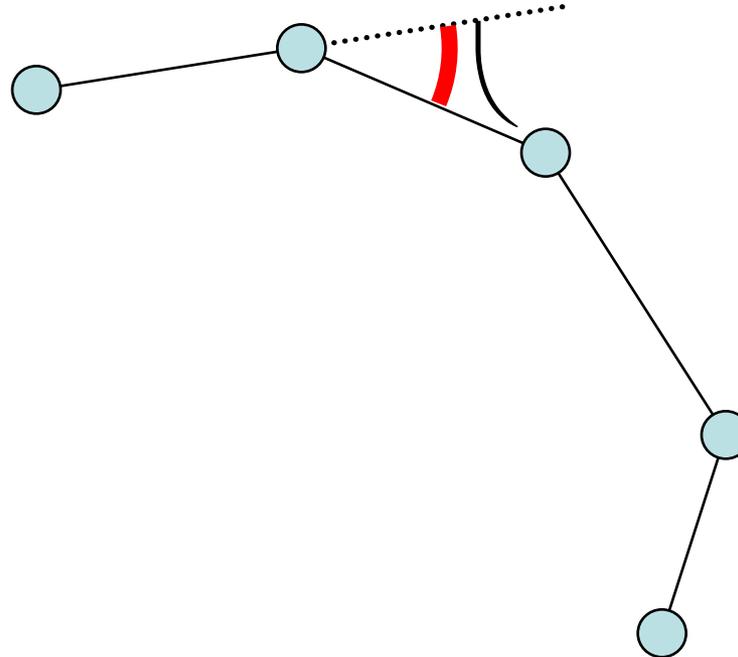
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# Hair

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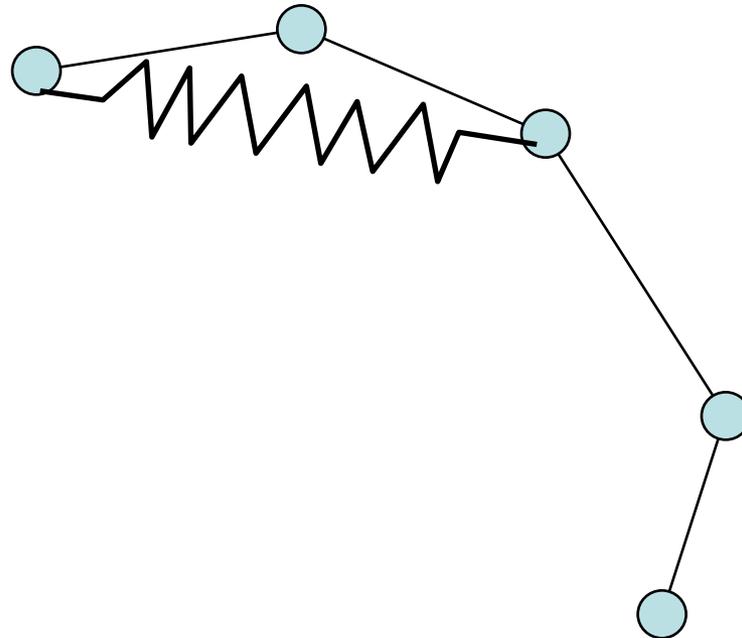
- Linear set of particles
- Length-preserving **structural** springs like before
- **Deformation** forces proportional to the angle between segments
- **External** forces



# Hair - Alternative Structural Forces

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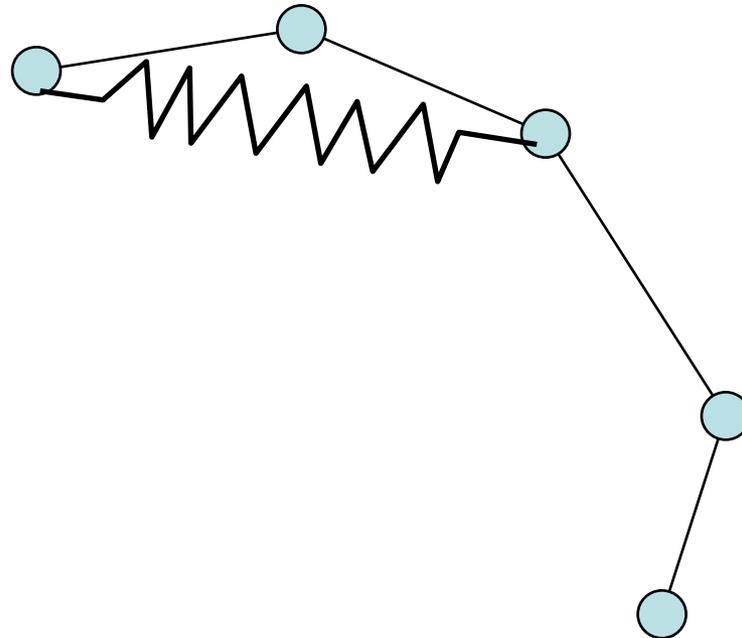
- Springs between mass  $n$  &  $n+2$  with rest length  $2L_0$ 
  - Wants to keep particles aligned



# Hair - Alternative Structural Forces

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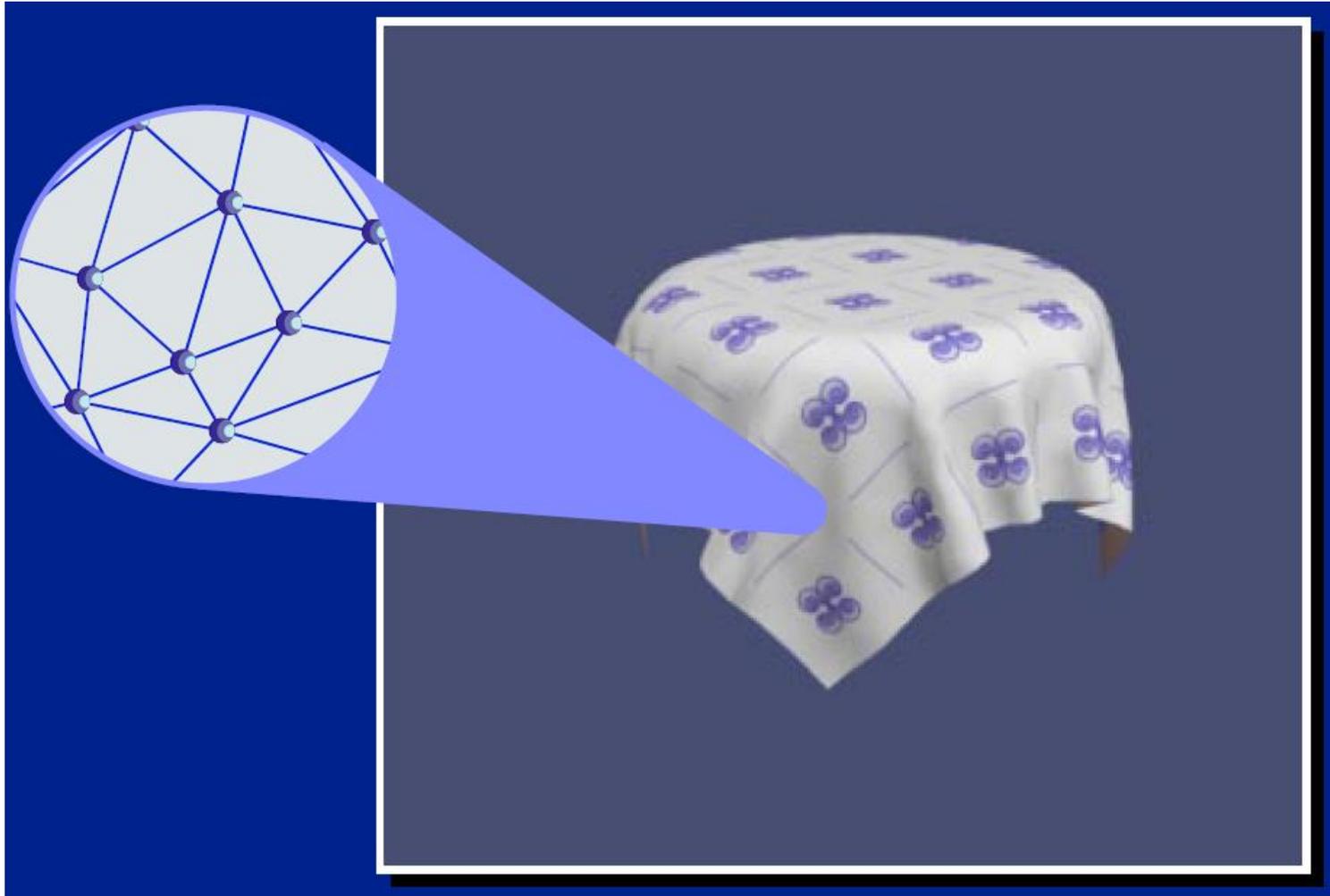
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Questions?

# Mass-Spring Cloth

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Michael Kass

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# Cloth – Three Types of Forces

- **Structural forces**

- Try to enforce invariant properties of the system

- E.g. force the distance between two particles to be constant

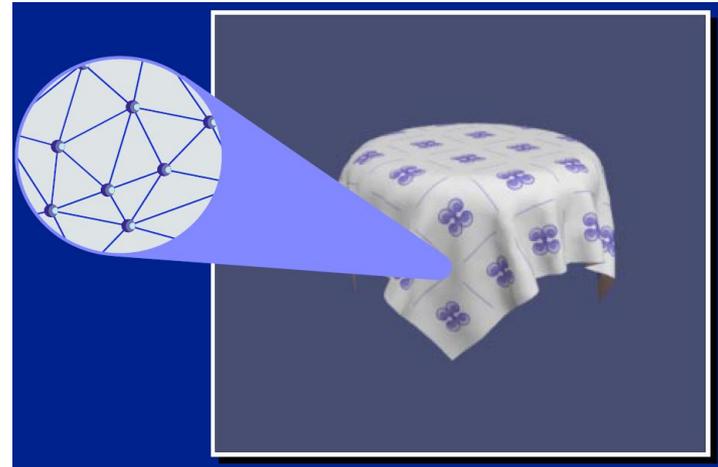
- Ideally, these should be *constraints*, not forces

- **Internal deformation forces**

- E.g. a string deforms, a spring board tries to remain flat

- **External forces**

- Gravity, etc.



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# Springs for Cloth

- Network of masses and springs
- **Structural** springs:
  - link  $(i, j)$  and  $(i+1, j)$ ;  
and  $(i, j)$  and  $(i, j+1)$
- **Deformation:**
  - Shear springs
    - $(i, j)$  and  $(i+1, j+1)$
  - Flexion springs
    - $(i, j)$  and  $(i+2, j)$ ;  
 $(i, j)$  and  $(i, j+2)$
- See Provot's Graphics Interface '95 paper for details

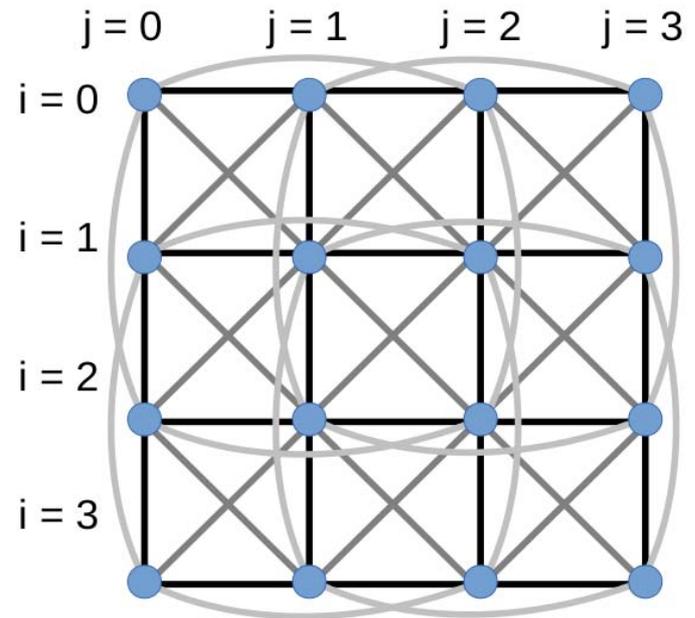


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Provot 95

# External Forces

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- Gravity  $G$
- Friction
- Wind, etc.

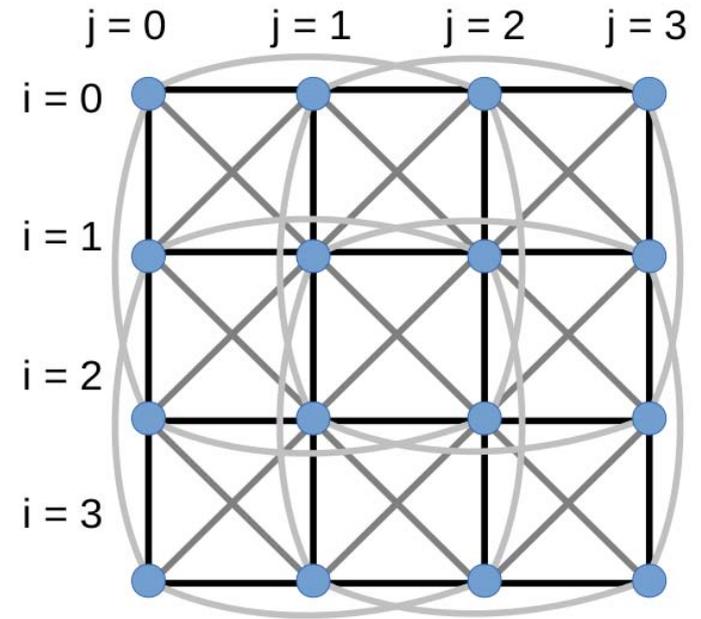


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Provot 95

# Cloth Simulation

- Then, the all trick is to set the stiffness of all springs to get realistic motion!
- Remember that forces depend on other particles (coupled system)
- But it is *sparse* (only near neighbors)
  - This is in contrast to e.g. the N-body problem.

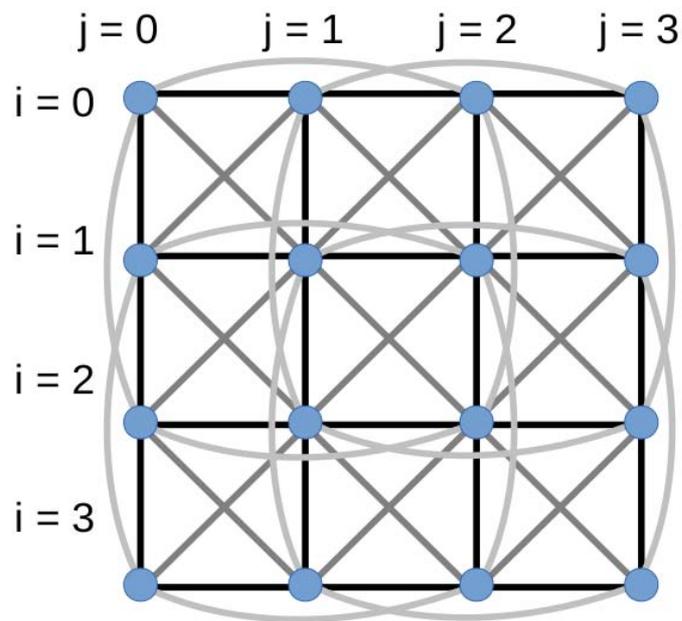


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Provot 95

# Forces: Structural vs. Deformation

- Structural forces are here just to enforce a constraint
- Ideally, the constraint would be enforced strictly
  - at least a lot more than we can afford
- We'll see that this is the root of a lot of problems
- In contrast, deformation forces actually correspond to physical forces

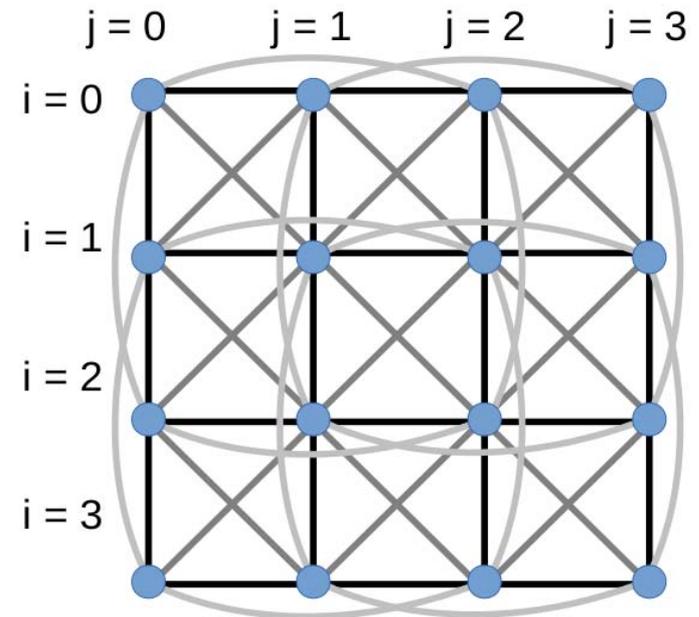


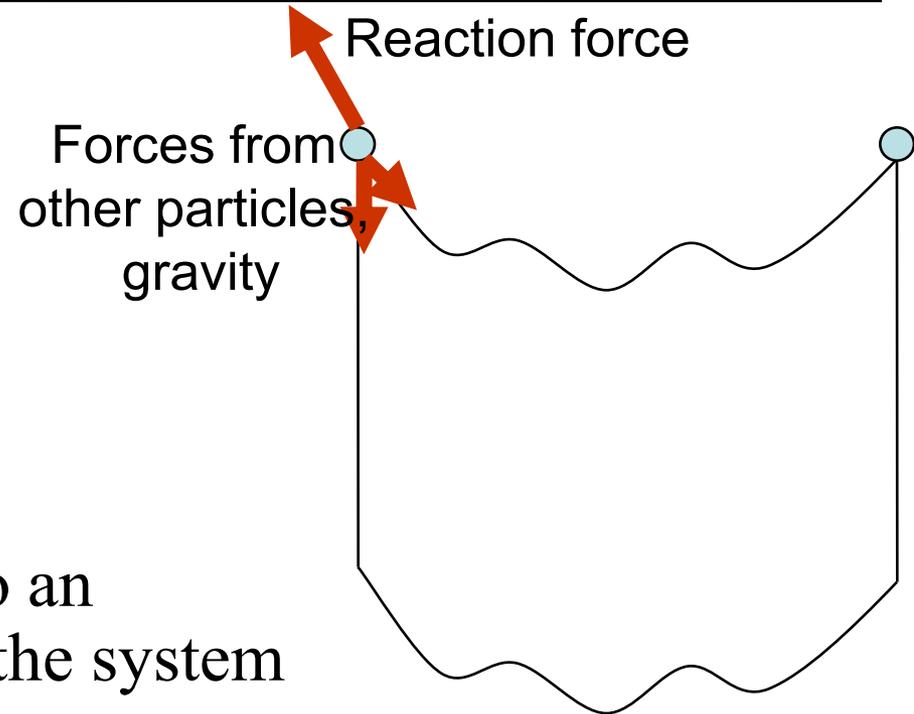
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Provot 95

# Contact Forces

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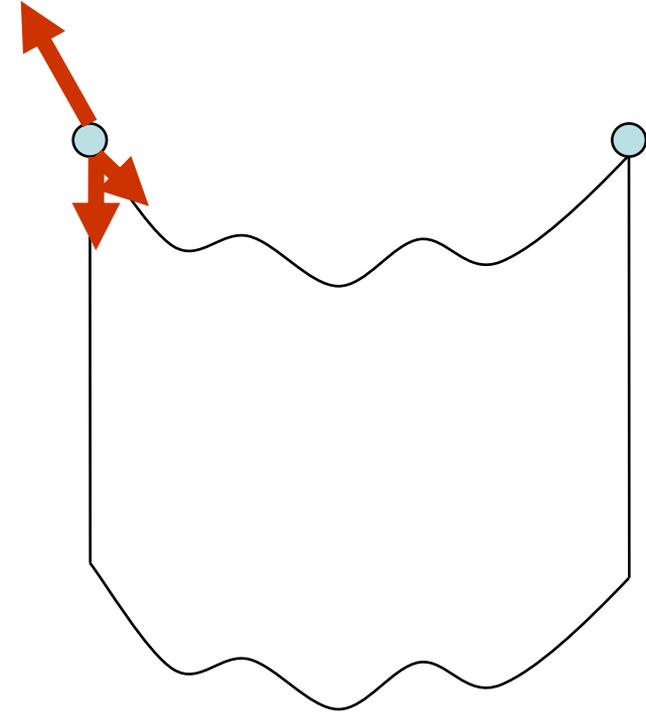
- Hanging curtain:
  - 2 contact points stay fixed
- What does it mean?
  - Sum of the forces is zero
- How so?
  - Because those point undergo an external force that balances the system
- What is the force at the contact?
  - Depends on all other forces in the system
  - Gravity, wind, etc.



# Contact Forces

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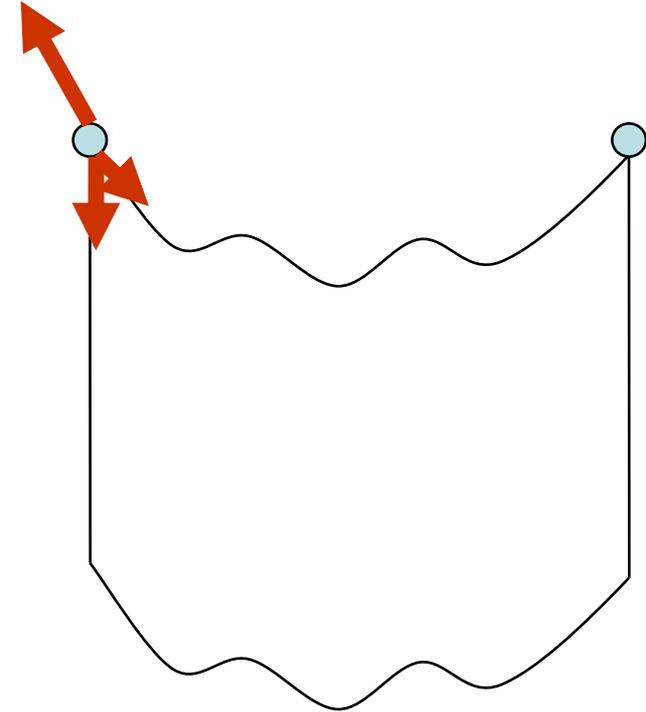
- How can we compute the external contact force?
  - Inverse dynamics!
  - Sum all other forces applied to point
  - Take negative
- Do we really need to compute this force?
  - Not really, just ignore the other forces applied to this point!



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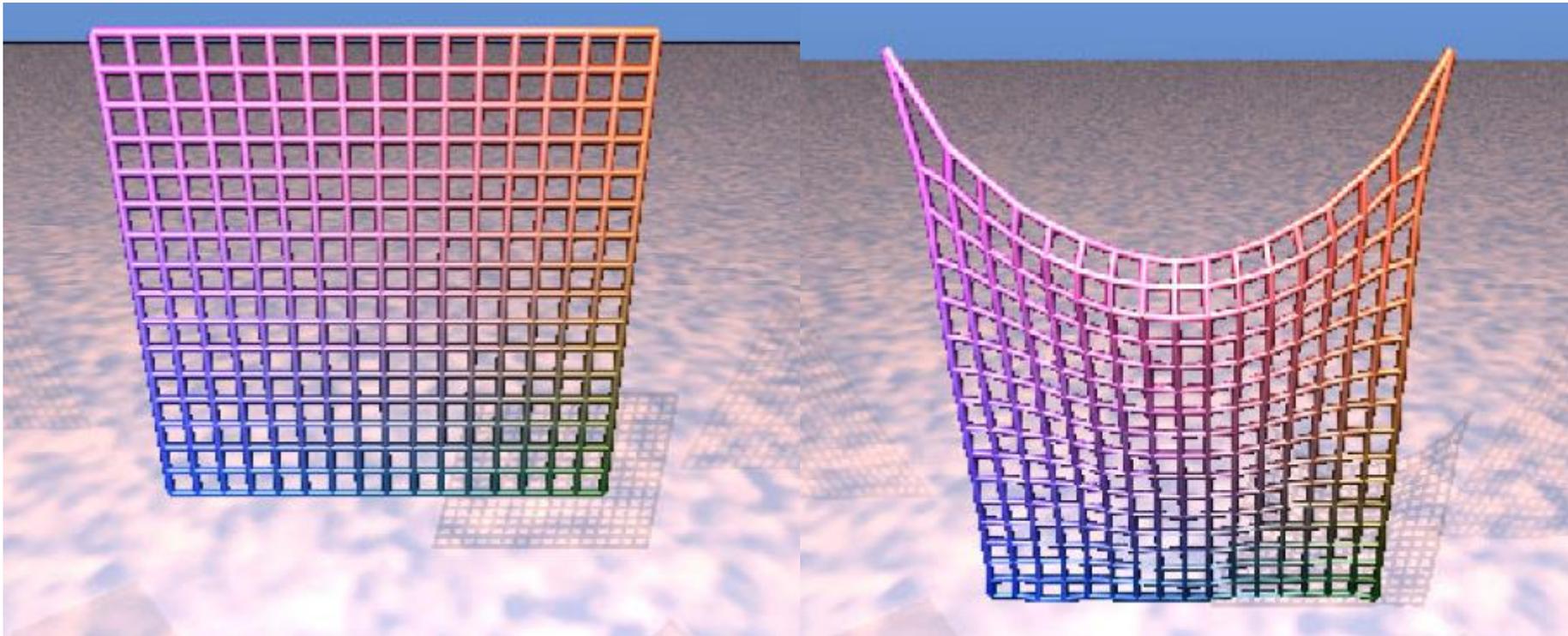


Questions?

# Example

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- Excessive rubbery deformation:  
the strings are not stiff enough



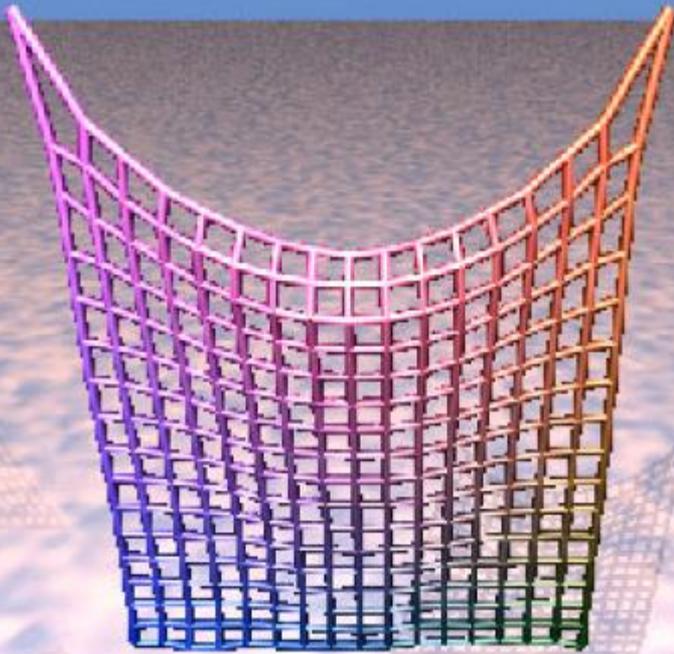
Initial position

After 200 iterations

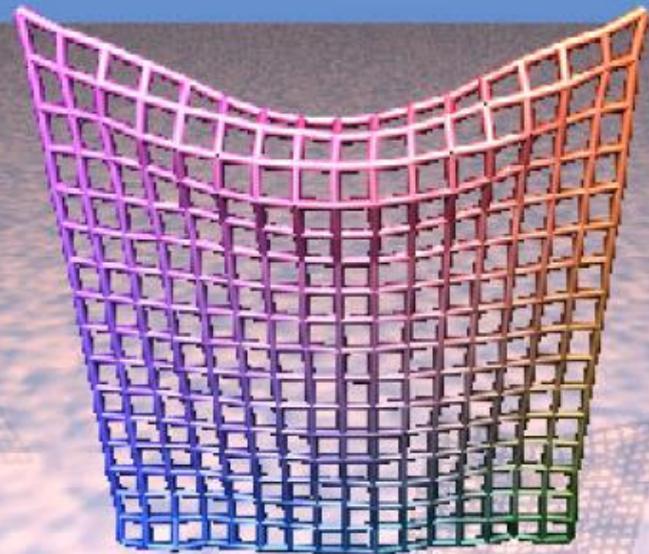
# One Solution

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- Constrain length to increase by less than 10%
  - A little hacky



Simple mass-spring system

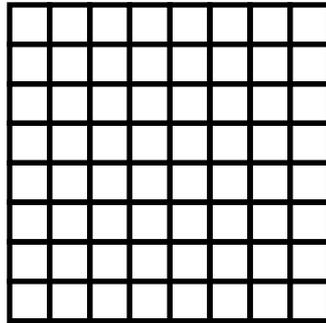
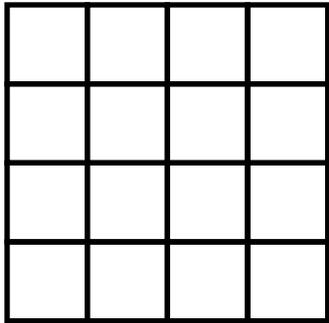


Improved solution  
(see Provot Graphics Interface 1995)

# The Discretization Problem

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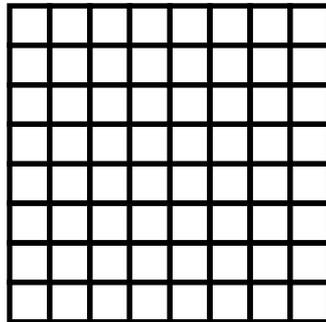
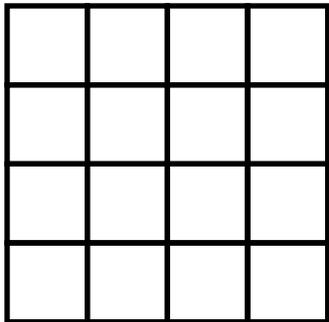
- What happens if we discretize our cloth more finely?
- Do we get the same behavior?
- Usually not! It takes a lot of effort to design a scheme that is mostly oblivious to the discretization.



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Questions?

# The Stiffness Issue

---

- We use springs while we really mean constraint
  - Spring should be super stiff, which requires tiny  $\Delta t$
  - Remember  $x' = -kx$  system and Euler speed limit!
    - The story extends to N particles and springs (unfortunately)
- Many numerical solutions
  - Reduce  $\Delta t$  (well, not a great solution)
  - Actually use constraints (see 6.839)
  - Implicit integration scheme (more next Thursday)

# Euler Has a Speed Limit!

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- $h > 1/k$ : oscillate.  $h > 2/k$ : explode!

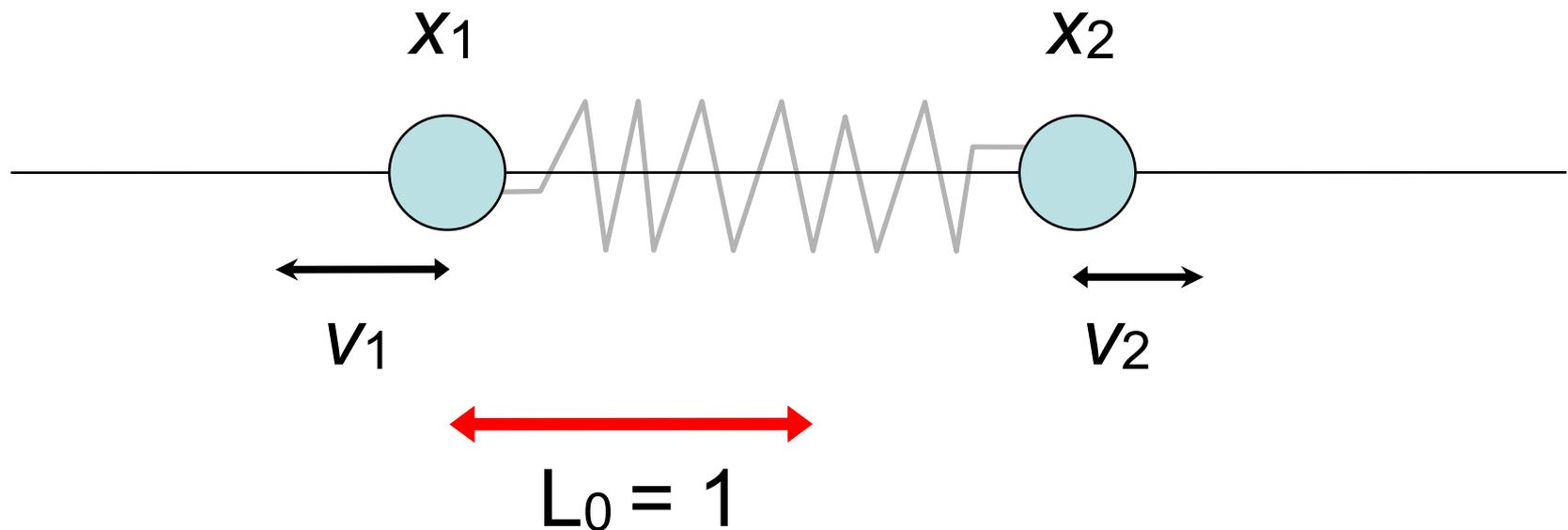
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From the SIGGRAPH PBM notes

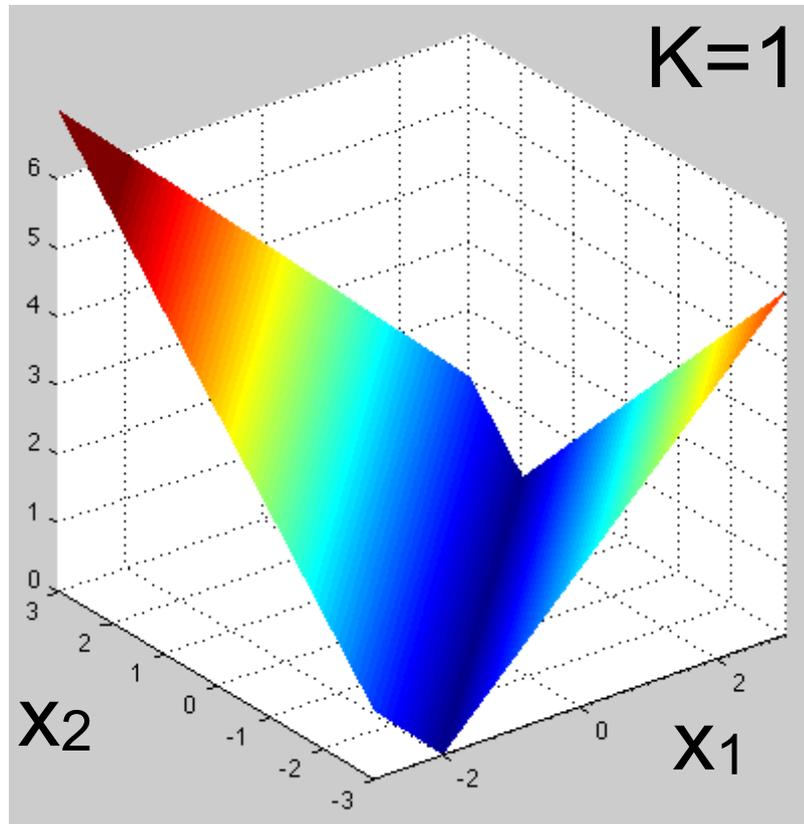
# Why Stiff Springs Are Difficult

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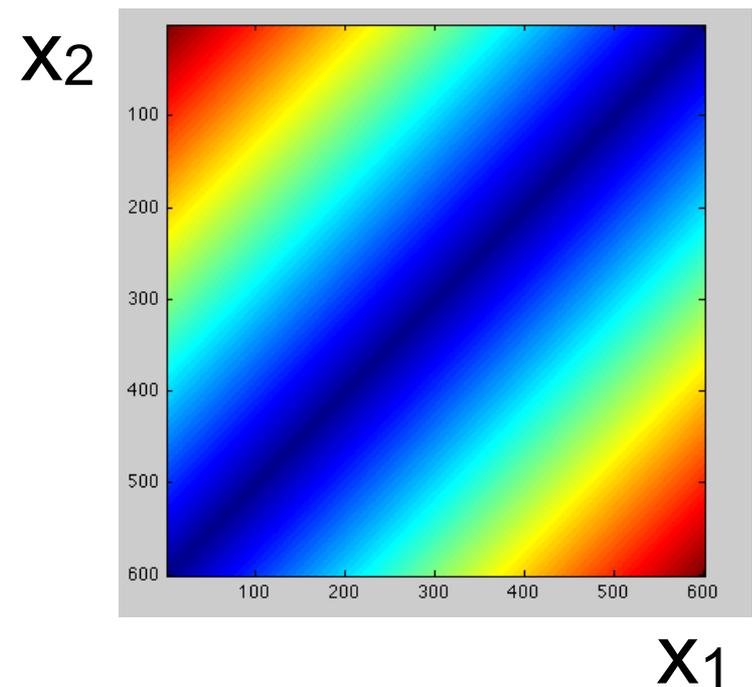
- 1D example, with two particles constrained to move along the  $x$  axis only, rest length  $L_0 = 1$
- Phase space is 4D:  $(x_1, v_1, x_2, v_2)$ 
  - But spring force only depends on  $x_1, x_2$  and  $L_0$ .



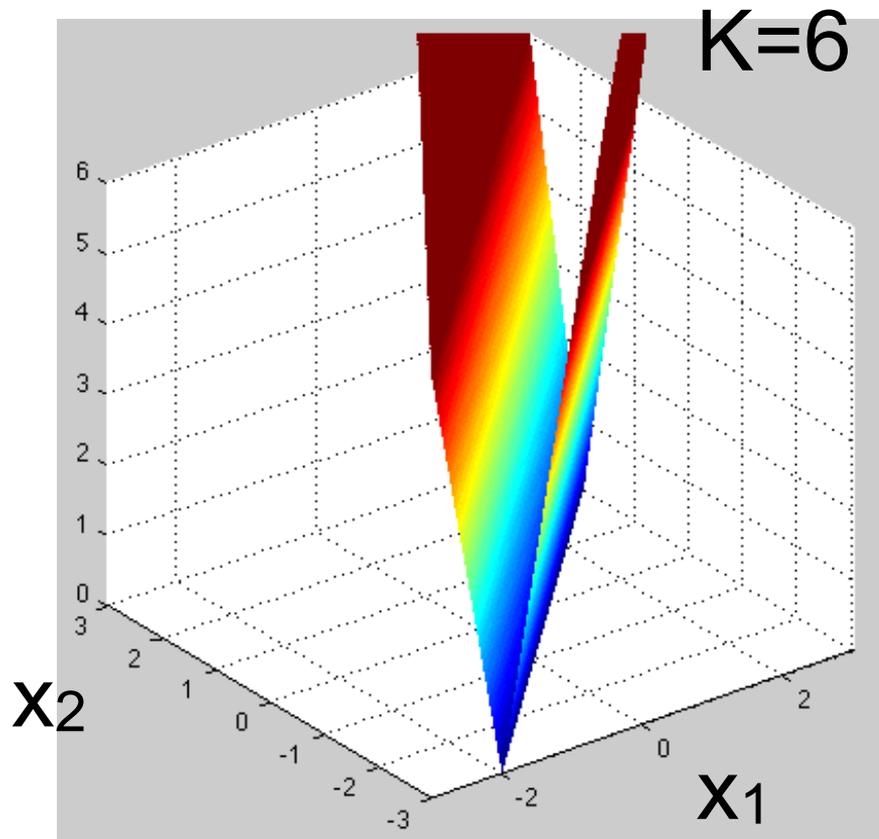
# Why Stiff Springs Are Difficult



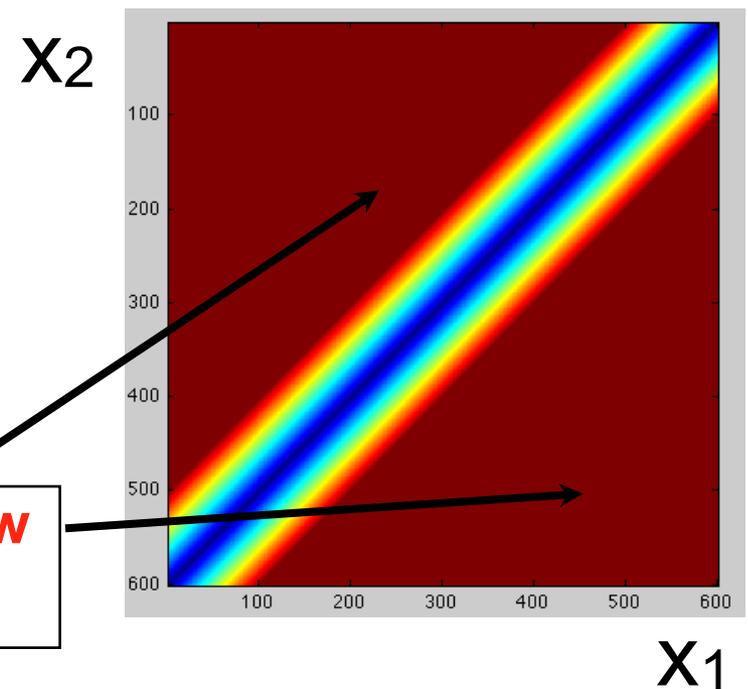
height=magnitude  
of spring force



# Why Stiff Springs Are Difficult

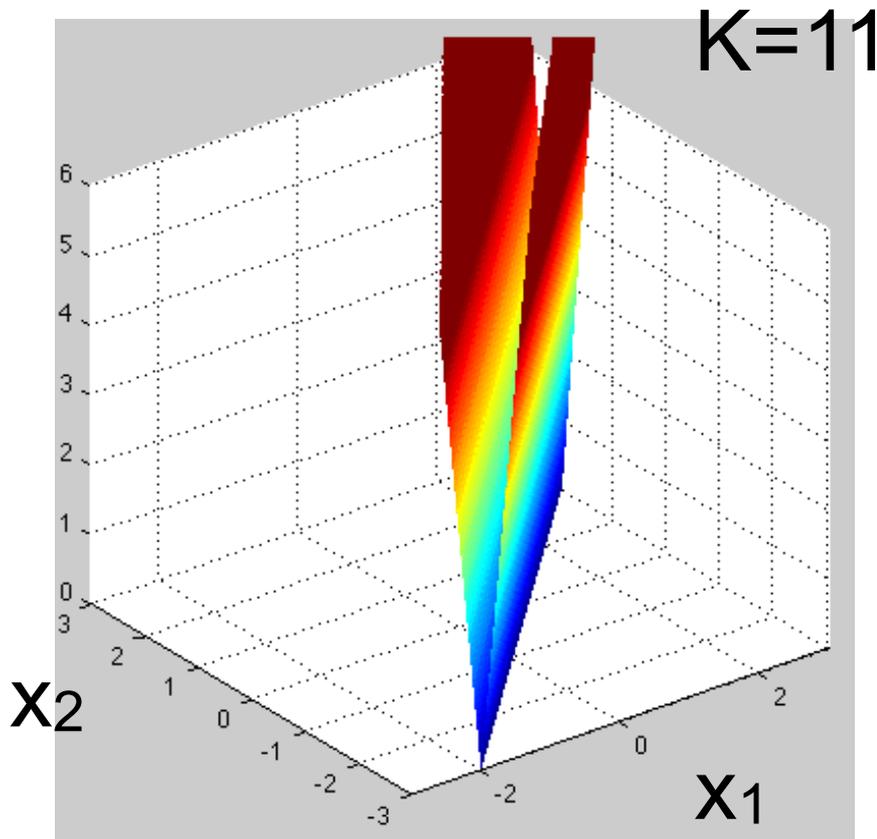


height=magnitude  
of spring force

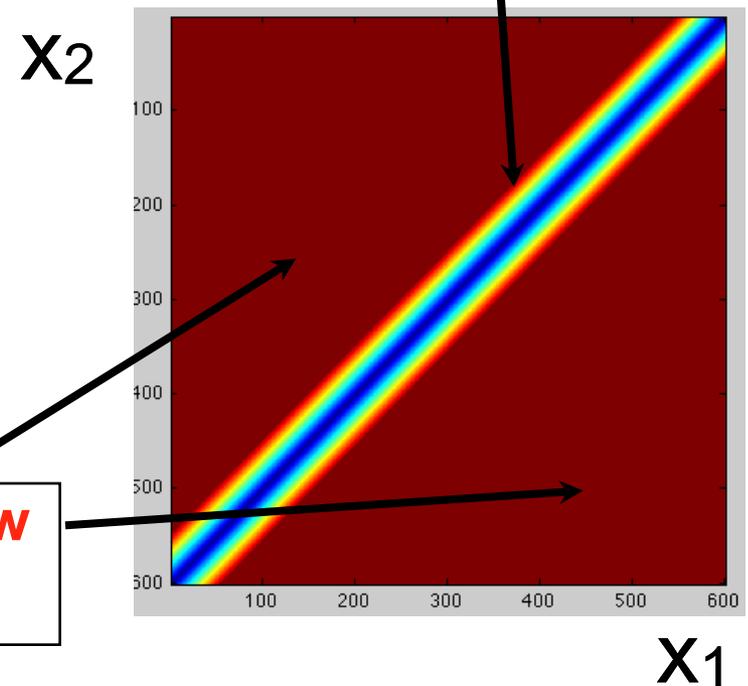


**Forces grow  
really big!**

# Why Stiff Springs Are Difficult

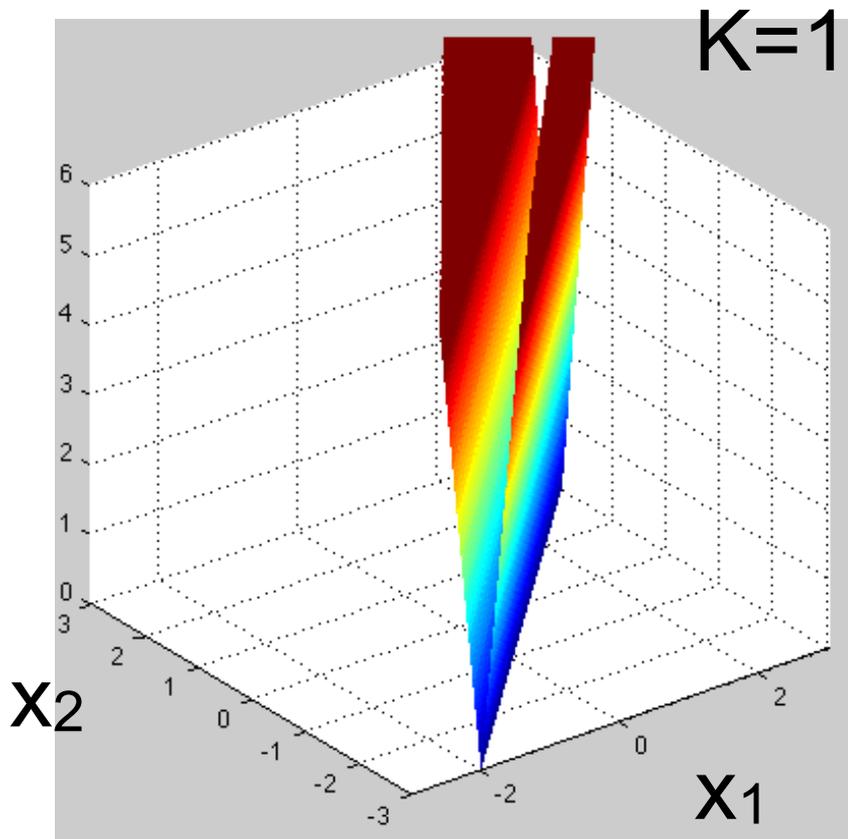


The “admissible region”  
shrinks towards the line  
 $x_1 - x_2 = 1$  as  $K$  grows



Forces grow  
really big!

# Why Stiff Springs Are Difficult

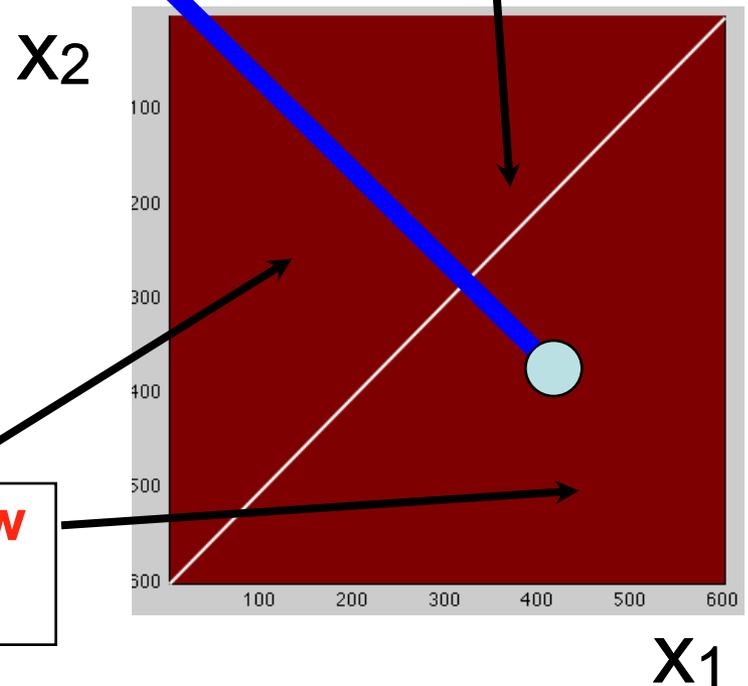


$K=11$

*off to the moon!*

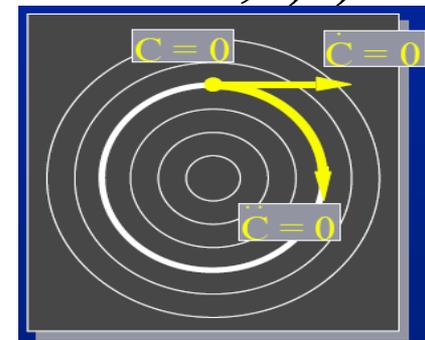
The “admissible region”  
shrinks towards the line  
 $x_1 - x_2 = 1$  as  $K$  grows

**Forces grow  
really big!**



# Constrained Dynamics

- In our mass-spring cloth, we have “encouraged” length preservation using springs that want to have a given length (unfortunately, they can refuse offer ;-)
- Constrained dynamic simulation: force it to be constant!
- How it works – **more in 6.839**
  - Start with constraint equation
    - E.g.,  $(x_2 - x_1) - l = 0$  in the previous 1D example
  - Derive extra forces that will exactly enforce constraint
    - This means *projecting* the external forces (like gravity) onto the “subspace” of phase space where constraints are satisfied
    - Fancy name for this: “Lagrange multipliers”
  - Again, see the SIGGRAPH 2001 Course Notes



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# Questions?

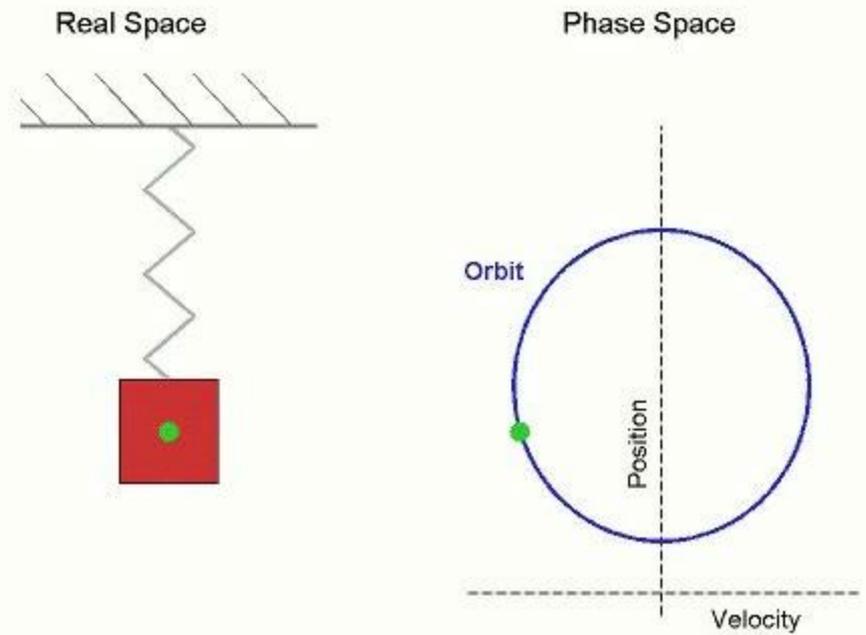
---

- Further reading
  - Stiff systems
  - Explicit vs. implicit solvers
  - Again, consult the 2001 course notes!

# Mass on a Spring, Phase Space

- State of system (phase) : velocity & position
  - similar to our  $X=(x \ v)$  to get 1st order

Wikipedia user Mazemaster

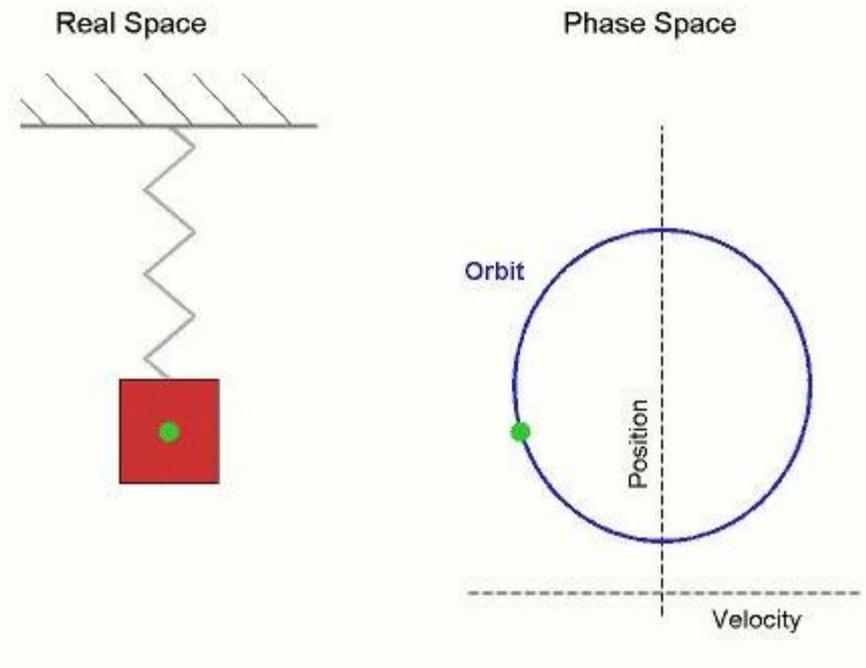


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# Mass on a Spring, Phase Space

- Guess how well Euler will do...  
always diverge

Wikipedia user Mazemaster



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# Difference with $x' = -kx$

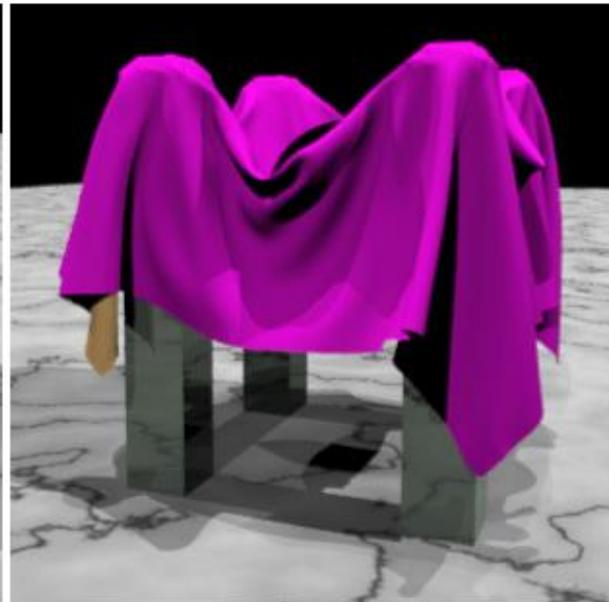
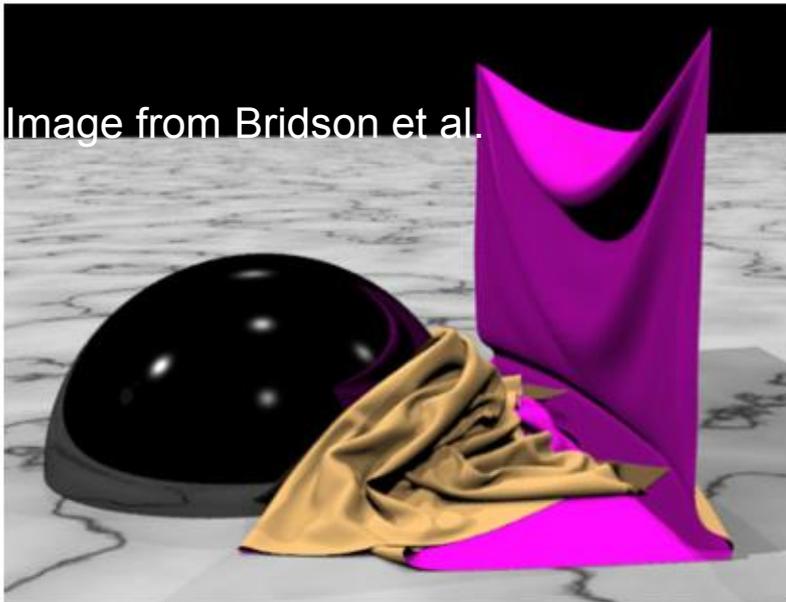
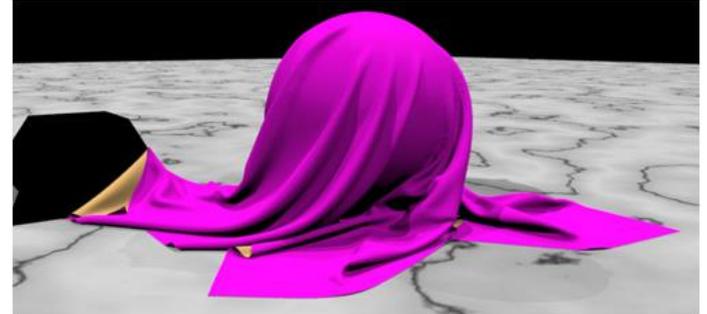
---

- $x' = -kx$  is a true 1st order ODE
- Energy gets dissipated
- In contrast, a spring is a second order system
- Energy does not get dissipated
  - It is just transferred between potential and kinetic energy
  - Unless you add damping
- This is why people always add damping forces and results look too viscous

- 
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  - Energy does not get dissipated
    - It is just transferred between potential and kinetic energy
    - Unless you add damping
  - This is why people always add damping forces and results look too viscous

# The Collision Problem

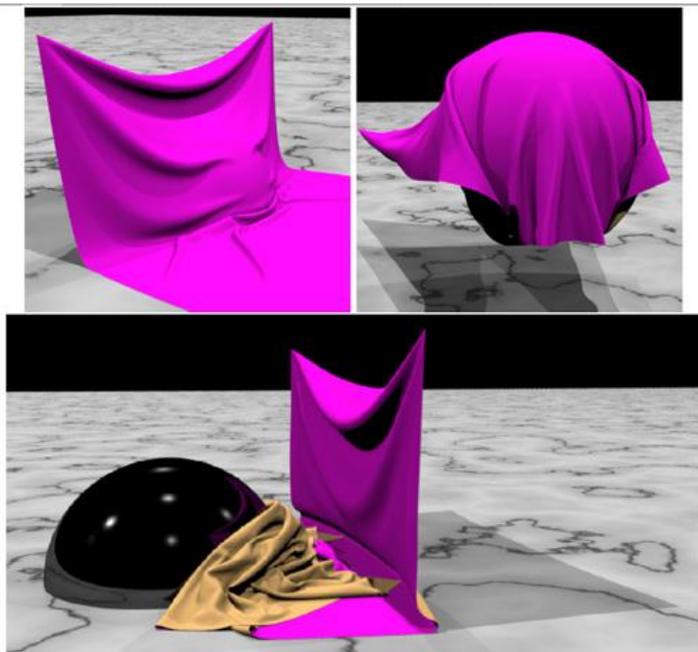
- A cloth has many points of contact
- Requires
  - Efficient collision detection
  - Efficient numerical treatment (stability)



# Collisions

Robert Bridson, Ronald Fedkiw & John Anderson  
*Robust Treatment of Collisions, Contact*  
*and Friction for Cloth Animation*  
SIGGRAPH 2002

- Cloth has many points of contact
- Need efficient collision detection and stable treatment



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# Cool Cloth/Hair Demos

---

- Robert Bridson, Ronald Fedkiw & John Anderson: Robust Treatment of Collisions, Contact and Friction for Cloth Animation SIGGRAPH 2002
- Selle. A, Su, J., Irving, G. and Fedkiw, R., "Robust High-Resolution Cloth Using Parallelism, History-Based Collisions, and Accurate Friction," IEEE TVCG 15, 339-350 (2009).
- Selle, A., Lentine, M. and Fedkiw, R., "A Mass Spring Model for Hair Simulation", SIGGRAPH 2008, ACM TOG 27, 64.1-64.11 (2008).

# Cool Cloth/Hair Demos

---

⇒ U[YfYa cj YX'Xi Y'hc'W&dnf][\hfYghf]V]cbg"

- Selle. A, Su, J., Irving, G. and Fedkiw, R., "Robust High-Resolution Cloth Using Parallelism, History-Based Collisions, and Accurate Friction," IEEE TVCG 15, 339-350 (2009).

=a U[ Y fYa cj YX Xi Y hc Wdrf][ \ hfYgrf]V]cbg"

- Selle. A, Su, J., Irving, G. and Fedkiw, R., "Robust High-Resolution Cloth Using Parallelism, History-Based Collisions, and Accurate Friction," IEEE TVCG 15, 339-350 (2009).

# Implementation Notes

---

- It pays off to abstract (as usual)
  - It's easy to design your “Particle System” and “Time Stepper” to be unaware of each other
- Basic idea
  - “Particle system” and “Time Stepper” communicate via floating-point vectors  $\mathbf{X}$  and a function that computes  $f(\mathbf{X},t)$ 
    - “Time Stepper” does not need to know anything else!

# Implementation Notes

---

- Basic idea
  - “Particle System” tells “Time Stepper” how many dimensions ( $N$ ) the phase space has
  - “Particle System” has a function to write its state to an  $N$ -vector of floating point numbers (and read state from it)
  - “Particle System” has a function that evaluates  $f(\mathbf{X},t)$ , given a state vector  $\mathbf{X}$  and time  $t$
  - “Time Stepper” takes a “Particle System” as input and advances its state

# Particle System Class

---

```
class ParticleSystem
{
    virtual int getDimension()
    virtual setDimension(int n)
    virtual float* getStatePositions()
    virtual setStatePositions(float* positions)
    virtual float* getStateVelocities()
    virtual setStateVelocities(float* velocities)
    virtual float* getForces(float* positions, float* velocities)
    virtual setMasses(float* masses)
    virtual float* getMasses()

    float* m_currentState
}
```

# Time Stepper Class

---

```
class TimeStepper
{
    virtual takeStep(ParticleSystem* ps, float h)
}
```

# Forward Euler Implementation

---

```
class ForwardEuler : TimeStepper
{
    void takeStep(ParticleSystem* ps, float h)
    {
        velocities = ps->getStateVelocities()
        positions = ps->getStatePositions()
        forces = ps->getForces(positions, velocities)
        masses = ps->getMasses()
        accelerations = forces / masses
        newPositions = positions + h*velocities
        newVelocities = velocities + h*accelerations
        ps->setStatePositions(newPositions)
        ps->setStateVelocities(newVelocities)
    }
}
```

# Mid-Point Implementation

---

```
class MidPoint : TimeStepper
{
    void takeStep(ParticleSystem* ps, float h)
    {
        velocities = ps->getStateVelocities()
        positions = ps->getStatePositions()
        forces = ps->getForces(positions, velocities)
        masses = ps->getMasses()
        accelerations = forces / masses
        midPositions = positions + 0.5*h*velocities
        midVelocities = velocities + 0.5*h*accelerations
        midForces = ps->getForces(midPositions, midVelocities)
        midAccelerations = midForces / masses
        newPositions = positions + 0.5*h*midVelocities
        newVelocities = velocities + 0.5*h*midAccelerations
        ps->setStatePositions(newPositions)
        ps->setStateVelocities(newVelocities)
    }
}
```

# Particle System Simulation

---

```
ps = new MassSpringSystem(particleCount, masses, springs, externalForces)
stepper = new ForwardEuler()
time = 0
while time < 1000
    stepper->takeStep(ps, 0.0001)
    time = time + 0.0001
    // render
```

# Particle System Simulation

---

```
ps = new MassSpringSystem(particleCount, masses, springs, externalForces)
stepper = new MidPoint()
time = 0
while time < 1000
    stepper->takeStep(ps, 0.0001)
    time = time + 0.0001
    // render
```

# Questions?

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# That's All for Today!

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6.837 Computer Graphics  
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