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6.013, Electromagnetic Fields, Forces, and Motion Prof. Markus Zahn, Sept. 13, 2005

Lecture 2: Electromagnetic Field Boundary Conditions

I. Boundary Conditions

1. Gauss' Continuity Condition

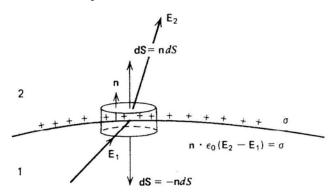


Figure 2-19 Gauss's law applied to a differential sized pill-box surface enclosing some surface charge shows that the normal component of $\varepsilon_0 \mathbf{E}$ is discontinuous in the surface charge density.

Zahn, Markus. Figs. 1.13-1.1.17, 1.19 (a) and (b), 1.23, 1.20, 2.19, 3.12 (a). Electromagnetic Field Theory: A Problem Solving Approach. Robert E. Krieger Publishing Company, Florida, 1987. Used with permission.

$$\begin{split} \oint_{S} \epsilon_{_{0}} \overline{E} \bullet \overline{da} &= \int_{S} \sigma_{_{S}} dS \Rightarrow \epsilon_{_{0}} \left(E_{_{2n}} - E_{_{1n}} \right) dS = \sigma_{_{S}} dS \\ \epsilon_{_{0}} \left(E_{_{2n}} - E_{_{1n}} \right) &= \sigma_{_{S}} \Rightarrow \overline{n} \bullet \left[\epsilon_{_{0}} \left(\overline{E}_{_{2}} - \overline{E}_{_{1}} \right) \right] = \sigma_{_{S}} \end{split}$$

2. Continuity of Tangential E

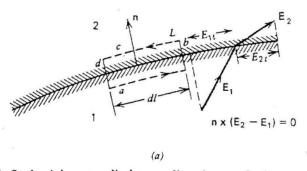


Figure 3-12 (a) Stokes' law applied to a line integral about an interface of discontinuity shows that the tangential component of electric field is continuous across the boundary.

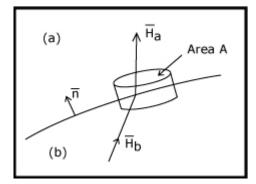
Zahn, Markus. Figs. 1.13-1.1.17, 1.19 (a) and (b), 1.23, 1.20, 2.19, 3.12 (a). Electromagnetic Field Theory: A Problem Solving Approach. Robert E. Krieger Publishing Company, Florida, 1987. Used with permission.

$$\begin{split} \oint_{c} \overline{E} \bullet \overline{ds} &= \left(E_{1t} - E_{2t}\right) dI = 0 \Rightarrow E_{1t} - E_{2t} = 0 \\ &= \overline{n} \times \left(\overline{E}_{1} - \overline{E}_{2}\right) = 0 \end{split}$$

Equivalent to $\Phi_1 = \Phi_2$ along boundary

3. Normal H

$$\oint_S \mu_0 \ \overline{H} \bullet \overline{da} = 0$$



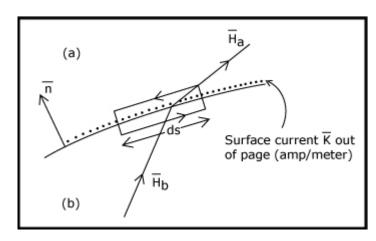
$$\mu_0 \left(H_{an} - H_{bn} \right) A = 0$$

$$H_{an} = H_{bn}$$

$$\overline{n} \cdot [\overline{H}_a - \overline{H}_b] = 0$$

4. Tangential H

$$\oint\limits_{C}\overline{H}\bullet\overline{ds}=\int\limits_{S}\overline{J}\bullet\overline{da}+\frac{d}{dt}\int\limits_{S}\epsilon_{0}\overline{E}\bullet\overline{da}$$

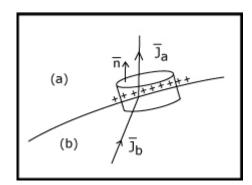


$$H_{bt}ds - H_{at}ds = Kds$$

$$H_{bt} - H_{at} = K$$

$$\overline{n} \times \left[\overline{H}_a - \overline{H}_b\right] = \overline{K}$$

5. Conservation of Charge Boundary Condition

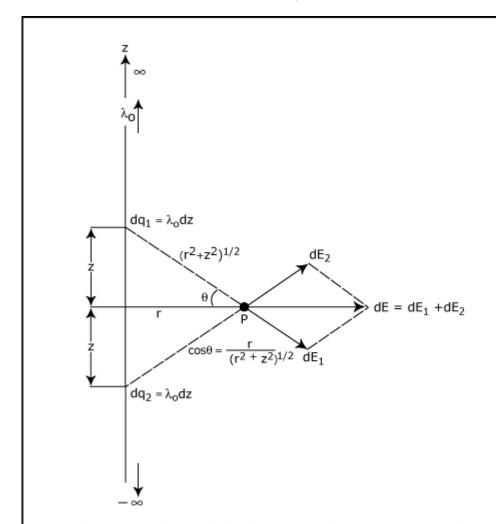


$$\oint_{S} \overline{J} \cdot \overline{da} + \frac{d}{dt} \int_{V} \rho dV = 0$$

$$\bar{n} \cdot [\bar{J}_a - \bar{J}_b] + \frac{\partial}{\partial t} \sigma_s = 0$$

II. Boundary Condition Problems

- 1. Electric Field from a Sheet of Surface Charge
 - a. Electric Field from a Line Charge



An infinitely long uniform distributaion of line charge only has a radially directed electric field because the z components of the electric field are canceled out by symmetrically located incremental charge elements as also shown in Figure 2-8a.

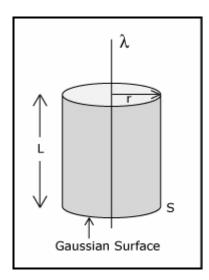
$$dE_{r}=\frac{dq}{4\pi\epsilon_{0}\left(r^{2}+z^{2}\right)}cos\theta=\frac{\lambda_{0}rdz}{4\pi\epsilon_{0}\left(r^{2}+z^{2}\right)^{3/2}}$$

$$E_r = \int\limits_{z=-\infty}^{+\infty} dE_r = \frac{\lambda_0 r}{4\pi\epsilon_0} \int\limits_{z=-\infty}^{+\infty} \frac{dz}{\left(r^2+z^2\right)^{\!\!\frac{3}{2}}}$$

$$= \frac{\lambda_0 r}{4\pi\epsilon_0} \frac{z}{r^2 \left(z^2 + r^2\right)^{1/2}} \Bigg|_{z=-\infty}^{+\infty}$$

$$=\frac{\lambda_0}{2\pi\epsilon_0 r}$$

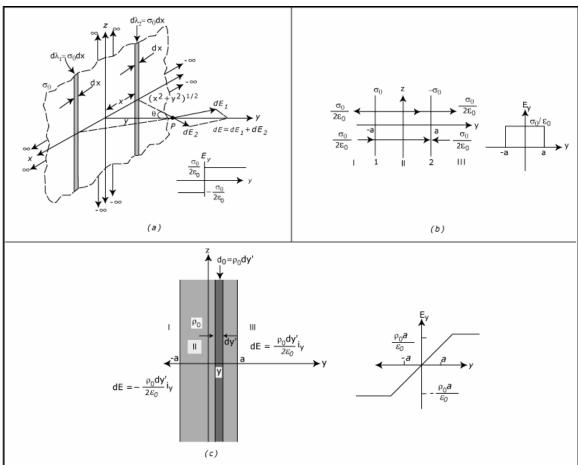
Another way: Gauss' Law



$$\int_{S} \epsilon_{0} \overline{E} \cdot \overline{da} = \epsilon_{0} E_{r} 2\pi r L = \lambda_{0} L$$

$$E_r \, = \frac{\lambda_0}{2\pi\,\epsilon_0\,r}$$

b. Electric Field from a Sheet Charge



(a) The electric field from a uniformly surface charged sheet of infinite extent is found by summing the contributions from each incremental line charge element. Symmetrically placed line charge elements have x field components that cancel, but y field components that add. (b) Two parallel but oppositely sheets of surface charge have fields that add in the region between the sheets but cancel outside. (c) The electric field from a volume charge distribution is obtained by summing the contributions from each incremental surface charge element.

$$\begin{split} dE_y &= \frac{d\lambda}{2\pi\epsilon_0 \left(x^2 + y^2\right)^{1/2}} cos \, \theta = \frac{\sigma_0 y dx}{2\pi\epsilon_0 \left(x^2 + y^2\right)} \\ E_y &= \int\limits_{x = -\infty}^{+\infty} dE_y = \frac{\sigma_0 y}{2\pi\epsilon_0} \int\limits_{x = -\infty}^{+\infty} \frac{dx}{x^2 + y^2} \\ &= \frac{\sigma_0 y}{2\pi\epsilon_0} \frac{1}{y} tan^{-1} \left. \frac{x}{y} \right|_{-\infty}^{+\infty} \\ &= \begin{cases} \frac{\sigma_0}{2\epsilon_0} & y > 0 \\ -\frac{\sigma_0}{2\epsilon_0} & y < 0 \end{cases} \end{split}$$

Checking Boundary condition at y=0

$$\mathsf{E}_{\mathsf{y}}\left(\mathsf{y}=\mathsf{O}_{\scriptscriptstyle{+}}\right)-\mathsf{E}_{\mathsf{y}}\left(\mathsf{y}=\mathsf{O}_{\scriptscriptstyle{-}}\right)=\frac{\sigma_{\mathsf{O}}}{\epsilon_{\mathsf{O}}}$$

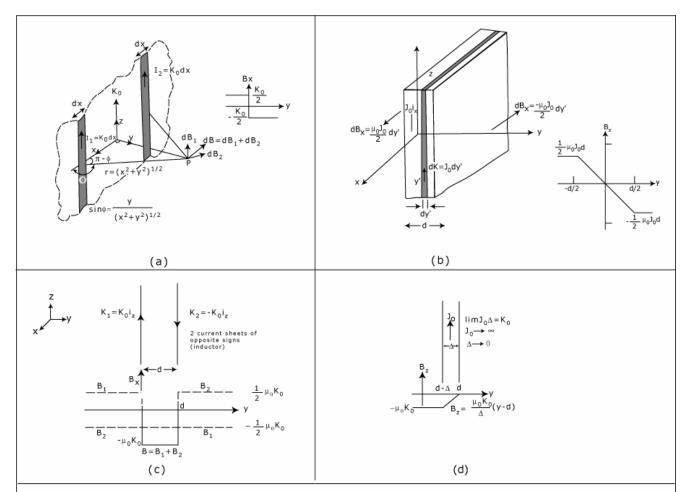
$$\frac{\sigma_0}{2\epsilon_0} - \left(-\frac{\sigma_0}{2\epsilon_0}\right) = \frac{\sigma_0}{\epsilon_0}$$

c. Two sheets of Surface Charge (Capacitor)

$$\bar{E}_1 = \begin{cases} \frac{\sigma_0}{2\epsilon_0} \bar{i}_y & y > -a \\ -\frac{\sigma_0}{2\epsilon_0} \bar{i}_y & y < -a \end{cases}, \ \bar{E}_2 = \begin{cases} -\frac{\sigma_0}{2\epsilon_0} \bar{i}_y & y > a \\ \frac{\sigma_0}{2\epsilon_0} \bar{i}_y & y < a \end{cases}$$

$$\vec{E} = \vec{E}_1 + \vec{E}_2 = \begin{cases} \frac{\sigma_0}{\epsilon_0} \vec{i}_y & |y| < a \\ 0 & |y| > a \end{cases}$$

2. Magnetic Field from a Sheet of Surface Current



A uniform surface current of infinite extent generates a uniform magnetic field oppositely directed on each side of the sheet. The magnetic field is perpendicular to the surface current but parellel to the plane of the sheet. (b) The magnetic field due to a slab of the volume current is found by superimposing the fields due to incremental surface currents. (c) Two parallel but oppositely directed surface current sheets have fields that add in the region between the sheets but cancel outside the sheet. (d) The force on a current sheet is due to the average field on each side of the sheet as found by modeling the sheet as a uniform volume current distributed over an infinitesimal thickness \$\Delta\$.

From a line current I

$$H_{\phi} = \frac{I}{2\pi r}$$

$$\bar{i}_{\phi} = -\sin\phi \, \bar{i}_{x} + \cos\phi \, \bar{i}_{y}$$

Thus from 2 symmetrically located line currents

$$dH_{x} = \frac{dI}{2\pi \left(x^{2} + y^{2}\right)^{1/2}} \left(-\sin\phi\right)$$

$$=-\frac{K_0dx}{2\pi}\frac{y}{x^2+y^2}$$

$$H_{x} = -\frac{K_{0}}{2\pi} y \int_{x=-\infty}^{+\infty} \frac{dx}{x^{2} + y^{2}}$$

$$= -\frac{K_0 y}{2\pi} \frac{1}{y} \tan^{-1} \frac{x}{y} \bigg|_{x=-\infty}^{+\infty}$$

$$= \begin{cases} -\frac{K_0}{2} & y > 0 \\ +\frac{K_0}{2} & y < 0 \end{cases}$$

Check boundary condition at y=0:

$$H_x(y = O_+) - H_x(y = O_-) = -K_0$$

$$-\frac{\mathsf{K}_0}{2} - \left(\frac{\mathsf{K}_0}{2}\right) = -\mathsf{K}_0$$