

Introduction to Simulation - Lecture 19

Laplace's Equation – FEM Methods

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Thanks to Deepak Ramaswamy, Michal Rewienski,
and Karen Veroy, Jaime Peraire and Tony Patera

Outline for Poisson Equation Section

- Why Study Poisson's equation
 - Heat Flow, Potential Flow, Electrostatics
 - Raises many issues common to solving PDEs.
- Basic Numerical Techniques
 - basis functions (FEM) and finite-differences
 - Integral equation methods
- Fast Methods for 3-D
 - Preconditioners for FEM and Finite-differences
 - Fast multipole techniques for integral equations

Outline for Today

- Why Poisson Equation
 - Reminder about heat conducting bar
- Finite-Difference And Basis function methods
 - Key question of convergence
- Convergence of Finite-Element methods
 - Key idea: solve Poisson by minimization
 - Demonstrate optimality in a carefully chosen norm

Drag Force Analysis of Aircraft

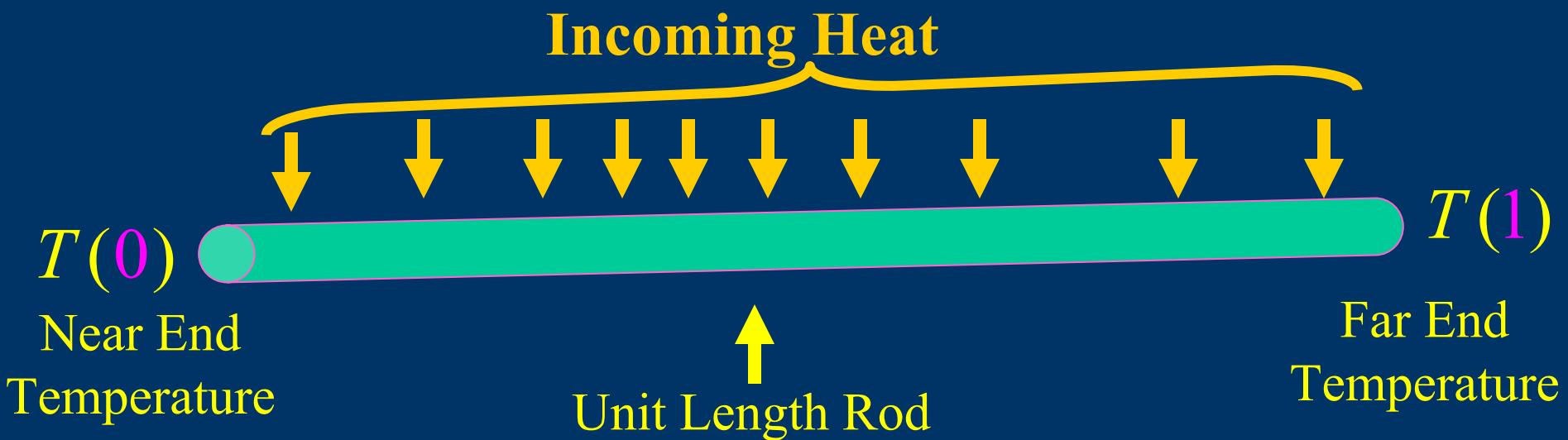
- Potential Flow Equations
 - Poisson Partial Differential Equations.

Engine Thermal Analysis

- Thermal Conduction Equations
 - The Poisson Partial Differential Equation.

Capacitance on a microprocessor Signal Line

- Electrostatic Analysis
 - The Laplace Partial Differential Equation.

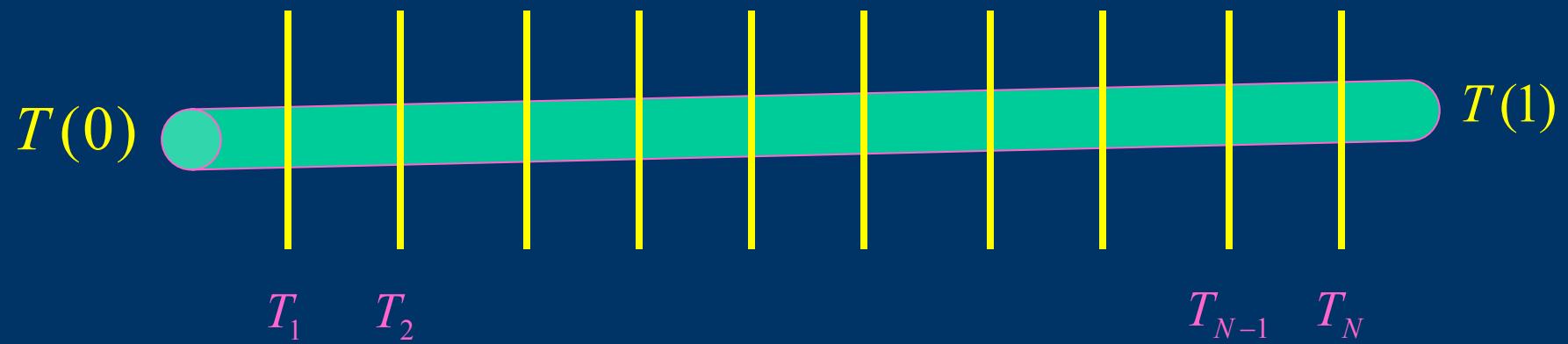


Question: What is the temperature distribution along the bar



1) Cut the bar into short sections

2) Assign each cut a temperature



Heat Flow through one section

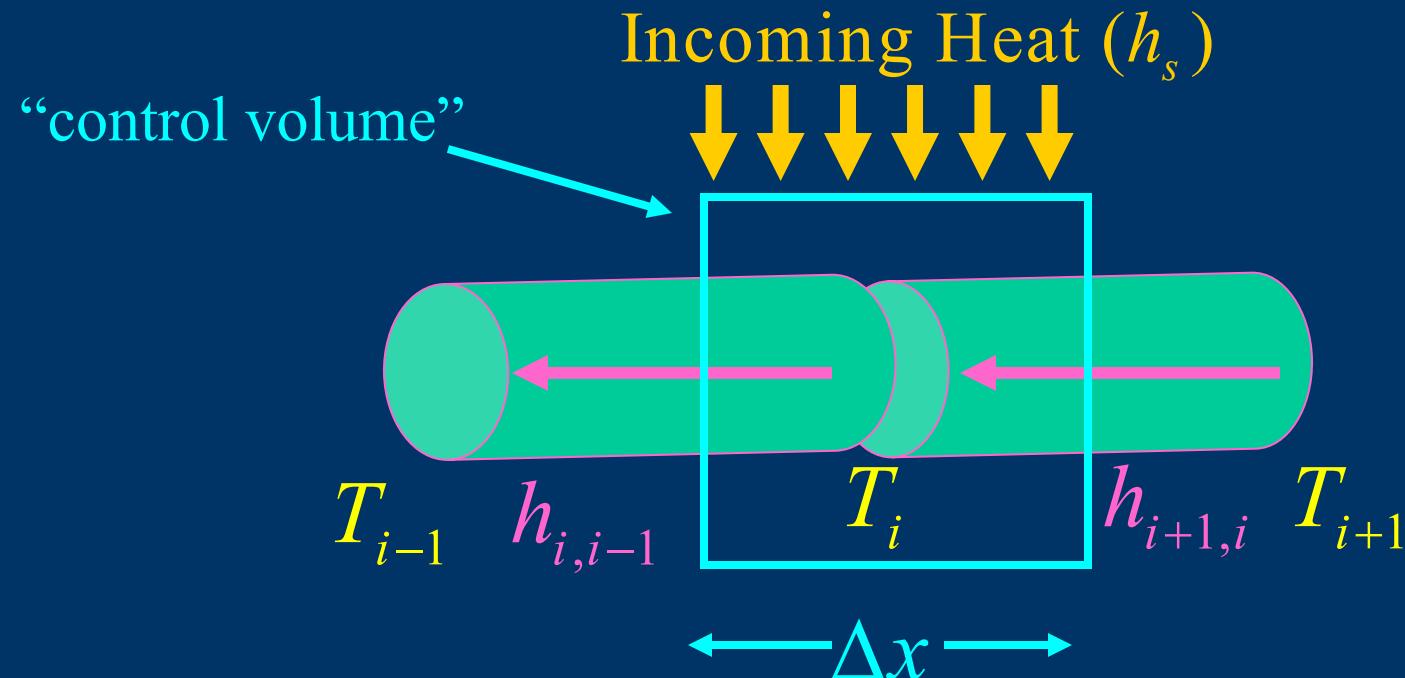

$$T_i \quad T_{i+1} \quad h_{i+1,i} = \text{heat flow} = \kappa \frac{T_{i+1} - T_i}{\Delta x}$$

A diagram showing a cylindrical section of width Δx and height $h_{i+1,i}$. The left face is at temperature T_i and the right face is at temperature T_{i+1} .

Limit as the sections become vanishingly small

$$\lim_{\Delta x \rightarrow 0} h(x) = \kappa \frac{\partial T(x)}{\partial x}$$

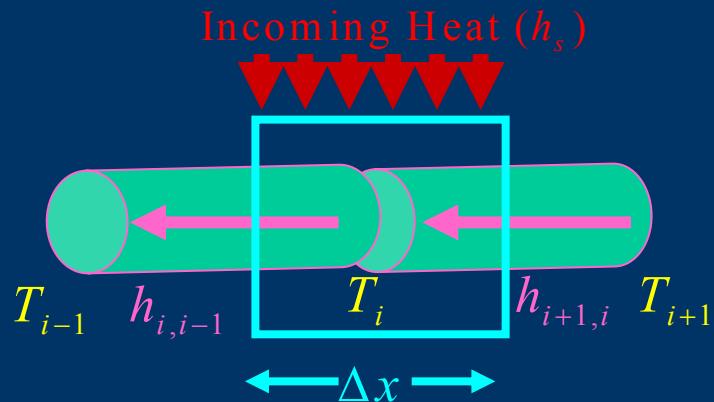
Two Adjacent Sections



Heat Flows into Control Volume Sums to zero

$$h_{i+1,i} - h_{i,i-1} = -h_s \Delta x$$

Heat Flows into Control Volume Sums to zero



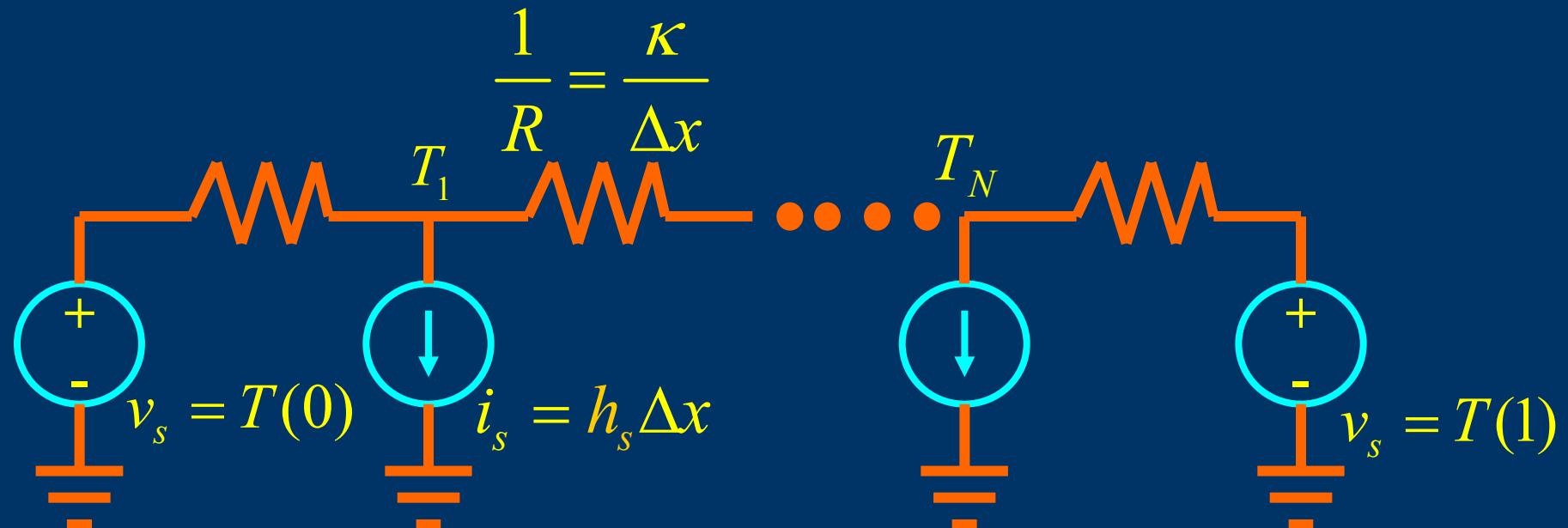
$$h_{i+1,i} - h_{i,i-1} = -h_s \Delta x$$

Heat in from left Heat out from right Incoming heat per unit length

Limit as the sections become vanishingly small

$$\lim_{\Delta x \rightarrow 0} h_s(x) = \frac{\partial h(x)}{\partial x} = \frac{\partial}{\partial x} K \frac{\partial T(x)}{\partial x}$$

Temperature analogous to Voltage
Heat Flow analogous to Current



Normalized Poisson Equation

$$\frac{\partial}{\partial x} \kappa \frac{\partial T(x)}{\partial x} = -h_s \Rightarrow -\frac{\partial^2 u(x)}{\partial x^2} = f(x)$$

$$-u_{xx}(x) = f(x)$$

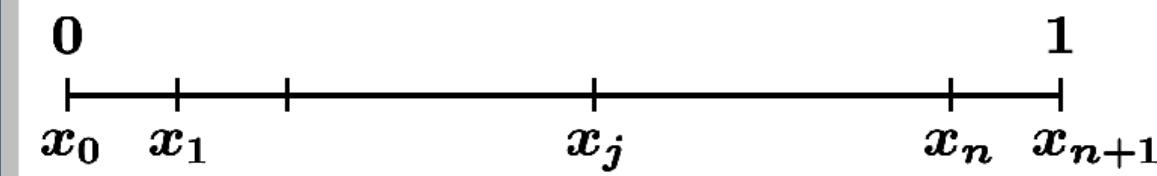
Numerical Solution

Finite Differences

Discretization

Subdivide interval $(0, 1)$ into $n + 1$ equal subintervals

$$\Delta x = \frac{1}{n+1}$$



$$x_j = j\Delta x, \quad \hat{u}_j \approx u_j \equiv u(x_j)$$

$$\text{for } 0 \leq j \leq n + 1$$

Numerical Solution

Finite Differences

Approximation

For example . . .

$$\begin{aligned} v''(x_j) &\approx \frac{1}{\Delta x} (v'(x_{j+1/2}) - v'(x_{j-1/2})) \\ &\approx \frac{1}{\Delta x} \left(\frac{v_{j+1} - v_j}{\Delta x} - \frac{v_j - v_{j-1}}{\Delta x} \right) \\ &= \frac{v_{j+1} - 2v_j + v_{j-1}}{\Delta x^2} \end{aligned}$$

for Δx small

Using Basis Functions

Residual Equation

Partial Differential Equation form

$$-\frac{\partial^2 u}{\partial x^2} = f \quad u(0) = 0 \quad u(1) = 0$$

Basis Function Representation

$$u(x) \approx u_h(x) = \sum_{i=1}^n \omega_i \underbrace{\varphi_i(x)}_{\text{Basis Functions}}$$

Plug Basis Function Representation into the Equation

$$R(x) = \sum_{i=1}^n \omega_i \frac{d^2 \varphi_i(x)}{dx^2} + f(x)$$

Using Basis Functions

Example Basis functions

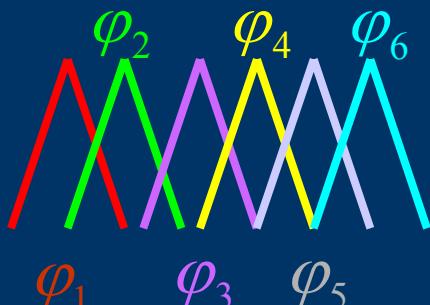
Introduce basis representation $u(x) \approx u_h(x) = \sum_{i=1}^n \omega_i \underbrace{\varphi_i(x)}_{\text{Basis Functions}}$
 $\Rightarrow u_h(x)$ is a weighted sum of basis functions

The basis functions define a space

$$X_h = \left\{ v \in X_h \mid v = \sum_{i=1}^n \beta_i \varphi_i \text{ for some } \beta_i \text{'s} \right\}$$

Example

“Hat” basis functions



Piecewise linear Space



Using Basis functions

Basis Weights

Galerkin Scheme

Force the residual to be “orthogonal” to the basis functions

$$\int_0^1 \varphi_l(x) R(x) dt = 0$$

Generates n equations in n unknowns

$$\int_0^1 \varphi_l(x) \left[\sum_{i=1}^n \omega_i \frac{d^2 \varphi_i(x)}{dx^2} + f(x) \right] dx = 0 \quad l \in \{1, \dots, n\}$$

Using Basis Functions

Basis Weights

Galerkin with integration by parts

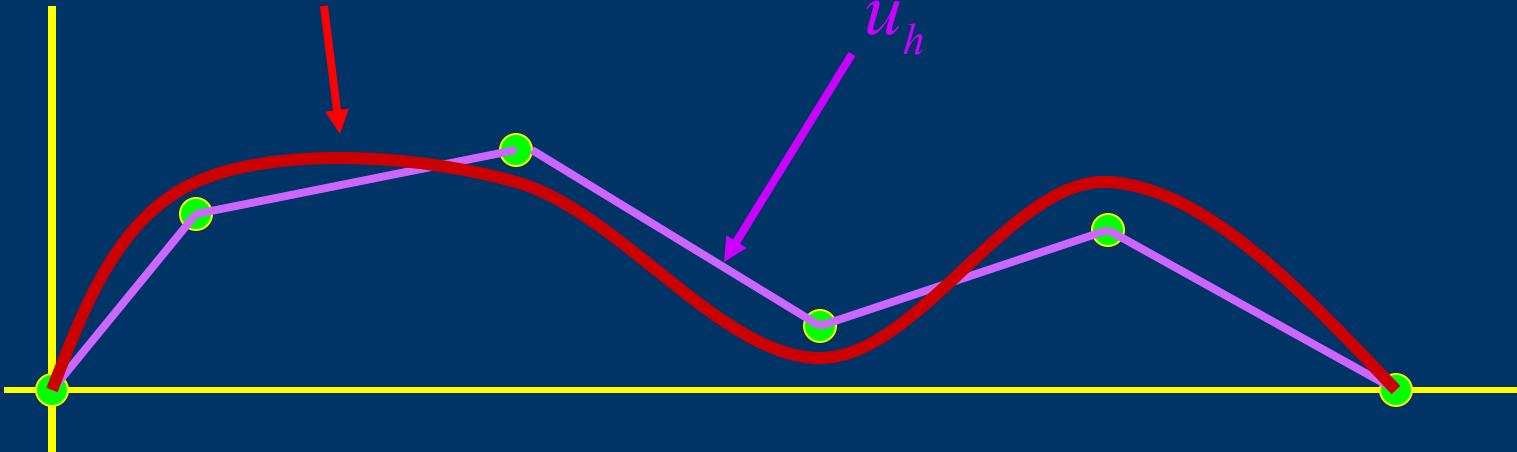
Only first derivatives of basis functions

$$\int_0^1 \frac{d\varphi_l(x)}{dx} \frac{d \sum_{i=1}^n \omega_i \varphi_i(x)}{dx} dx - \int_0^1 \varphi_l(x) f(x) dx = 0$$

$$l \in \{1, \dots, n\}$$

Convergence Analysis

The question is



How does $\underbrace{\|u - u_h\|}_{error}$ decrease with refinement?

- This time – Finite-element methods
- Next time – Finite-difference methods

Partial Differential Equation form

$$-\frac{\partial^2 u}{\partial x^2} = f \quad u(0) = 0 \quad u(1) = 0$$

“Nearly” Equivalent weak form

$$\underbrace{\int_{\Omega} \frac{\partial u}{\partial x} \frac{\partial v}{\partial x} dx}_{a(u,v)} = \underbrace{\int_{\Omega} f v dx}_{l(v)} \quad \text{for all } v$$

Introduced an abstract notation for the equation u must satisfy

$$a(u, v) = l(v) \quad \text{for all } v$$

Introduce basis representation $u(x) \approx u_h(x) = \sum_{i=1}^n \omega_i \underbrace{\varphi_i(x)}_{\text{Basis Functions}}$

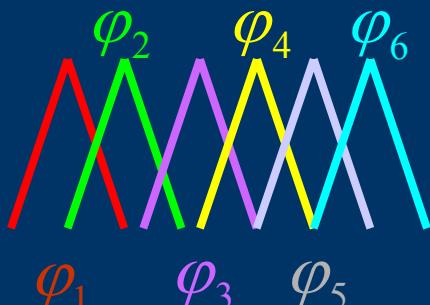
$\Rightarrow u_h(x)$ is a weighted sum of basis functions

The basis functions define a space

$$X_h = \left\{ v \in X_h \mid v = \sum_{i=1}^n \beta_i \varphi_i \text{ for some } \beta_i \text{'s} \right\}$$

Example

“Hat” basis functions



Piecewise linear Space



Key Idea

$a(u, u)$ defines a norm $a(u, u) \equiv \|u\|$

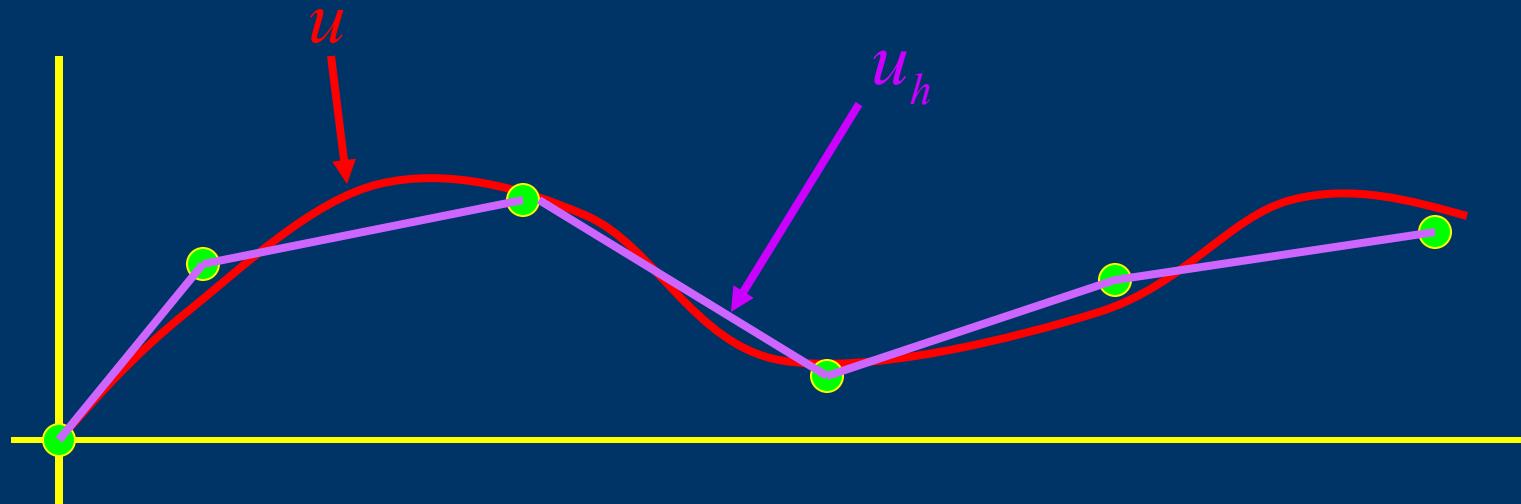
U is restricted to be 0 at 0 and 1!!

Using the norm properties, it is possible to show

If $a(u_h, \varphi_i) = l(\varphi_i)$ for all $\varphi_i \in \{\varphi_1, \varphi_2, \dots, \varphi_n\}$

Then $\underbrace{\|u - u_h\|}_{\text{Solution Error}} = \min_{w_h \in X_h} \underbrace{\|u - w_h\|}_{\text{Projection Error}}$

The question is only



How well can you fit u with a member of X_h

But you must measure the error in the $\| \cdot \|$ norm

For piecewise linear:

$$\underbrace{\|u - u_h\|}_{\text{error}} = O\left(\frac{1}{n}\right)$$

Problem of interest

Helmholtz Equation in 1D

Boundary Value Problem (BVP) - Strong Form

$$-u''(x) + \alpha u(x) = f(x) \quad \alpha \geq 0$$

$$x \in (0, 1), \quad u(0) = u(1) = 0$$

Describes many physical phenomena (e.g.) :

- Temperature distribution in a bar *
- Deformation of an elastic bar
- Deformation of a string under tension

N1

N2

Problem of interest

Solution Properties

- the solution $u(x)$ always *exists*
- $u(x)$ is always smoother than the data $f(x)$
- given $f(x)$ the solution $u(x)$ is *unique*

Minimization Principle

Statement

Find

$$u = \arg \min_{w \in X} J(w)$$

where

$$X = \{v \text{ sufficiently smooth} \mid v(0) = v(1) = 0\},$$

and

$$J(w) = \frac{1}{2} \int_0^1 (w_x w_x + \alpha w w) dx - \int_0^1 f w dx$$

Minimization Principle

Statement

In words:

Over all functions \mathbf{w} in \mathbf{X} ,

\mathbf{u} that satisfies

$$-\mathbf{u}_{xx} + \alpha \mathbf{u} = f \quad \text{in } \Omega$$

$$\mathbf{u}(0) = \mathbf{u}(1) = 0$$

makes $J(\mathbf{w})$ as small as possible.

N4

Minimization Principle

Statement

Proof...

Let $\mathbf{w} = \mathbf{u} + \mathbf{v}$.

Then

$$\begin{aligned} J(\underbrace{\mathbf{u}}_{\in X} + \underbrace{\mathbf{v}}_{\in X}) &= \frac{1}{2} \int_0^1 (\mathbf{u} + \mathbf{v})_x (\mathbf{u} + \mathbf{v})_x \, dx \\ &\quad + \frac{\alpha}{2} \int_0^1 (\mathbf{u} + \mathbf{v})(\mathbf{u} + \mathbf{v}) \, dx \\ &\quad - \int_0^1 f(\mathbf{u} + \mathbf{v}) \, dx . \end{aligned}$$

Minimization Principle

Statement

...Proof...

$$J(u+v) = \frac{1}{2} \int_0^1 (u_x u_x + \alpha u u) dx - \int_0^1 f u dx \quad J(u)$$

$$+ \int_0^1 (u_x v_x + \alpha u v) dx - \int_0^1 f v dx \quad \delta J_v(u)$$

first variation

$$+ \frac{1}{2} \int_0^1 (v_x v_x + \alpha v v) dx > 0 \text{ for } v \neq 0$$

Minimization Principle

Statement

...Proof...

$$\begin{aligned}\delta J_v(u) &= \int_0^1 (u_x v_x + \alpha u v) dx - \int_0^1 f v dx \\&= \vec{x}^0(0) u_x(0) - \vec{x}^0(1) u_x(1) - \int_0^1 u_{xx} v dx \\&\quad + \alpha \int_0^1 u v dx - \int_0^1 f v dx \\&= \int_0^1 v \underbrace{\{-u_{xx} + \alpha u - f\}}_n dx = 0, \quad \forall v \in X\end{aligned}$$

Minimization Principle

Statement

...Proof

$$J(\underbrace{u+v}_w) = J(u) + \frac{1}{2} \underbrace{\int_0^1 (v_x v_x + \alpha v v) dx}_{> 0 \text{ unless } v = 0}, \forall v \in X$$

$$\Rightarrow \boxed{J(w) > J(u), \quad \forall w \in X, w \neq u}$$

\Updownarrow

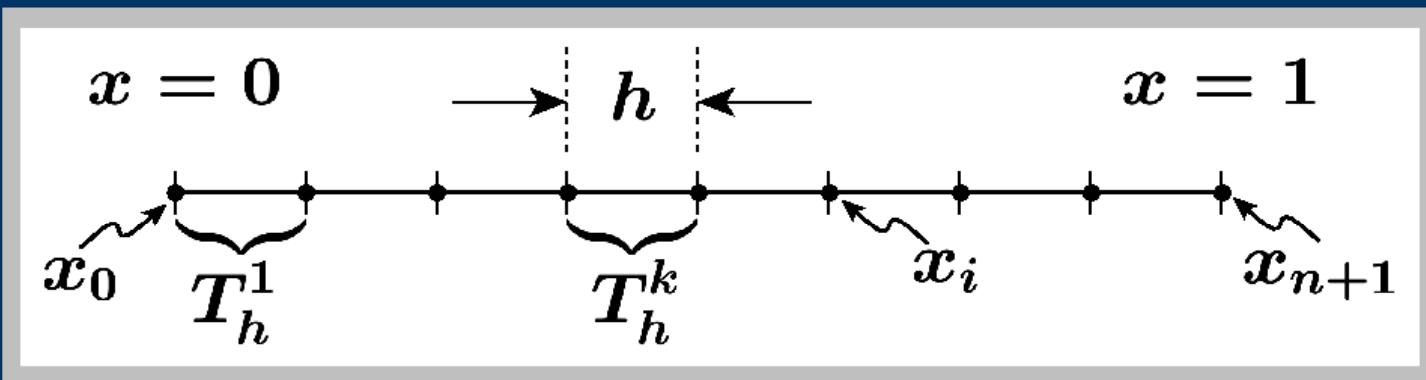
w is the minimizer of $J(w)$

E1 N5

Rayleigh-Ritz Approach

Approximation

Mesh


$$\overline{\Omega} = \bigcup_{k=1}^K \overline{T}_h^k \quad T_h^k, k = 1, \dots, K = n + 1: \text{elements}$$
$$x_i, i = 0, \dots, n + 1: \text{nodes}$$

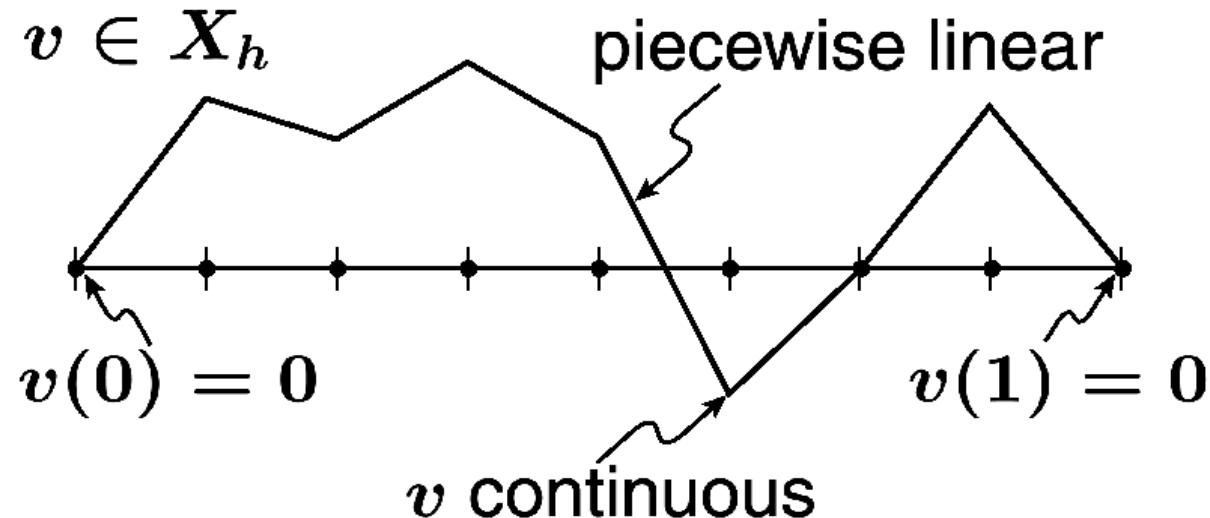
N6

Rayleigh-Ritz Approach

Approximation

Space $X_h \subset X$

$$X_h = \left\{ v \in X \mid v|_{T_h^k} \in \mathbb{P}_1(T_h^k), \quad k = 1, \dots, K \right\}$$



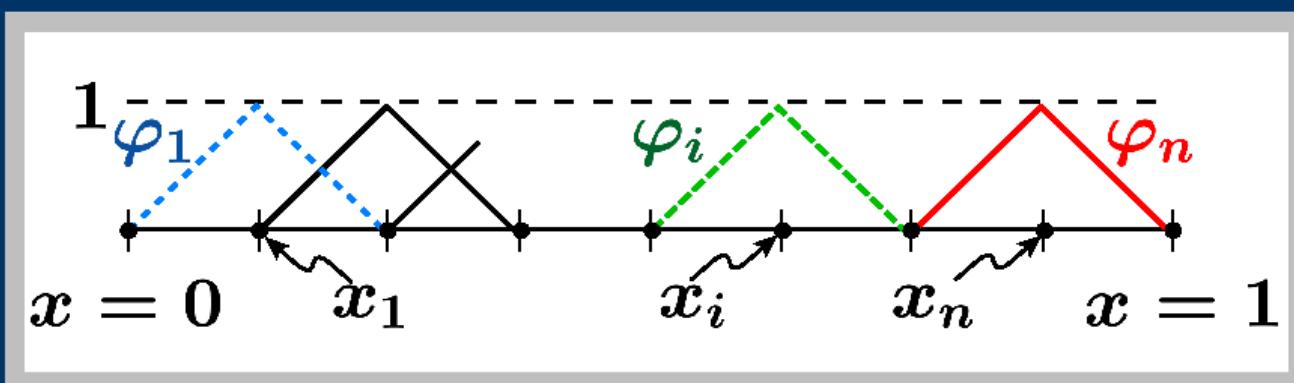
Rayleigh-Ritz Approach

Approximation

Basis

Nodal basis for X_h :

$$\varphi_j, j = 1, \dots, n = \dim(X_h)$$



φ_i nonzero only on $\overline{T}_h^i \cup \overline{T}_h^{i+1}$

N7

N8

Rayleigh-Ritz Approach

“Projection”

Plan...

Let

$$\underbrace{\mathbf{u}_h \ (\in X_h)}_{\text{RR/FE Approximation}} = \sum_{j=1}^n \mathbf{u}_{hj} \varphi_j(\mathbf{x}) ;$$

set $\mathbf{u}_{hj} = \mathbf{w}_j$ that minimize

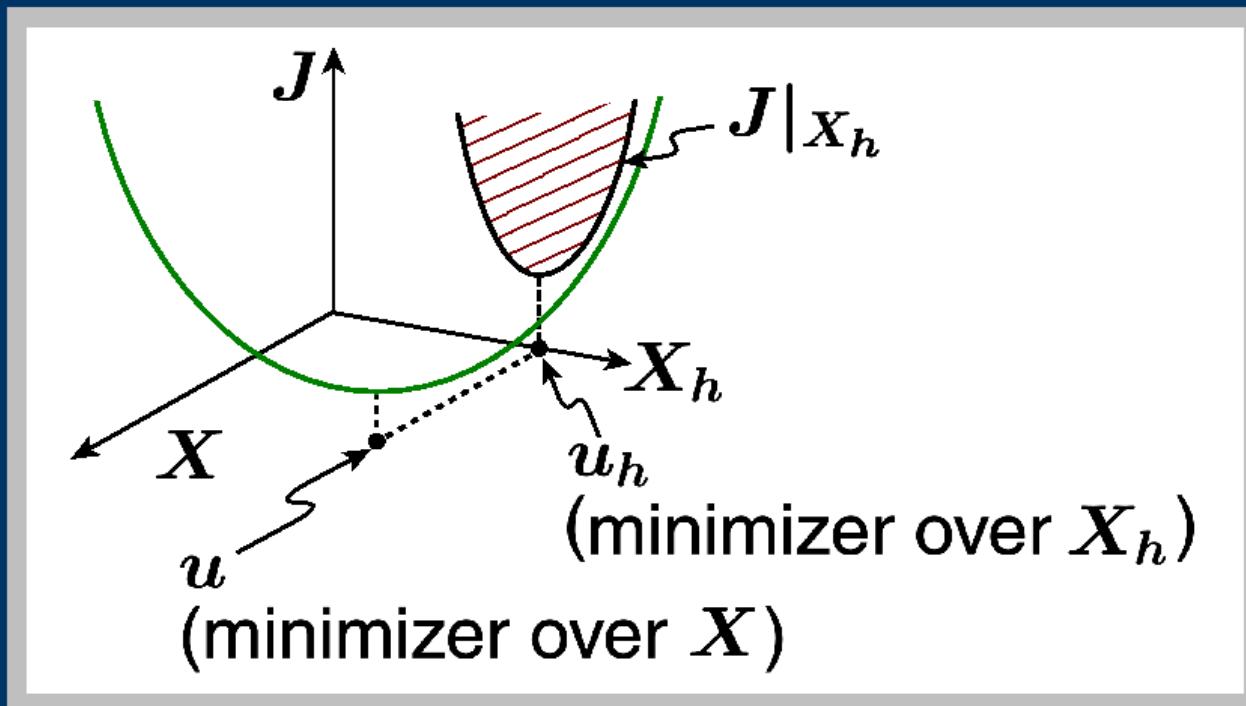
$$J \left(\sum_{j=1}^n \mathbf{w}_j \varphi_j \right) .$$

Rayleigh-Ritz Approach

“Projection”

...Plan

Geometric Picture:



Rayleigh-Ritz Approach

“Projection”

$J|_{X_h}$...

$$\begin{aligned} J \left(\sum_{j=1}^n w_j \varphi_j \right) &= \frac{1}{2} \int_0^1 \frac{d}{dx} \left(\sum_{i=1}^n w_i \varphi_i \right) \frac{d}{dx} \left(\sum_{j=1}^n w_j \varphi_j \right) \\ &\quad + \frac{\alpha}{2} \int_0^1 \sum_{i=1}^n (w_i \varphi_i) \sum_{j=1}^n (w_j \varphi_j) - \int_0^1 f \sum_{j=1}^n w_j \varphi_j \\ &= \frac{1}{2} \sum_{i=1}^n \sum_{j=1}^n w_i w_j \int_0^1 \left(\frac{d\varphi_i}{dx} \frac{d\varphi_j}{dx} + \alpha \varphi_i \varphi_j \right) dx - \sum_{j=1}^n w_j \int_0^1 f \varphi_j dx \end{aligned}$$

by *bilinearity* and *linearity*.

Rayleigh-Ritz Approach

“Projection”

$\dots J|_{X_h}$

$$\begin{aligned} \underline{J}^R(\underline{w} \in \mathbb{R}^n) &\equiv J \left(\sum_{j=1}^n w_j \varphi_j \right) \\ &= \frac{1}{2} \underline{w}^T \underline{A}_h \underline{w} - \underline{w}^T \underline{F}_h . \end{aligned}$$

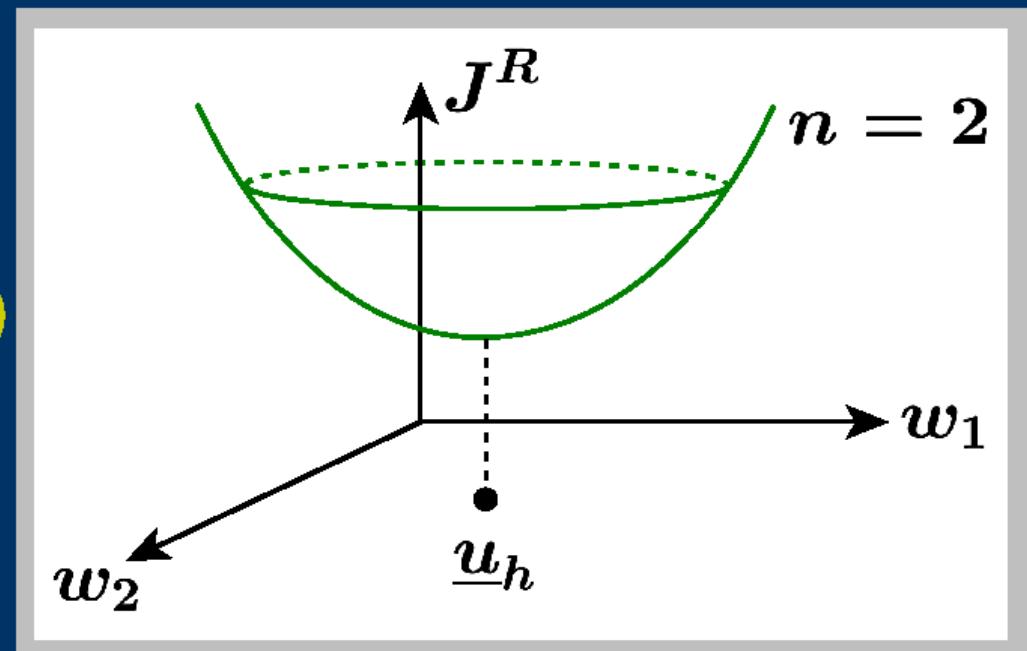
$$\begin{aligned} \underline{F}_h \in \mathbb{R}^n: F_{h,i} &= \int_0^1 f \varphi_i dx \\ \underline{A}_h \in \mathbb{R}^{n \times n}: A_{h,ij} &= \int_0^1 \left(\frac{d\varphi_i}{dx} \frac{d\varphi_j}{dx} + \alpha \varphi_i \varphi_j \right) dx \end{aligned}$$

Rayleigh-Ritz Approach

“Projection”

Minimization...

$$\underline{u}_h = \arg \min_{\underline{w} \in \mathbb{R}^n} J^R(\underline{w})$$



Expand $J(\underline{w} = \underline{u}_h + \underline{v})$; require $J(\underline{w}) > J(\underline{u}_h)$ unless $\underline{v} = 0$.

Rayleigh-Ritz Approach

“Projection”

...Minimization...

$$J^R(\underline{u}_h + \underline{v})$$

$$\begin{aligned} &= \frac{1}{2} (\underline{u}_h + \underline{v})^T \underline{A}_h (\underline{u}_h + \underline{v}) - (\underline{u}_h + \underline{v})^T \underline{F}_h \\ &= \frac{1}{2} \underline{u}_h^T \underline{A}_h \underline{u}_h - \underline{u}_h^T \underline{F}_h \\ &\quad + \frac{1}{2} \underline{v}^T \underline{A}_h \underline{u}_h + \frac{1}{2} \underline{u}_h^T \underline{A}_h \underline{v} - \underline{v}^T \underline{F}_h \\ &\quad + \frac{1}{2} \underline{v}^T \underline{A}_h \underline{v} \end{aligned}$$

Rayleigh-Ritz Approach

“Projection”

...Minimization...

$$\mathbf{J}^R(\underline{\mathbf{u}}_h + \underline{\mathbf{v}}) = \mathbf{J}(\underline{\mathbf{u}})$$

$$+ \underbrace{(\underline{\mathbf{A}}_h \underline{\mathbf{u}}_h - \underline{\mathbf{F}}_h)^T \underline{\mathbf{v}}}_{\nabla J^R(\underline{\mathbf{u}}_h)} \quad \delta J_{\underline{\mathbf{v}}}^R(\underline{\mathbf{u}}_h) \quad SPD$$

$$+ \frac{1}{2} \underbrace{\underline{\mathbf{v}}^T \underline{\mathbf{A}}_h \underline{\mathbf{v}}}_{>0, \forall \underline{\mathbf{v}} \neq 0} \quad SPD$$

Rayleigh-Ritz Approach

“Projection”

...Minimization

If (and only if)

$$\delta J_{\underline{v}}^R(\underline{u}_h) = \mathbf{0}, \quad \forall \underline{v} \in \mathbb{R}^n$$

\Updownarrow

$$\nabla J^R(\underline{u}_h) = \underline{A}_h \underline{u}_h - \underline{F}_h = \mathbf{0}$$

then

$$J(\underline{w} = \underline{u}_h + \underline{v}) > J(\underline{u}_h), \quad \forall \underline{v} \neq \mathbf{0}.$$

Rayleigh-Ritz Approach

“Projection”

Final Result

Find $\underline{u}_h \in \mathbb{R}^n$ such that

$$\underline{A}_h \underline{u}_h = \underline{F}_h \quad \Rightarrow \quad u_h(x) = \sum_{j=1}^n u_{hj} \varphi_j(x) .$$

SPD \Rightarrow existence and uniqueness.

Energy norm

Error Analysis

Remember

$$J(u+v) = J(u) + \frac{1}{2} \underbrace{\int_0^1 (v_x v_x + \alpha v v) dx}_{\geq 0, SPD}, \quad \forall v \in X$$

Define

$$|||v||| = \left[\int_0^1 (v_x v_x + \alpha v v) dx \right]^{\frac{1}{2}}$$

Energy norm

Energy norm

Error Analysis

Therefore

$$J(\mathbf{u} + \mathbf{v}) = J(\mathbf{u}) + \frac{1}{2}|||\mathbf{v}|||^2, \quad \forall \mathbf{v} \in \mathbf{X}$$

Choose any $\mathbf{w}_h \in \mathbf{X}_h$, $\mathbf{v} \rightarrow (\mathbf{w}_h - \mathbf{u}) \in \mathbf{X}$

$$J(\mathbf{w}_h) = J(\mathbf{u}) + \frac{1}{2}|||\mathbf{u} - \mathbf{w}_h|||^2, \quad \forall \mathbf{w}_h \in \mathbf{X}_h$$

For $\mathbf{w}_h = \mathbf{u}_h$

$$J(\mathbf{u}_h) = J(\mathbf{u}) + \frac{1}{2}|||\mathbf{u} - \mathbf{u}_h|||^2$$

Energy norm

Error Analysis

$$J(u_h) < J(w_h), \forall w_h \in X_h, w_h \neq u_h$$

if $e = u - u_h$

$$\|\underline{u - u_h}\|_e < \|u - w_h\|, \forall w_h \in X_h, w_h \neq u_h$$

and

$$\|e\|_e = \inf_{w_h \in X_h} \|u - w_h\|$$

Energy norm

Error Analysis

In words: even if you *knew* \mathbf{u} ,

you could not find a \mathbf{w}_h in \mathbf{X}_h

more accurate than \mathbf{u}_h

in the energy norm.

A priori theory

Error Analysis

A priori error estimates

N9

Energy norm:

$$|||e||| \leq C_1 h$$

L_2 norm:

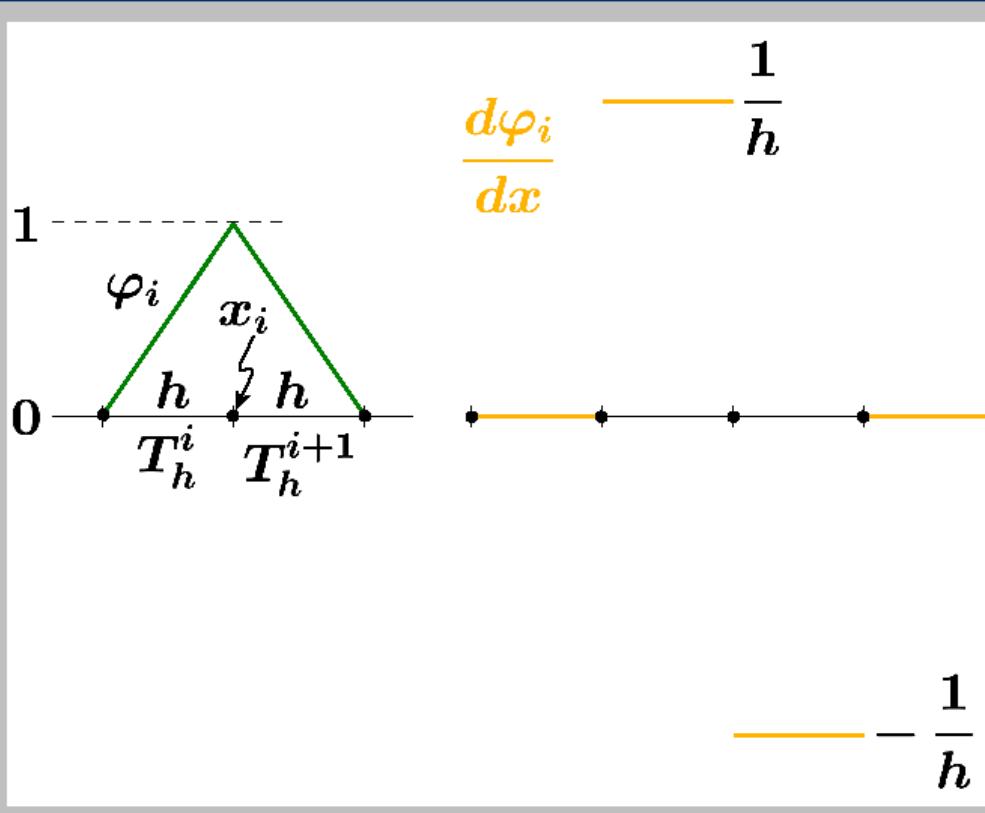
$$\|e\| = \left(\int_0^1 e \cdot e \, dx \right)^{1/2} \leq C_2 h^2$$

$$C_{1,2} = \mathcal{F}(\Omega, \text{problem parameters, smoothness of } u)$$

Discrete Equations

Matrix Elements: \underline{A}_h^1

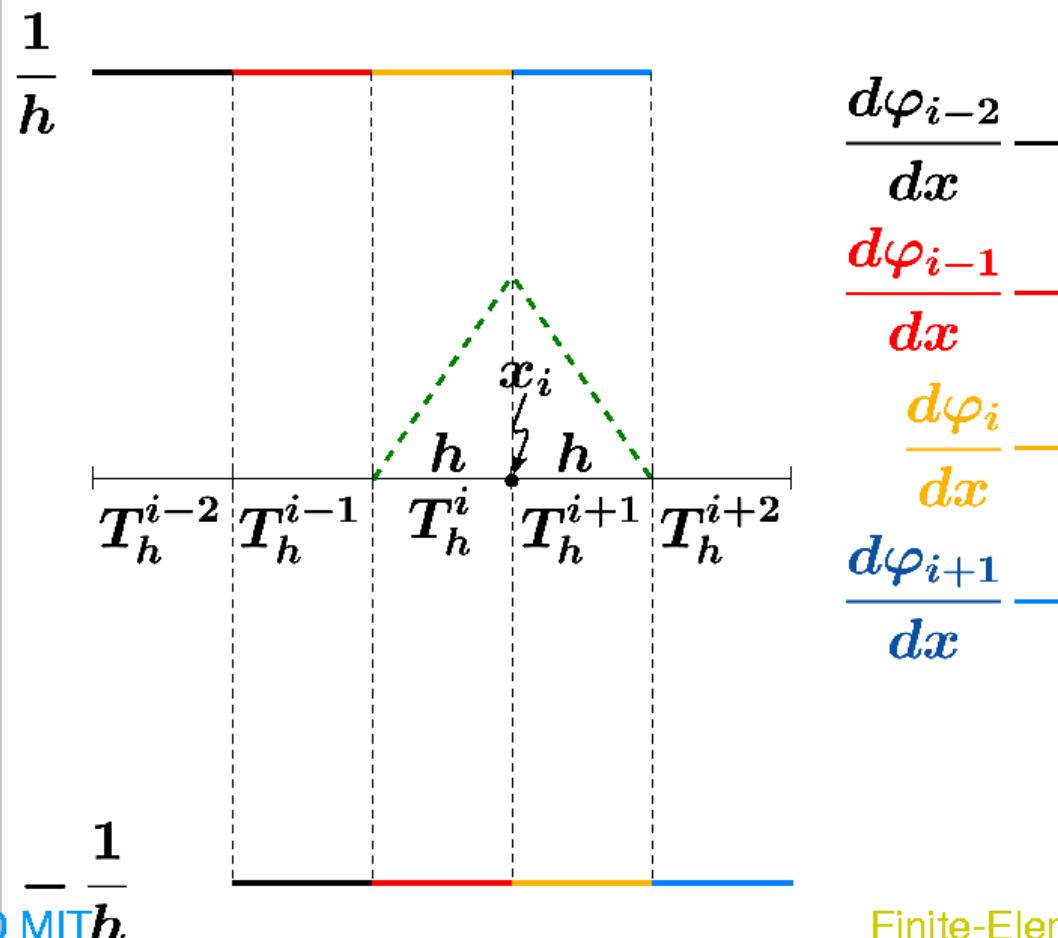
φ_i and $d\varphi_i/dx\dots$



Discrete Equations

Matrix Elements: \underline{A}_h^1

$\dots \varphi_i$ and $d\varphi_i/dx$



Discrete Equations

Matrix Elements: \underline{A}_h^1

Typical Row

$$A_{h i j}^1 = \int_{\Omega} \frac{d\varphi_i}{dx} \frac{d\varphi_j}{dx} dx = \int_{T_h^i} \frac{d\varphi_i}{dx} \frac{d\varphi_j}{dx} dx + \int_{T_h^{i+1}} \frac{d\varphi_i}{dx} \frac{d\varphi_j}{dx} dx$$

is nonzero only for $i = j - 1, j, j + 1$

$$A_{h i i}^1 = \frac{1}{h^2} (h) + \frac{1}{h^2} (h) = \frac{2}{h}$$

$$A_{h i i-1}^1 = \frac{1}{h} \left(-\frac{1}{h} \right) (h) = -\frac{1}{h}$$

$$A_{h i i+1}^1 = \left(-\frac{1}{h} \right) \frac{1}{h} (h) = -\frac{1}{h}$$

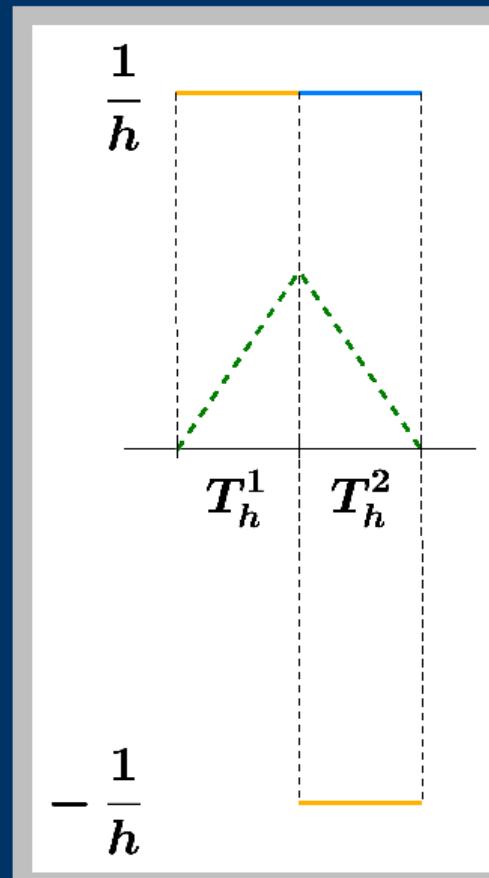
Discrete Equations

Matrix Elements: A_h^1

Boundary Rows

$$A_{h11}^1 = \frac{2}{h}, \quad A_{h12}^1 = -\frac{1}{h},$$

$$A_{hn_n}^1 = \frac{2}{h}, \quad A_{hn_{n-1}}^1 = -\frac{1}{h}.$$



Discrete Equations

Matrix Elements: \underline{A}_h^1

Properties of \underline{A}_h

$$\underline{A}_h^1 = \frac{1}{h} \begin{pmatrix} 2 & -1 & & & & 0 \\ -1 & 2 & -1 & & & \dots \\ & & & \ddots & & \\ 0 & & & & 2 & -1 \\ & & & & -1 & 2 \end{pmatrix}$$

\underline{A}_h^1 is SPD; *and*
diagonally dominant; *and*
sparse; *and*
tridiagonal.

Discrete Equations

Mass Matrix

$\underline{\boldsymbol{M}_h} \in \mathbf{R}^{n \times n}$:

$$\boldsymbol{M}_{h\ ij} = \int_0^1 \varphi_i \varphi_j \, dx$$

the finite element “identity” (\boldsymbol{I}) operator

Is nonzero only for $i = j - 1, j, j + 1$

$$\boldsymbol{M}_{h\ ij} = \int_{T_h^i} \varphi_i \varphi_j \, dx + \int_{T_h^{i+1}} \varphi_i \varphi_j \, dx$$

Discrete Equations

Mass Matrix

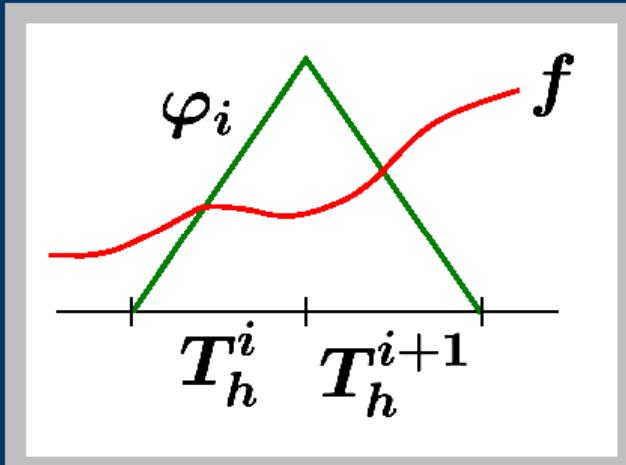
For linear elements, nodal basis:

$$\underline{\mathbf{M}}_h = \mathbf{h} \begin{pmatrix} \frac{2}{3} & \frac{1}{6} & & & & 0 \\ \frac{1}{6} & \frac{2}{3} & \frac{1}{6} & & & \dots \\ & & & \ddots & & \\ 0 & & & & \frac{2}{3} & \frac{1}{6} \\ & & & & \frac{1}{6} & \frac{2}{3} \end{pmatrix}$$

sparse, banded, tri-diagonal — “close” to $\underline{\mathbf{I}}$.

Discrete Equations

“Load” Vector Elements: \underline{F}_h



$$\mathbf{F}_{h,i} = \int_0^1 f \varphi_i \, dx$$

$$\mathbf{F}_{h,i} = \int_{T_h^i} \mathbf{f} \cdot \boldsymbol{\varphi}_i \, dx + \int_{T_h^{i+1}} \mathbf{f} \cdot \boldsymbol{\varphi}_i \, dx, \quad i = 1, \dots, n;$$

Discrete Equations

Summary

$\underline{u}_h \in \mathbb{R}^n$ satisfies

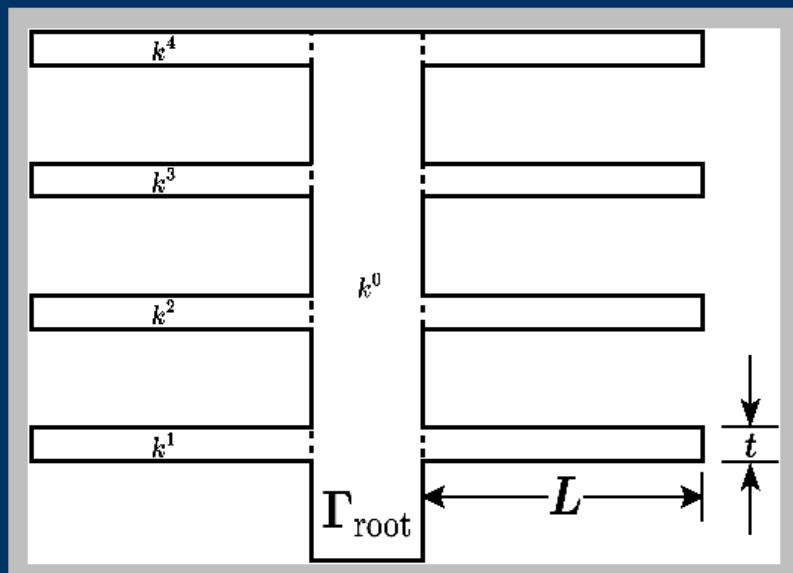
$$[\underline{A}_h^1 + \alpha \underline{M}_h] \begin{pmatrix} \underline{u}_{h,1} \\ \vdots \\ \underline{u}_{h,n} \end{pmatrix} = \begin{pmatrix} \underline{F}_{h,1} \\ \vdots \\ \underline{F}_{h,n} \end{pmatrix}$$

N10

Heat Transfer Problem

Example

Non-dimensional form



\mathbf{k}^i : Thermal conductivity
for Ω_i , $i = 0, \dots, 4$

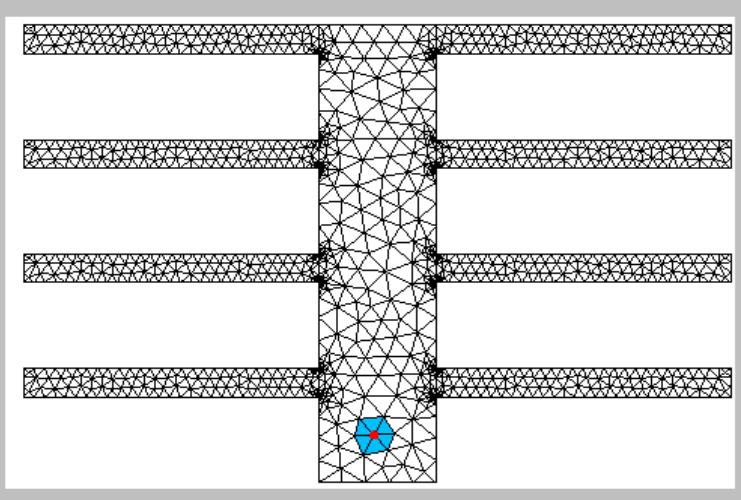
\mathbf{Bi} : Heat transfer coefficient

$\frac{\mathbf{t}}{\mathbf{L}}$: Geometric parameters

Finite element method

Example

$$X_h = \text{span}\{\varphi_1, \dots, \varphi_n\}$$



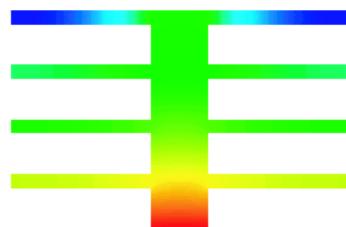
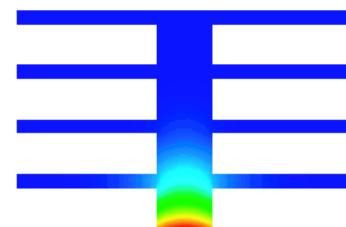
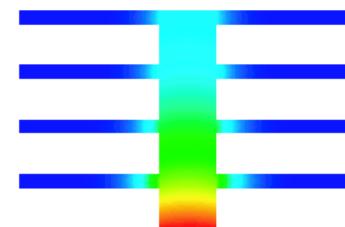
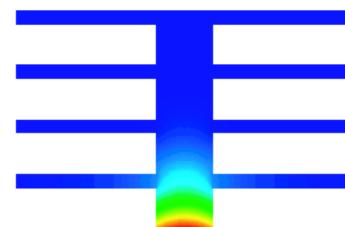
$\varphi_i(x)$:
Nodal basis functions

– First order elements

$$\dim(X_h) = n$$

Possible solutions

Example

 μ_1  μ_2  μ_3  μ_4 

Extensions

Example

- Complicated geometries
- General classes of problems
(Good mathematical properties)
- Wider class of operators

Summary

- Why Poisson Equation
 - Reminder about heat conducting bar
- Finite-Difference And Basis function methods
 - Key question of convergence
- Convergence of Finite-Element methods
 - Key idea: solve Poisson by minimization
 - Demonstrate optimality in a carefully chosen norm