

## Recitation 6: p-n junction

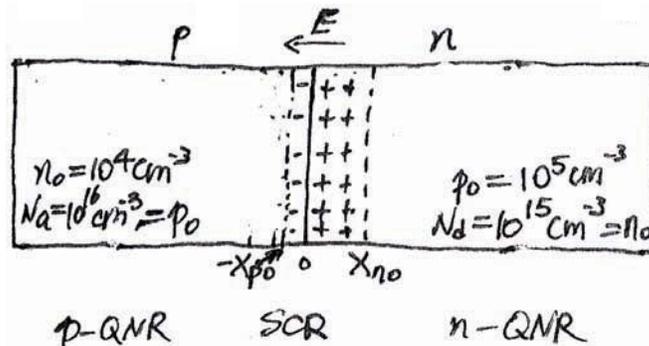
Today's agenda:

1. p-n junction under thermal equilibrium (T.E.):
  - space charge region (depletion layer)
  - quasi-neutral region
2. p-n junction under reverse bias
3. depletion capacitance

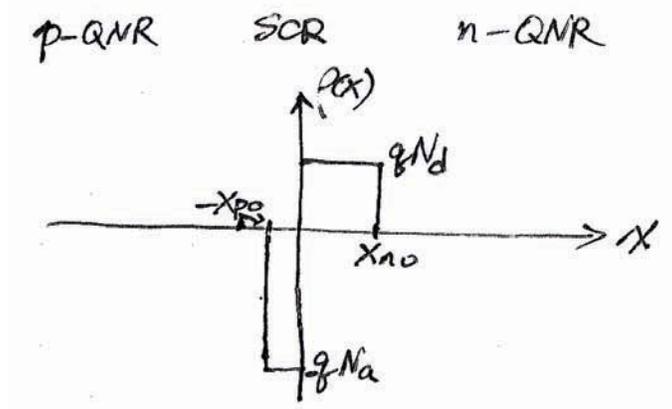
### p-n junction under T.E.

Yesterday we talked about p-n junction. We continue here:

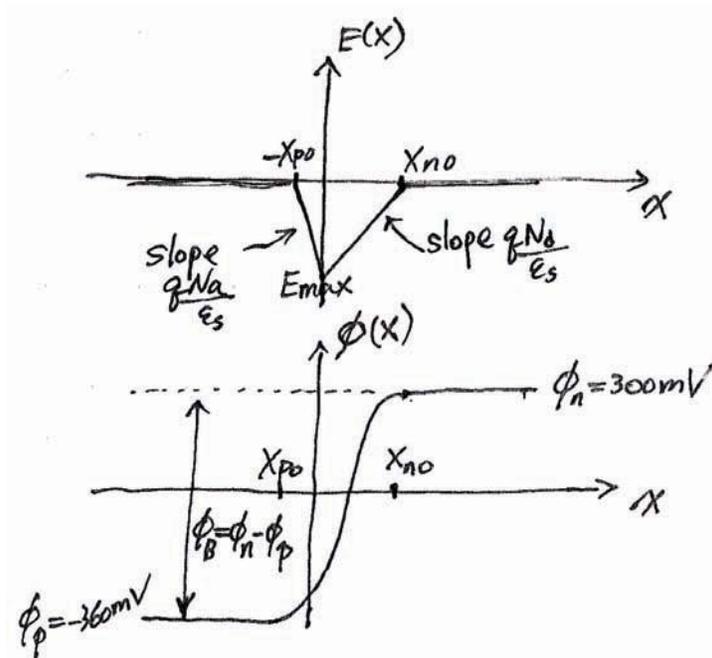
1. When we bring the p and n interface together, what will happen?  
Holes move from p-side to n-side and electrons diffuse from n-side to p-side, leaving behind space charges. The result is an electric field counteracting diffusion.



2. *Space Charge Region (SCR)*: is a depletion region.  $x_{no}, x_{po}, x_{no} + x_{po} = x_{do}$ .  
What is the charge density in SCR?  $q N_a(-)$  and  $q N_d(+)$   
And  $q N_a \cdot x_{po} = q N_d \cdot x_{no} \implies \frac{x_{po}}{x_{no}} = \frac{N_d}{N_a}$ .  
*Quasi-Neutral Region (QNR)*



3. Summaries of equations:



$$x_{no} = \sqrt{\frac{2\epsilon_s \phi_B N_a}{q(N_a + N_d) \cdot N_d}}$$

$$x_{po} = \sqrt{\frac{2\epsilon_s \phi_B N_d}{q(N_a + N_d) \cdot N_a}}$$

$x_{no}, x_{po}$  are determined by doping on *both* sides!

$$\begin{aligned} \text{Built-in potential } \phi_B &= \phi_n - \phi_p = \frac{kT}{q} \ln \frac{N_a N_d}{n_i^2} \\ x_{do} &= x_{no} + x_{po} = \sqrt{\frac{2\epsilon_s \phi_B}{q} \left( \frac{1}{N_a} + \frac{1}{N_d} \right)} \\ \text{Maximum field } |E_{\max}| &= |E_o| = \sqrt{\frac{2q\phi_B}{\epsilon_s} \frac{N_a \cdot N_d}{N_a + N_d}} \\ &= \frac{q N_a}{\epsilon_s} x_{po} = \frac{q N_d}{\epsilon_s} x_{no} \end{aligned}$$

4. If strongly asymmetric, the lowly doped side controls the electrostatics.

$$p^+n : x_{po} \ll x_{no} \simeq x_{do} = \sqrt{\frac{2\epsilon_s \phi_B}{q} \frac{1}{N_d}} \quad |E_o| \simeq \sqrt{\frac{2q\phi_B}{\epsilon_s} N_d}$$

## Examples

Now let us do some exercises before moving on to the next topic:

$N_d$ [cm <sup>-3</sup> ]	$N_a$ [cm <sup>-3</sup> ]	$x_{no}$	$x_{po}$	$E_o$ [V/cm]	$\phi_B$
10 <sup>16</sup>	10 <sup>16</sup>	216 nm	216 nm	$3.3 \times 10^4$ V/cm	720 mV
10 <sup>19</sup>	10 <sup>16</sup>	3.14 Å	341 nm	$5.26 \times 10^4$ V/cm	900 mV

$\phi_B$  is easier to calculate first.

$$1. \quad \phi_n = 360 \text{ mV}, \phi_p = -360 \text{ mV} \quad \phi_B = \phi_n - \phi_p = 720 \text{ mV}$$

$$\begin{aligned} x_{no} &= \sqrt{\frac{2\epsilon_s \phi_B}{q} \frac{N_a}{N_a + N_d}} \\ &= \sqrt{\frac{2 \times 1 \times 10^{-12} \text{ (F/cm)} \times 0.72 \text{ (V)}}{1.6 \times 10^{-19} \text{ (C)} \times 2 \times 10^{16} \text{ (cm}^{-3}\text{)}}} = \sqrt{4.66} \times 10^{-5} \text{ cm} = 216 \text{ nm} \\ x_{po} &= x_{no} \quad (\text{symmetric}) \\ |E_o| &= \sqrt{\frac{2q\phi_B}{\epsilon_s} \frac{N_a \cdot N_d}{N_a + N_d}} \\ &= \sqrt{2 \times \frac{1.6 \times 10^{-19} \text{ (C)} \times 0.72 \text{ (V)}}{1 \times 10^{-12} \text{ (F/cm)}} \times \frac{10^{16} \times 10^{16} \text{ (cm}^{-3}\text{)}^2}{2 \times 10^{16} \text{ (cm}^{-3}\text{)}}} = 3.3 \times 10^4 \text{ V/cm} \end{aligned}$$

$$2. \phi_B = \phi_n - \phi_p = 540 \text{ mV} + 360 \text{ mV} = 900 \text{ mV}$$

$$x_{no} = \sqrt{\frac{2\epsilon_s \phi_B}{q(N_a + N_d)} \frac{N_a}{N_d}} = 3.41 \text{ \AA} \quad (\text{really thin})$$

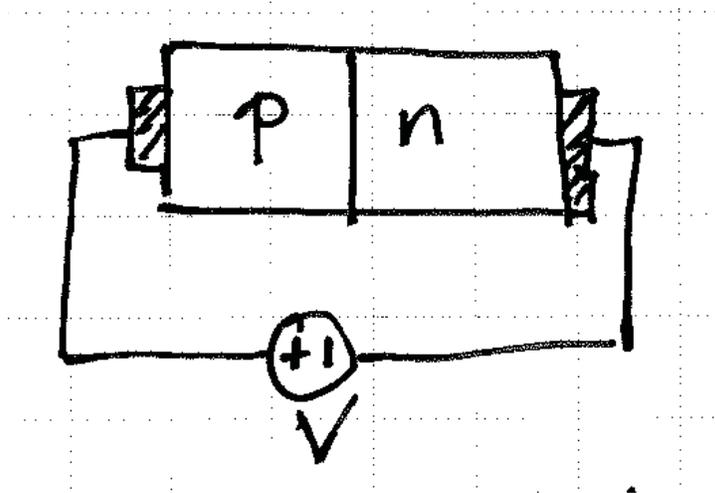
$$x_{po} = \sqrt{\frac{2\epsilon_s \phi_B}{q(N_a + N_d)} \frac{N_d}{N_a}} = \sqrt{\frac{2\epsilon_s \phi_B}{q} \frac{1}{N_a}} = 344 \text{ nm}$$

$$|E_o| = \frac{qN_a}{\epsilon_s x_{po}} = \frac{1.6 \times 10^{-19} \text{ (C)} \cdot 10^{16} \text{ (cm}^{-3}\text{)}}{1 \times 10^{-12} \text{ (F/cm)}} \times 3.41 \times 10^{-7} \text{ cm} = 5.26 \times 10^4 \text{ (V/cm)}$$

## Reverse Bias

Now what happens when we apply a reverse bias?

First: Bias convention for pn junction:



use n-side as reference:

$V > 0$ : forward bias

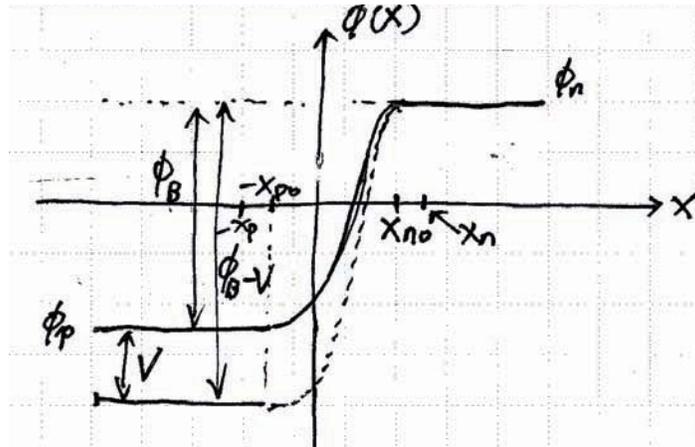
$V < 0$ : reverse bias

In your weblab experiment, when you have  $V < 0$ , what current do you measure?  $6 \times 10^{-15} \text{ A}$  (very small).

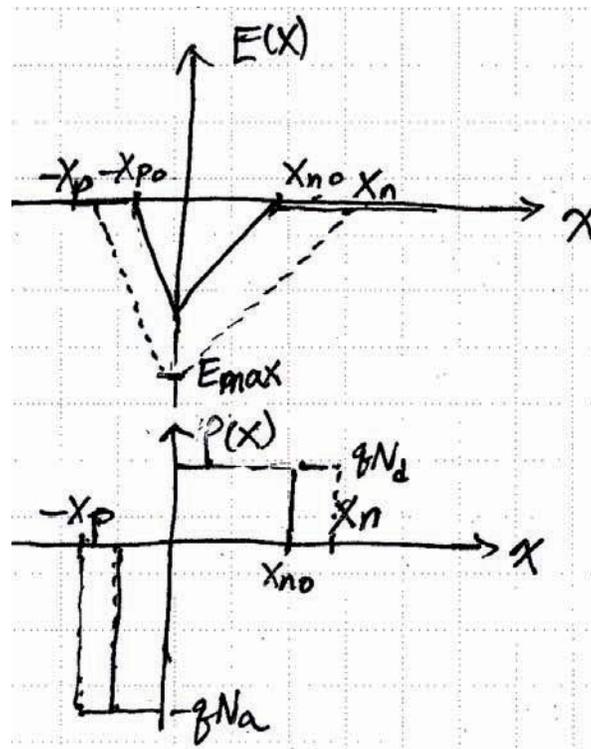
With such a small current, the voltage drop across the QNRs will be really small, safely to ignore.

Then, where is the voltage drop?  $\implies$  across SCR

Forward bias means  $V > 0$ , moving " $\phi_p$ " up and reverse bias means  $V < 0$ , moving " $\phi_p$ " down. Reverse bias  $V$ : barrier becomes  $\phi_B - V > \phi_B$  ( $V$  is negative, so barrier increases).



Now, what happened to the electric field? Still zero in QNRs. Larger barriers means higher E-field, and  $\therefore |E_{\max}|$  increases. This will mean we need more space charges to support the larger E-field. However, the charge density  $qN_a$  or  $qN_d$  can not change,  $\implies x_n, x_p$  widens.

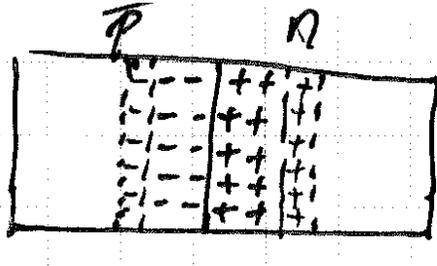


It turns out a straightforward way to do the new calculations is: replace  $\phi_B$  with  $\phi_B - V$ .  
So:

$$\begin{aligned} x_n &= x_{no} \sqrt{\frac{\phi_{rmB} - V}{\phi_B}} = x_{no} \sqrt{1 - \frac{V}{\phi_B}} \\ &= \sqrt{\frac{2\epsilon_s(\phi_B - V) N_a}{q(N_a + N_d) N_d}} \\ x_p &= x_{po} \sqrt{1 - \frac{V}{\phi_B}} \\ x_d &= x_{do} \sqrt{1 - \frac{V}{\phi_B}} \\ E_{\max} &= E_o \sqrt{1 - \frac{V}{\phi_B}} \end{aligned}$$

## The concept of depletion capacitance

Now let us think about the junction region a bit further.



1. It stores charges
2. By changing the voltage (bias) across the junction, the amount of charges stored there changes  $\implies$  Capacitor!
3. The charges that it stores:

$$Q_{jo} = |-q N_a \cdot x_{po}| = |q N_d \cdot x_{no}| = \sqrt{\frac{2q\epsilon_s N_a N_d}{N_a + N_d}} \cdot \phi_B$$

As voltage changes,  $Q_j$  becomes:

$$Q_j(V) = \sqrt{\frac{2q\epsilon_s N_a N_d}{N_a + N_d}} (\phi_B - V) = Q_{jo} \sqrt{1 - \frac{V}{\phi_B}}$$

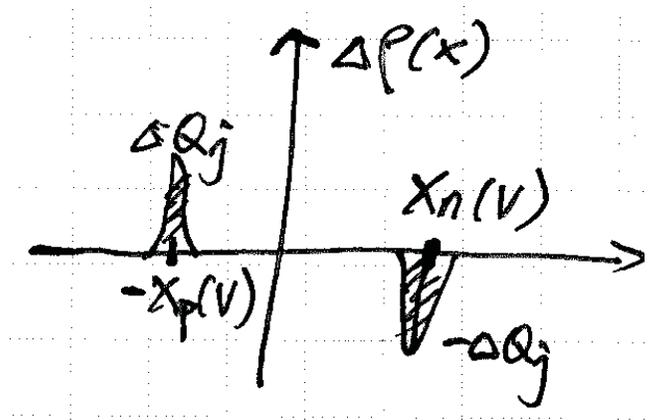
Previously, we had  $C = \frac{Q}{V}$ , but that was for constant capacitance, a more rigorous definition should be:

$$C = \frac{dQ}{dV}$$

Therefore,

$$\begin{aligned} C_j &= C_j(V_D) = \left. \frac{dQ_j}{dV} \right|_{V_D} = \left. Q_{j0} \frac{d\left(\sqrt{1 - \frac{V}{\phi_B}}\right)}{dV} \right|_{V_D} \\ &= \frac{+qN_a x_{po}}{2\sqrt{1 - \frac{V_D}{\phi_B}}} \cdot \frac{1}{\phi_B} \\ &= \frac{qN_a x_{po}}{2\phi_B} \cdot \frac{1}{\sqrt{1 - \frac{V_D}{\phi_B}}} \\ C_j(V=0) &= C_{j0} = \frac{qN_a x_{po}}{2\phi_B} \\ &= \frac{qN_a}{2\phi_B} \sqrt{\frac{2\epsilon_s \phi_B}{q(N_a + N_d)}} \times \frac{N_d}{N_a} = \sqrt{\frac{q\epsilon_s}{2\phi_B} \frac{N_a N_d}{N_a + N_d}} \\ \therefore x_{do} &= x_{no} + x_{po} = \sqrt{\frac{2\epsilon_s \phi_B}{q} \frac{N_a + N_d}{N_a N_d}} \\ \Rightarrow C_{j0} &= \frac{\epsilon_s}{x_{do}}, C_j(V_D) = \frac{C_{j0}}{\sqrt{1 - \frac{V_D}{\phi_B}}} \\ C_j(V_D) &= \frac{\epsilon_s}{x_{do} \sqrt{1 - \frac{V_D}{\phi_B}}} = \frac{\epsilon_s}{x_d} \end{aligned}$$

Can be considered as “parallel plate capacitor”.



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