

Recitation 18: BJT-Regions of Operation & Small Signal Model

BJT: Regions of Operation

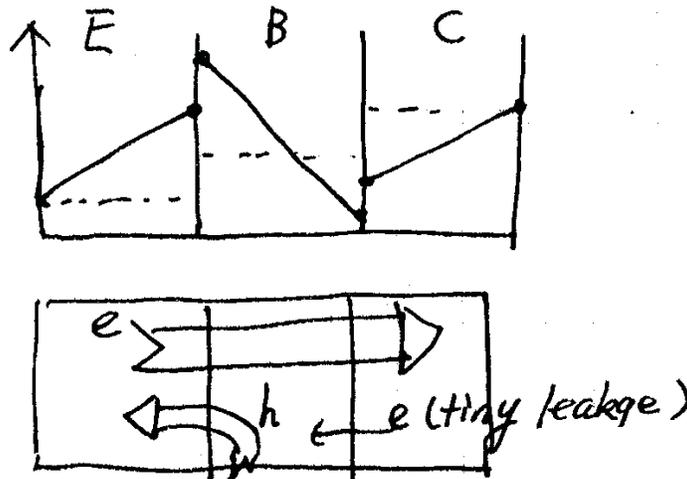
System of equations that describes BJT operation:

$$\begin{aligned}
 I_c &= I_s \left(\exp\left(\frac{qV_{BE}}{kT}\right) - \exp\left(\frac{qV_{BC}}{kT}\right) \right) - \frac{I_s}{\beta_R} \left(\exp\left(\frac{qV_{BC}}{kT}\right) - 1 \right) \\
 I_B &= \frac{I_s}{\beta_F} \left(\exp\left(\frac{qV_{BE}}{kT}\right) - 1 \right) + \frac{I_s}{\beta_R} \left(\exp\left(\frac{qV_{BC}}{kT}\right) - 1 \right) \\
 I_E &= -I_s \left(\exp\left(\frac{qV_{BE}}{kT}\right) - \exp\left(\frac{qV_{BC}}{kT}\right) \right) - \frac{I_s}{\beta_F} \left(\exp\left(\frac{qV_{BE}}{kT}\right) - 1 \right)
 \end{aligned}$$

$$\begin{aligned}
 I_s &= \frac{qA_E n_i^2 D_{nB}}{N_{aB} w_B} \\
 \beta_F &= \frac{N_{dF} D_{nB} w_E}{N_{aB} D_{pE} w_B} \\
 \beta_R &= \frac{N_{dC} D_{nB} w_C}{N_{aB} D_{pC} w_B}
 \end{aligned}$$

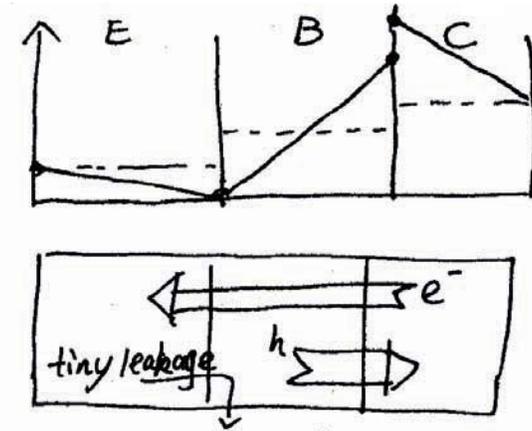
This set of equations can describe all four regimes of operation for BJT

Forward Active: $V_{BE} > 0, V_{BC} < 0$



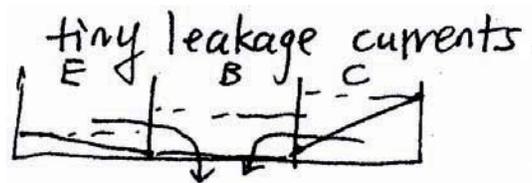
Reverse Active (RAR)

$V_{BE} < 0, V_{BC} > 0$



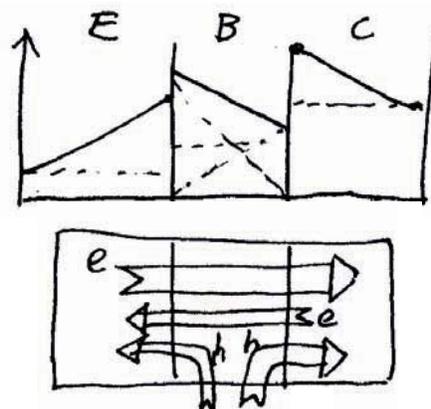
Cut-off

$V_{BE} < 0, V_{BC} < 0$

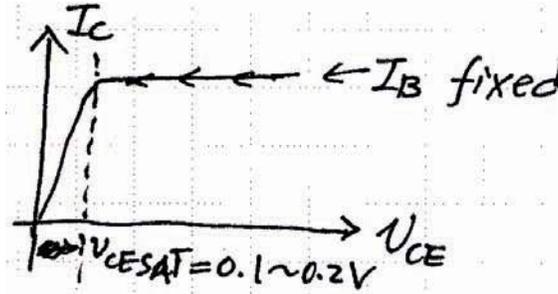


Saturation

$V_{BE} > 0, V_{BC} > 0$



Understanding the I_C vs. V_{CE} curve: I_C drops rapidly below $V_{CE,SAT} \simeq 0.1$ to 0.2 V.

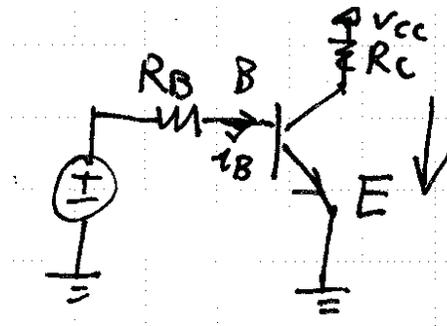


Why?

- Each curve I_B is fixed
- $V_{CE} = V_{BE} - V_{BC}$, $\implies V_{BC} = V_{BE} - V_{CE}$
- When V_{CE} is large, $V_{BC} < 0$, FAR. As we reduce V_{CE} , V_{BC} reduces, at some point, V_{BC} starts to become forward biased. Now, hole flux from B \rightarrow C increases exponentially from Law of Junction; to keep I_B constant, hole flux into emitter must be reduced, $\implies V_{BE}$ drops, $\implies I_C$ drops quickly.

Small Signal Model of BJT

(Next week we will be using BJT & MOSFET for amplifier circuits) Want to know the



small signal circuit model of BJT

$$1. \text{ Transconductance } g_m = \left. \frac{\delta i_c}{\delta V_{BE}} \right|_Q$$

$$I_C = I_s e^{qV_{BE}/kT} \implies g_m = \frac{q}{kT} I_s e^{qV_{BE}/kT} = \frac{I_C}{V_{th}}$$

Note, different from MOSFET: $g_m \simeq \sqrt{2 \frac{w}{L}} I_D$ (depends upon device size), but not for bipolar case.

2. Input resistance:

$$I_B = \frac{I_S}{\beta_F} e^{qV_{BE}/kT}$$

$$g_m = \frac{1}{\gamma_\pi} = \frac{\delta i_B}{\delta V_{BE}} = \frac{I_B}{V_{th}} = \frac{g_m}{\beta_F}$$

or $\gamma_\pi = \frac{\beta_F}{g_m}$

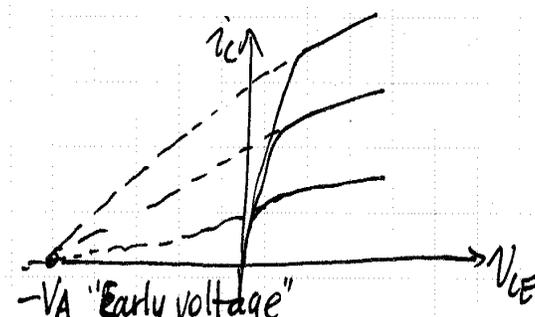
The input resistance of MOSFET is ∞ . In order to have a high input resistance for BJT, need high current gain β_F

Example: npn with $\beta_F = 150$, $I_c = \text{mA}$

$$g_m = \frac{I_c}{V_{th}} = \frac{1 \times 10^{-3} \text{ A}}{0.025 \text{ V}} = 40 \text{ mS}$$

$$g_\pi = \frac{g_m}{\beta_F} = \frac{40 \text{ mS}}{150} = 0.267 \text{ mS} \quad (\gamma_\pi = 3.7 \text{ k}\Omega)$$

3. Output resistance: Ebers-Moll model have perfect current source in FAR. Real characteristics show some increase in i_c with V_{CE}



$$g_o = \frac{\delta i_c}{\delta V_{CE}} \quad \text{where does } g_o \text{ come from?}$$

$$\text{In FAR: } I_c = I_S e^{qV_{BE}/kT} = \frac{qA_E n_i^2 D_{nB}}{N_{aB} w_B} e^{qV_{BE}/kT}$$

w_B shrinks as $|V_{BC}| \uparrow$, thus $I_c \uparrow$.

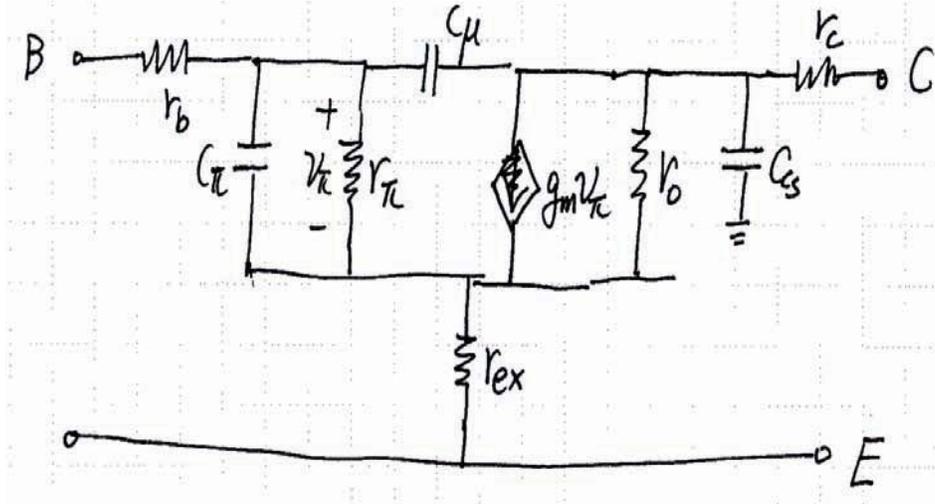
$$\text{Model: } g_o = \text{slope} = \frac{I_c}{V_{CE} + V_A} \simeq \frac{I_c}{V_A} \quad (V_A \gg V_{CE})$$

$$g_o = \frac{1}{\gamma_o} = \frac{I_c}{V_A}$$

Example: $I_c = 100 \mu\text{A}$, $V_A = 35 \text{ V}$, $\implies \gamma_o = 350 \text{ k}\Omega$

V_A increases with increasing base width and increasing base doping. This is also why N_{aB} usually $\gg N_{dC}$

Now what do we have so far? Need to add capacitances...



Junction Capacitance (depletion capacitance)

$$(B-E): C_{jE} = \sqrt{\frac{q\epsilon_s N_{aB} N_{dE}}{2(N_{aB} + N_{dE})(\phi_{BE} - V_{BE})}} \quad (\because N_{dE} \gg N_{aB})$$

$$(B-E): C_{jC} = \sqrt{\frac{q\epsilon_s N_{aB} N_{dC}}{2(N_{aB} + N_{dC})(\phi_{BC} - V_{BC})}} \approx \sqrt{\frac{q\epsilon_s}{2}} \frac{N_{dC}}{(\phi_{BC} - V_{BE})}$$

- Both are functions of bias
- Since $N_{aB} \gg N_{dC}$, $C_{jE} \gg C_{jC}$. C_{jC} is often called C_μ .

Diffusion Capacitance

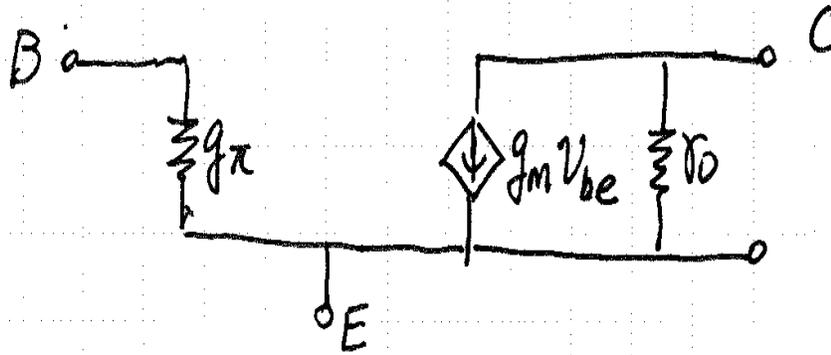
$$C_b = \frac{\delta}{\delta V_{BE}} |Q_{nB}|$$

$$|Q_{nB}| = \frac{1}{2} q A_E w_B n_{pBO} e^{qV_{BE}/kT}$$

$$= \frac{1}{2} w_B \left(\frac{w_B}{D_{nB}} \right) \left(\frac{q A_E D_{nB}}{w_B} \right) n_{pBO} e^{qV_{BE}/kT} = \left(\frac{w_B^2}{2D_{nB}} \right) I_c$$

$$C_b = \frac{\delta}{\delta V_{BE}} \left(\left(\frac{w_B^2}{2D_{nB}} \right) i_c \right) = \left(\frac{w_B^2}{2D_n} \right) g_m$$

$$\frac{w_B^2}{2D_n} = \tau_F \quad \text{base diffusion transit time}$$



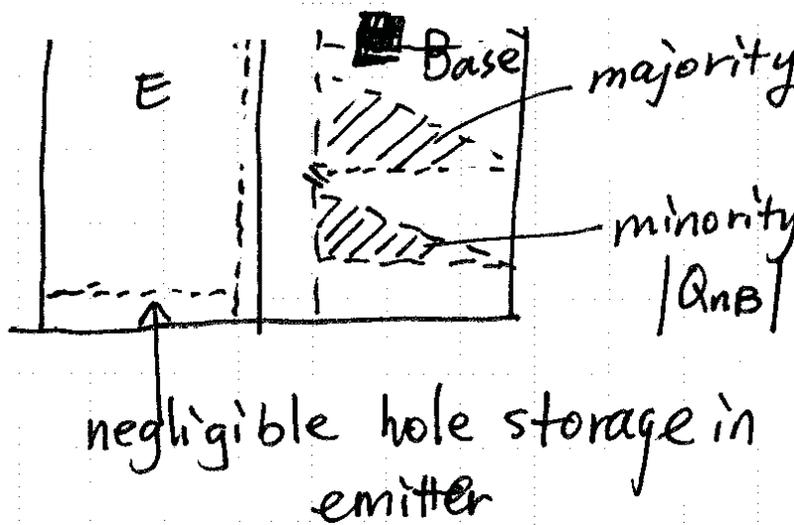
C_b is in between base and emitter:

$$C_b + C_{jE} = C_\pi$$

Add the following

- depletion capacitance: collector to bulk C_{CS}
- parasitic resistances: γ_b of base, γ_{ex} of emitter, γ_c of collector

Complete small signal model



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6.012 Microelectronic Devices and Circuits
Spring 2009

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