

Lecture 8 - Carrier Drift and Diffusion (*cont.*), Carrier Flow

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1. Quasi-Fermi levels
2. Continuity equations
3. Surface continuity equations

Reading assignment:

del Alamo, Ch. 4, §4.6; Ch. 5, §§5.1, 5.2

Key questions

- Is there something equivalent to the Fermi level that can be used outside equilibrium?
- How do carrier distributions *in energy* look like outside equilibrium?
- In the presence of carrier flow in the bulk of a semiconductor, how does one formulate bookkeeping relationships for carriers?
- How about at surfaces?

1. Quasi-Fermi levels

□ Interested in energy band diagram representations of complex situations in semiconductors *outside thermal equilibrium*.

□ In TE, Fermi level makes statement about energy distribution of carriers in bands

⇒ E_F relates n_o with N_c and p_o with N_v :

$$n_o = N_c \mathcal{F}_{1/2}\left(\frac{E_F - E_c}{kT}\right) \quad p_o = N_v \mathcal{F}_{1/2}\left(\frac{E_v - E_F}{kT}\right)$$

Outside TE, E_F cannot be used.

Define two "quasi-Fermi levels" such that:

$$n = N_c \mathcal{F}_{1/2}\left(\frac{E_{fe} - E_c}{kT}\right) \quad p = N_v \mathcal{F}_{1/2}\left(\frac{E_v - E_{fh}}{kT}\right)$$

Under Maxwell-Boltzmann statistics ($n \ll N_c$, $p \ll N_v$):

$$n = N_c \exp\left(\frac{E_{fe} - E_c}{kT}\right)$$

$$p = N_v \exp\left(\frac{E_v - E_{fh}}{kT}\right)$$

What are quasi-Fermi levels good for?

□ Take derivative of $n = f(E_{fe})$ with respect to x :

$$n = N_c \exp \frac{E_{fe} - E_c}{kT}$$

$$\frac{dn}{dx} = \frac{n}{kT} \left(\frac{dE_{fe}}{dx} - \frac{dE_c}{dx} \right) = \frac{n}{kT} \frac{dE_{fe}}{dx} - \frac{q}{kT} n \mathcal{E}$$

Plug into current equation:

$$J_e = q\mu_e n \mathcal{E} + qD_e \frac{dn}{dx}$$

To get:

$$J_e = \mu_e n \frac{dE_{fe}}{dx}$$

Similarly for holes:

$$J_h = \mu_h p \frac{dE_{fh}}{dx}$$

Gradient of quasi-Fermi level: *unifying driving force for carrier flow.*

□ Physical meaning of ∇E_f

For electrons,

$$J_e = \mu_e n \frac{dE_{fe}}{dx} = -qnv_e$$

Then:

$$\frac{dE_{fe}}{dx} = -\frac{q}{\mu_e} v_e$$

∇E_{fe} linearly proportional to electron velocity!

Similarly for holes:

$$\frac{dE_{fh}}{dx} = \frac{q}{\mu_h} v_h$$

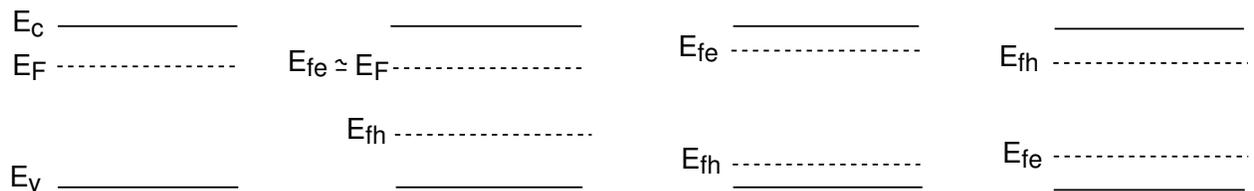
□ Quasi-Fermi levels: effective way to *visualize* carrier phenomena outside equilibrium in energy band diagram

1. Visualize carrier concentrations and net recombination

$$np = n_i^2 \exp \frac{E_{fe} - E_{fh}}{kT}$$

- If $E_{fe} > E_{fh} \Rightarrow np > n_i^2 \Rightarrow U > 0$
- If $E_{fe} < E_{fh} \Rightarrow np < n_i^2 \Rightarrow U < 0$
- If $E_{fe} = E_{fh} \Rightarrow np = n_i^2 \Rightarrow U = 0$ (carrier conc's in TE)

Examples (same semiconductor):



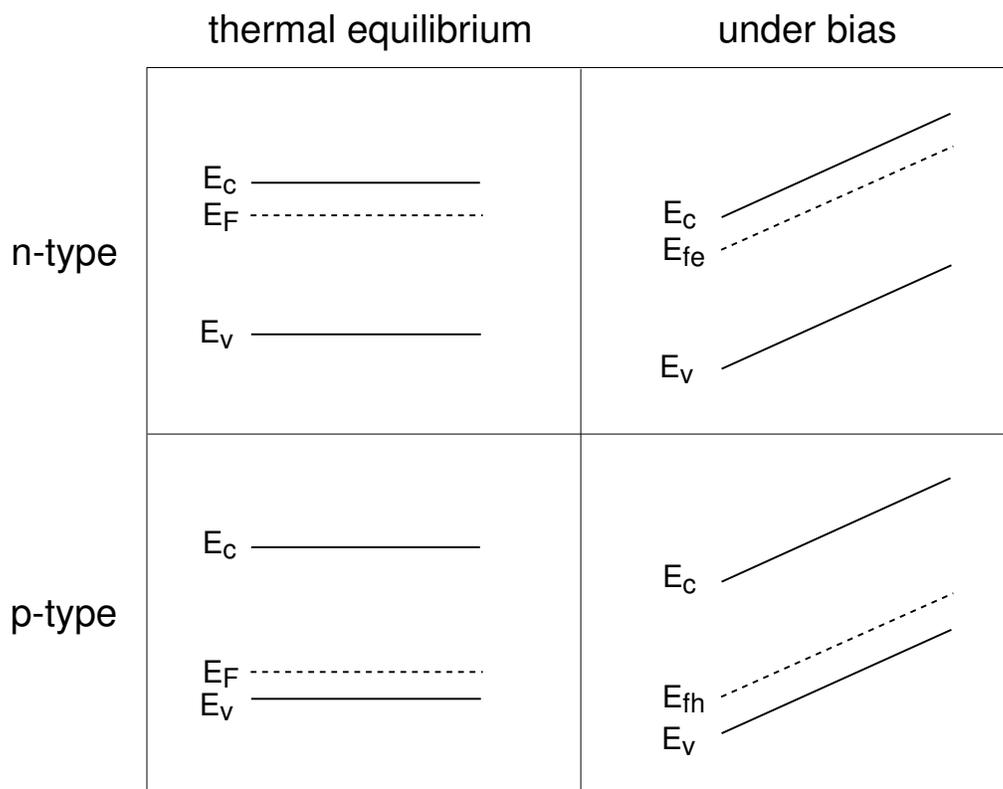
But can't visualize G_{ext} .

2. Visualize currents

$$J_e = \mu_e n \frac{dE_{fe}}{dx}$$

- $\nabla E_{fe} = 0 \Rightarrow J_e = 0$
- $\nabla E_{fe} \neq 0 \Rightarrow J_e \neq 0$
- if n high, ∇E_{fe} small to maintain a certain current level
- if n low, ∇E_{fe} large to maintain a certain current level

Examples:



□ The concept of Quasi-Fermi level hinges on notion of:

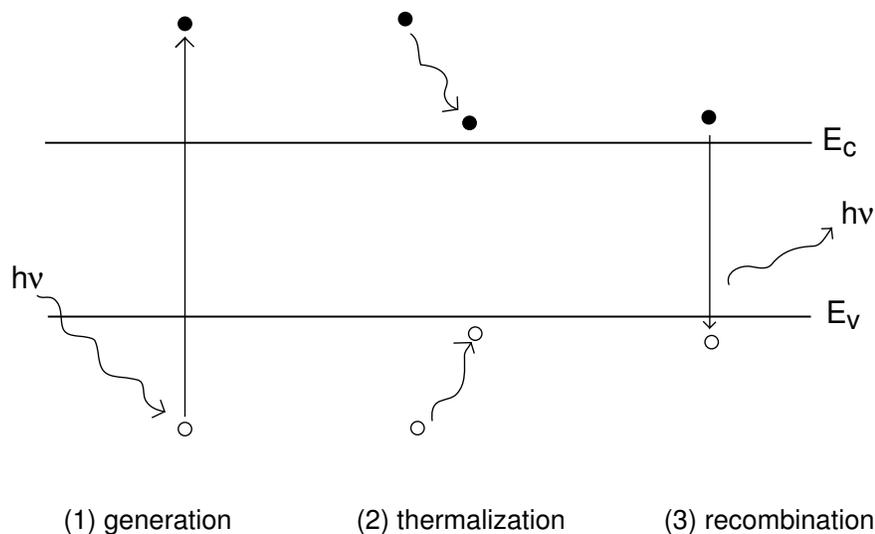
Quasi-equilibrium: *carrier distributions in energy never depart too far from TE in times scales of practical interest.*

Quasi-equilibrium appropriate if:

$$\text{scattering time} \ll \text{dominant device time constant}$$

⇒ carriers undergo many collisions and attain thermal quasi-equilibrium *with the lattice and among themselves* very quickly.

In time scales of interest, carrier distribution is close to *Maxwellian* (i.e., well described by a Fermi level).



2. Continuity Equations

Semiconductor physics so far:

Gauss' law:
$$\vec{\nabla} \cdot \vec{\mathcal{E}} = \frac{q}{\epsilon}(p - n + N_D^+ - N_A^-)$$

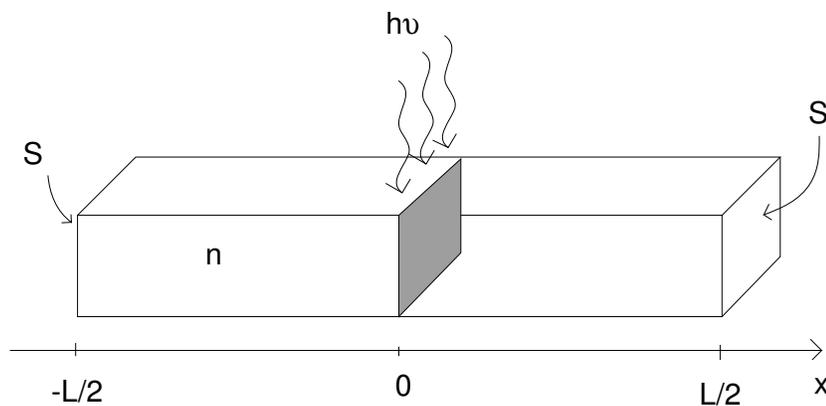
Electron current equation:
$$\vec{J}_e = -qn\vec{v}_e^{drift} + qD_e\vec{\nabla}n$$

Hole current equation:
$$\vec{J}_h = qp\vec{v}_h^{drift} - qD_h\vec{\nabla}p$$

Total current equation:
$$\vec{J}_t = \vec{J}_e + \vec{J}_h$$

Carrier dynamics:
$$\frac{dn}{dt} = \frac{dp}{dt} = G - R$$

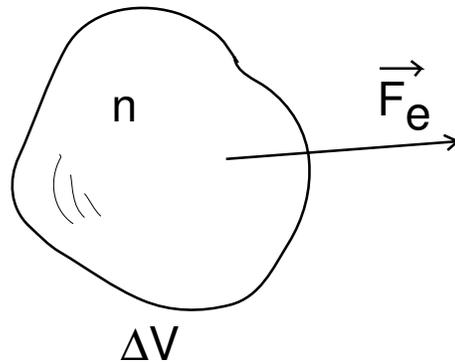
Still, can't solve problems like this:



Equation system does not capture:

- impact of carrier movement on carrier concentration (*i.e.* when carriers move away from a point, their concentration drops!)
- boundary conditions (surfaces are not infinitely far away)

Need "book-keeping relationships" for particles:



rate of increase of number of electrons in $\Delta V =$
 rate of electron generation in ΔV
 - rate of electron recombination in ΔV
 - net flow of electrons leaving ΔV per unit time

$$\frac{\partial(n\Delta V)}{\partial t} = G\Delta V - R\Delta V - \int \vec{F}_e \cdot d\vec{S}$$

Dividing by ΔV and taking the limit of small ΔV :

$$\frac{\partial n}{\partial t} = G - R - \vec{\nabla} \cdot \vec{F}_e$$

In terms of current density:

$$\frac{\partial n}{\partial t} = G - R + \frac{1}{q} \vec{\nabla} \cdot \vec{J}_e$$

For holes:

$$\frac{\partial p}{\partial t} = G - R - \frac{1}{q} \vec{\nabla} \cdot \vec{J}_h$$

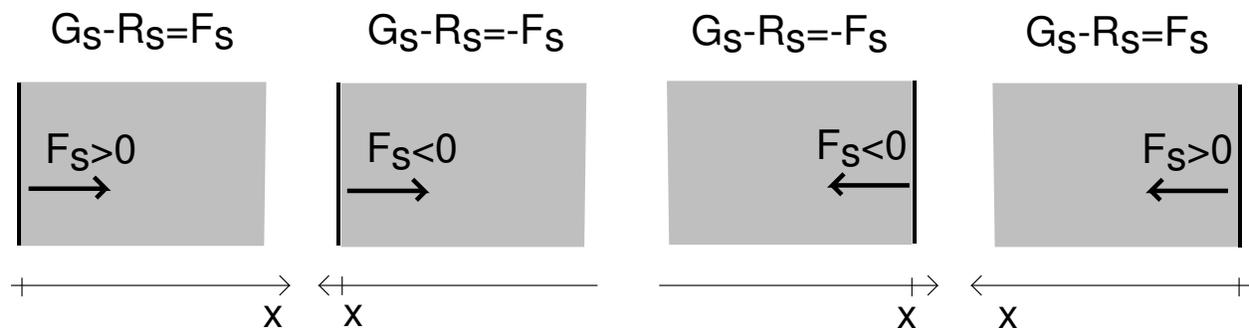
3. Surface continuity equations

□ "Free" surface: cannot "store" carriers:

*Surface generation - surface recombination =
carrier flow out of surface*

$$|G_s - R_s| = |F_s|$$

This equation is axis sensitive.



Rewrite in terms of current densities normal to surface:

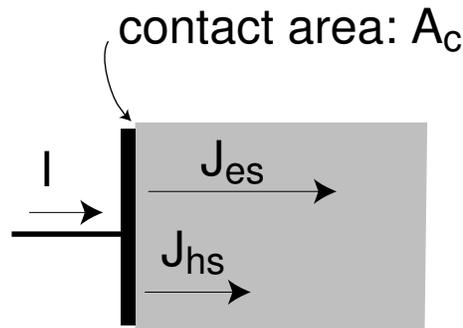
$$|U_s| = \frac{1}{q} |J_{es}| = \frac{1}{q} |J_{hs}|$$

Always, no net current into surface:

$$J_s = J_{es} + J_{hs} = 0 \implies J_{es} = -J_{hs}$$

□ *Ohmic contact*: provides path for current to flow out of device.

For n-type:



Kirchoff's law demands current continuity at metal-semiconductor interface:

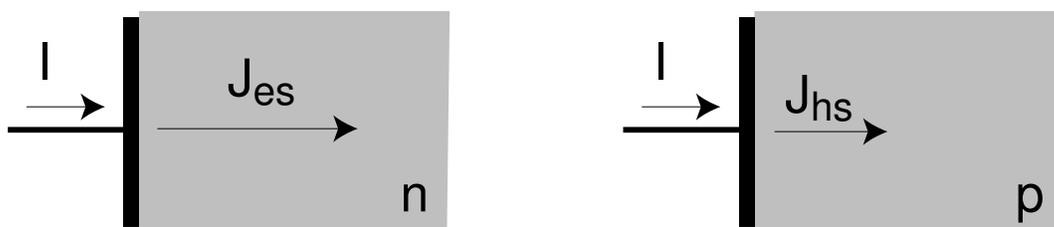
$$|I| = A_c |J_{es} + J_{hs}| = qA_c |F_{es} - F_{hs}|$$

Equation is sign sensitive:

- IEEE convention: I entering into device is positive
- sign of J_{es} and J_{hs} depend of choice of axis in semiconductor

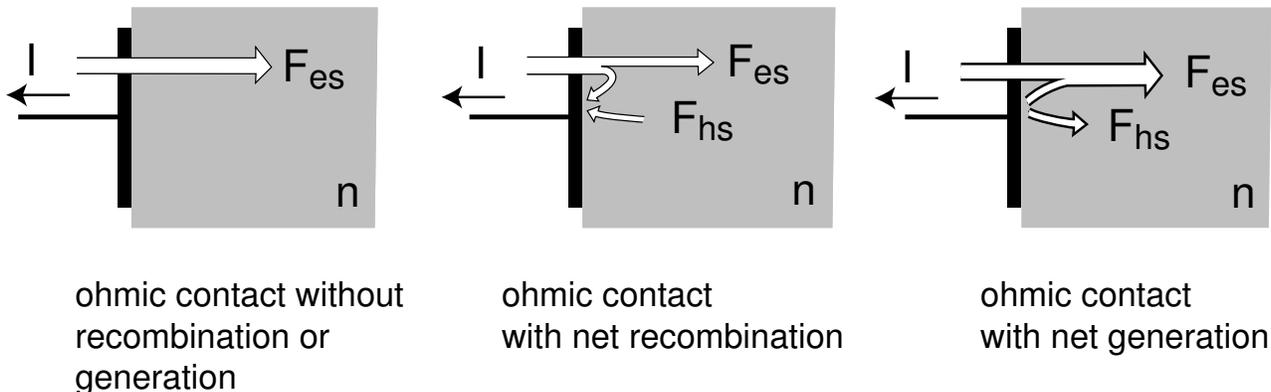
Key result from metal-semiconductor junction: Metal in ohmic contact only communicates with majority carriers in semiconductor \Rightarrow *majority carrier current*

- in n-type material, current supported by electrons
- in p-type material, current supported by holes



□ If there is net generation or recombination at surface right below ohmic contact \Rightarrow *minority carrier current* in addition to majority carrier current

Three possible cases (examples for n-type):



- ohmic contact with equilibrium carrier concentrations:

$$U_s = 0 \quad |I| = qA_c |F_{es}|$$

- ohmic contact with net recombination:

$$U_s = |F_{hs}| \quad |I| = qA_c |F_{es} - F_{hs}| > qA_c |F_{es}|$$

- ohmic contact with net generation:

$$U_s = -|F_{hs}| \quad |I| = qA_c |F_{es} - F_{hs}| < qA_c |F_{es}|$$

Important boundary condition (will understand soon): at metal-semiconductor interface:

$$n'_s = p'_s = 0 \quad \text{or} \quad S = \infty$$

Key conclusions

- *Quasi-Fermi levels* describe carrier statistics outside equilibrium \Rightarrow powerful way to visualize carrier concentrations and currents in energy band diagrams.
- *Quasi-equilibrium*: carrier distributions in energy not significantly different from TE in time scales of interest for device operation.
- *Continuity equations*: "book-keeping" relations for carriers.
- Surfaces cannot store carriers: at all times must have current balance at surface.
- At "free" surface: electron and hole currents result from carrier generation or recombination at surface (but net current is zero).
- At ohmic contact:
 - additional majority carrier current to support terminal current
 - excess carrier concentrations are zero

Self-study

- Do examples 4.8, 4.9, 5.1, 5.2.
- Study integral form of continuity equations and corollary (§5.1).