

**Lecture 13**  
**The Levinson-Durbin Recursion**

In the previous lecture we looked at all-pole signal modeling, linear prediction, and the stochastic inverse-whitening problem. Each scenario was related in concept to the problem of processing a signal  $s[n]$  by:

$$s[n] \longrightarrow \boxed{\frac{1}{A} [1 - \sum_{k=1}^p a_k z^{-k}]} \longrightarrow g[n],$$

such that

$$e[n] \equiv g[n] - \delta[n]$$

was minimized in some sense. In the deterministic case, we chose to minimize  $\varepsilon = \sum_n e^2[n]$ , and in the stochastic case, we chose to minimize  $\varepsilon = E [e^2[n]]$ . The causal solution for the parameters  $a_k$  was found in each case to be of the form

$$\begin{bmatrix} \phi_s[0] & \phi_s[1] & \cdots & \phi_s[p-1] \\ \phi_s[1] & \phi_s[0] & \cdots & \phi_s[p-2] \\ \vdots & \vdots & \ddots & \vdots \\ \phi_s[p-1] & \phi_s[p-2] & \cdots & \phi_s[0] \end{bmatrix} \begin{bmatrix} a_1 \\ a_2 \\ \vdots \\ a_p \end{bmatrix} = \begin{bmatrix} \phi_s[1] \\ \phi_s[2] \\ \vdots \\ \phi_s[p] \end{bmatrix}, \quad (1)$$

or equivalently,

$$T_p \alpha_p = r_p. \quad (2)$$

In this lecture, we address the issue of how exactly to solve for the parameters  $a_k$  (or equivalently for the vector  $\alpha_p$ ). The fact that  $\alpha_p$  can be found simply by calculating  $\alpha_p = T_p^{-1} r_p$  is worth a mention, but the focus of this lecture is on an elegant technique for finding  $\alpha_p$  which is generally less computationally expensive than taking a matrix inversion and which is also recursive in the filter order  $p$ . This method for finding  $\alpha_p$  is called the *Levinson-Durbin* recursion. In addition to being recursive in  $p$ , we'll see that certain intermediate results from the recursion give coefficients for a lattice implementation of an LTI filter  $h[n]$  which implements our signal model as

$$H(z) = S'(z) = \frac{A}{1 - \sum_{k=1}^p a_k z^{-k}}.$$

We'll also see that the Levinson-Durbin recursion prescribes a similar implementation for the inverse filter  $1/H(z)$ .

Let's first prove the Levinson-Durbin recursion. The basic idea of the recursion is to find the solution  $\alpha_{p+1}$  for the  $(p+1)$ st order case from the solution  $\alpha_p$  for the  $p$ th order case. We'll see that this hinges on the fact that  $T_p$  is a  $p \times p$  Toeplitz matrix, that is, that  $T_p$  is symmetric, and that all entries along a given diagonal are equal. This property of  $T_p$  is a result of its definition:

$$(T_p)_{ij} \equiv \phi_s[i - j].$$

To begin the proof, consider what happens in the  $p = 1$  case. Equation 2 becomes for  $p = 1$

$$T_1 \alpha_1 = r_1,$$

or

$$\phi_s[0] a_1 = \phi_s[1],$$

giving the solution

$$a_1 = \frac{\phi_s[1]}{\phi_s[0]}. \quad (3)$$

The recursion will now be developed by evaluating Equation 2 for order  $p + 1$ :

$$T_{p+1} \alpha_{p+1} = r_{p+1}.$$

Because we will be dealing with a solution vector  $\alpha_\ell$  for multiple orders  $\ell$ , we will adopt the notation  $a_k^{(\ell)}$  to refer to the parameter  $a_k$  for the  $\ell$ th-order model. (Note that  $a_k^{(\ell)}$  is generally not equal to  $a_k^{(m)}$ .) Equation 2 is therefore represented in matrix form as

$$\begin{bmatrix} \phi_s[0] & \phi_s[1] & \cdots & \phi_s[p-1] & \phi_s[p] \\ \phi_s[1] & \phi_s[0] & \cdots & \phi_s[p-2] & \phi_s[p-1] \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ \phi_s[p-1] & \phi_s[p-2] & \cdots & \phi_s[0] & \phi_s[1] \\ \phi_s[p] & \phi_s[p-1] & \cdots & \phi_s[1] & \phi_s[0] \end{bmatrix} \begin{bmatrix} a_1^{(p+1)} \\ a_2^{(p+1)} \\ \vdots \\ a_p^{(p+1)} \\ a_{p+1}^{(p+1)} \end{bmatrix} = \begin{bmatrix} \phi_s[1] \\ \phi_s[2] \\ \vdots \\ \phi_s[p] \\ \phi_s[p+1] \end{bmatrix}.$$

The matrix  $T_{p+1}$  relates to  $T_p$  as

$$T_{p+1} = \left[ \begin{array}{c|c} & \begin{matrix} \phi_s[p] \\ \phi_s[p-1] \\ \vdots \\ \phi_s[1] \end{matrix} \\ \hline \begin{matrix} \phi_s[p] & \phi_s[p-1] & \cdots & \phi_s[1] \end{matrix} & \begin{matrix} \phi_s[0] \end{matrix} \end{array} \right],$$

and the vector  $r_{p+1}$  relates to  $r_p$  as

$$r_{p+1} = \begin{bmatrix} r_p \\ \phi_s[p+1] \end{bmatrix}.$$

Defining a new vector  $\rho_p$  as  $r_p$  upside-down, or

$$\rho_p = \begin{bmatrix} \phi_s[p] \\ \phi_s[p-1] \\ \vdots \\ \phi_s[1] \end{bmatrix},$$

gives a more compact representation of  $T_{p+1}$  in terms of  $T_p$ :

$$T_{p+1} = \left[ \begin{array}{c|c} T_p & \rho_p \\ \hline (\rho_p)^T & \phi_s[0] \end{array} \right].$$

Let's now represent the  $(p+1)$ st-order parameter vector  $\alpha_{p+1}$  in terms of the  $p$ th order vector  $\alpha_p$ , a correction term  $k_{p+1}$  and a correction vector  $\varepsilon_p$  as

$$\alpha_{p+1} = \boxed{\begin{bmatrix} \alpha_p \\ 0 \end{bmatrix} + \begin{bmatrix} \varepsilon_p \\ k_{p+1} \end{bmatrix}}. \quad (4)$$

Equation 2 for order  $p+1$  is therefore represented in terms of the  $p$ th order equation as

$$\left[ \begin{array}{c|c} T_p & \rho_p \\ \hline (\rho_p)^T & \phi_s[0] \end{array} \right] \left\{ \begin{bmatrix} \alpha_p \\ 0 \end{bmatrix} + \begin{bmatrix} \varepsilon_p \\ k_{p+1} \end{bmatrix} \right\} = \begin{bmatrix} r_p \\ \phi_s[p+1] \end{bmatrix},$$

which, sorting the equations out, implies

$$T_p \alpha_p + T_p \varepsilon_p + \rho_p k_{p+1} = r_p \quad (5)$$

and

$$(\rho_p)^T \alpha_p + (\rho_p)^T \varepsilon_p + \phi_s[0] k_{p+1} = \phi_s[p+1]. \quad (6)$$

Substituting Equation 2 for order  $p=1$  into Equation 5 then gives

$$T_p \varepsilon_p k_{p+1}^{-1} = -\rho_p. \quad (7)$$

Note now that since  $T_p$  is Toeplitz, the matrix realization of the causal Yule-Walker equations for order  $p$  (Equation 1) implies also

$$\begin{bmatrix} \phi_s[0] & \phi_s[1] & \cdots & \phi_s[p-1] \\ \phi_s[1] & \phi_s[0] & \cdots & \phi_s[p-2] \\ \vdots & \vdots & \ddots & \vdots \\ \phi_s[p-1] & \phi_s[p-2] & \cdots & \phi_s[0] \end{bmatrix} \begin{bmatrix} a_p^{(p)} \\ a_{p-1}^{(p)} \\ \vdots \\ a_1^{(p)} \end{bmatrix} = \begin{bmatrix} \phi_s[p] \\ \phi_s[p-1] \\ \vdots \\ \phi_s[1] \end{bmatrix},$$

or

$$T_p \beta_p = \rho_p,$$

where  $\beta_p$  is  $\alpha_p$  upside-down. This then implies, after substituting in Equation 5,

$$\boxed{\varepsilon_p = -k_{p+1}\beta_p}. \quad (8)$$

We'll now use this and Equation 6 to find  $k_{p+1}$  and thus have all the pieces needed to complete the recursion.

Pre-multiplying both sides of Equation 8 by  $(\rho_p)^T$  gives the scalar relation

$$(\rho_p)^T \varepsilon_p = -(\rho_p)^T \beta_p k_{p+1}.$$

Note, however, that the scalar relation  $(\rho_p)^T \beta_p = (r_p)^T \alpha_p$  also holds, since  $r_p$  is a re-ordered version of  $\rho_p$  in the same way that  $\alpha_p$  is a re-ordered version of  $\beta_p$ . This allows us to re-write the above equation as

$$(\rho_p)^T \varepsilon_p = -(r_p)^T \alpha_p k_{p+1},$$

which implies from Equation 6

$$(\rho_p)^T \alpha_p - (r_p)^T \alpha_p k_{p+1} + \phi_s[0]k_{p+1} = \phi_s[p+1],$$

or

$$\boxed{k_{p+1} = \frac{\phi_s[p+1] - (\rho_p)^T \alpha_p}{\phi_s[0] - (r_p)^T \alpha_p}}. \quad (9)$$

Equations 4, 8, and 9 therefore define a recursion formula for finding the  $(p+1)$ st-order model parameters  $a_k^{(p+1)}$  in terms of the previously-obtained  $p$ th-order solution. Equation 3 gives a closed-form solution for  $a_1^{(1)}$ , which provides a starting point for the recursion. Since this ends a considerable algebraic detour, we'll now summarize what was concluded.

## The Levinson-Durbin recursion

- Problem statement: for order  $p$ , solve  $\sum_{k=1}^p a_k^{(p)} \phi_s[i-k] = \phi_s[i]$
- Definitions:

$$\begin{aligned}\alpha_p &= [a_1^{(p)}, a_2^{(p)}, \dots, a_p^{(p)}]^T \\ \beta_p &= [a_p^{(p)}, a_{p-1}^{(p)}, \dots, a_1^{(p)}]^T \\ r_p &= [\phi_s[1], \phi_s[2], \dots, \phi_s[p]]^T \\ \rho_p &= [\phi_s[p], \phi_s[p-1], \dots, \phi_s[1]]^T\end{aligned}$$

- Solution for  $p = 1$ :  $a_1^{(1)} = \frac{\phi_s[1]}{\phi_s[0]}$
- Recursion:

$$\begin{aligned}k_{p+1} &= \frac{\phi_s[p+1] - (\rho_p)^T \alpha_p}{\phi_s[0] - (r_p)^T \alpha_p} \\ \varepsilon_p &= -k_{p+1} \beta_p \\ \alpha_{p+1} &= \begin{bmatrix} \alpha_p \\ 0 \end{bmatrix} + \begin{bmatrix} \varepsilon_p \\ k_{p+1} \end{bmatrix} = \begin{bmatrix} a_1^{(p)} \\ a_2^{(p)} \\ \vdots \\ a_p^{(p)} \\ 0 \end{bmatrix} - k_{p+1} \begin{bmatrix} a_p^{(p)} \\ a_{p-1}^{(p)} \\ \vdots \\ a_1^{(p)} \\ -1 \end{bmatrix}\end{aligned}$$

In parting, note that when these coefficients  $k_\ell$  are used to implement an all-pole lattice filter with reflection coefficients  $k_\ell$ , the lattice filter's response is described by

$$H(z) = S'(z) = \frac{A}{1 - \sum_{m=1}^p a_m z^{-m}}.$$

In other words, the impulse response of an all-pole lattice filter with reflection coefficients  $k_\ell$ ,  $\ell = 1, \dots, p$ , as prescribed by the Levinson-Durbin recursion, is the  $p$ th-order model of our original signal  $s[n]$ . (Likewise, the corresponding all-zero lattice filter implements  $1/S'(z)$ .) Note further that it is straightforward to determine the stability of an all-pole filter when implemented in lattice form. Specifically, if the coefficients  $k_\ell$  all have magnitude  $< 1$ , the all-pole filter is guaranteed to be stable. (It is left as an exercise to gain intuition for why this is so.) This stability criterion is of particular interest in applications where the all-pole filter is

implemented with finite-precision coefficients, since determining the stability of a comparable direct-form all-pole implementation requires the generally more-expensive process of determining all of the filter's pole locations. In addition to the computational benefits gained from the Levinson-Durbin recursion, its connection with the lattice structure (and the structure's associated stability metric) enables a wide array of applications, including real-time speech coders and synthesizers.