

Recitation 5: Review on Electrostatics

This review is aimed at getting ready for p-n junction. Before the class, have an exercise on the Boltzman relationship:

1. If the doping of Si is 10^{16} Boron, what should be the corresponding electrostatic potential be? ($\phi = -6 \times 60 = -360$ mV)
2. If $\phi = 480$ mV, what is equilibrium electron concentration? ($\frac{480}{60} = 8, n_o = 10^{18} \text{ cm}^{-3}$)
3. If doping is $N_d = 10^{20} \text{ cm}^{-3}, \phi = ?$ ($\phi \neq 600$ mV, $\phi = \phi_{\max} = 550$ mV)

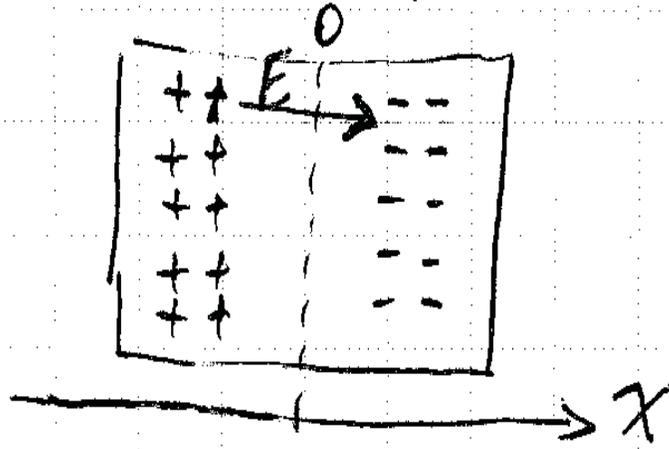
The following is the set of equations which relates $\rho(x), E(x), \phi(x)$ needed for this class (everything is 1D):

1. Relating charge density to the field:

$$\frac{dE}{dx} = \frac{\rho(x)}{\epsilon}$$

$$\text{or } E(x) - E(0) = \frac{1}{\epsilon} \int_0^x \rho(x') dx'$$

$\rho(x)$ = charge density in coulomb/cm³, ϵ is permittivity or dielectric constant of the material in F/cm, $\epsilon_o(\text{vacuum}) = 8.85 \times 10^{-14}$ F/cm.



2. Relating electrostatic potential to the field:

$$\frac{d\phi}{dx} = -E(x)$$

$$\text{or } \phi(x) - \phi(0) = - \int_0^x E(x') dx'$$

$E(x)$ has units V/cm and $\phi(x)$ has units of V or mV.

3. Boundary conditions:

- Continuation of potential at a boundary (infinite field inside semiconductor not possible)

$$\phi(x_b^-) = \phi(x_b^+)$$

- E-field at the boundary can (usually) have a jump, due to:
 - Materials change:

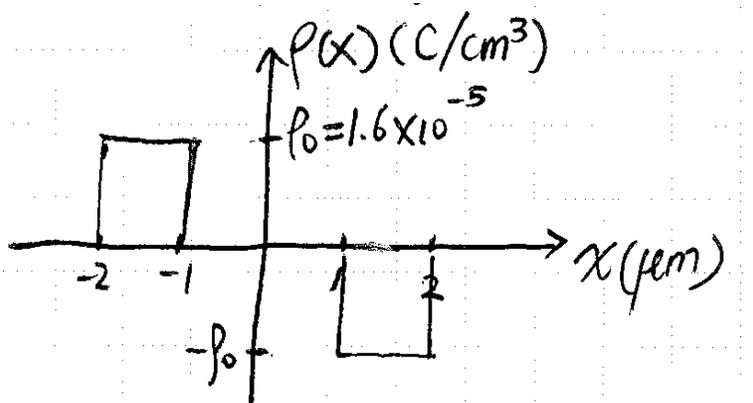
$$\begin{aligned} \int_{-\Delta}^{\Delta} d(\epsilon E(x)) &= \epsilon_2 E(x = \Delta) - \epsilon_1 E(x = -\Delta) \\ &= \int_{-\Delta}^{\Delta} \rho(x) dx = 0 \text{ no charge at interface} \\ E(0^+) &= \frac{\epsilon_1}{\epsilon_2} E(0^-) \end{aligned}$$

- A sheet of charge at the interface:

$$\begin{aligned} \int_{-\Delta}^{\Delta} d(\epsilon E(x)) &= \epsilon_2 E(x = \Delta) - \epsilon_1 E(x = -\Delta) \\ &= \int_{-\Delta}^{\Delta} \rho(x) dx = Q \text{ coulomb/cm}^2 \\ E(0^+) &= \frac{\epsilon_1}{\epsilon_2} E(0^-) + \frac{Q}{\epsilon_2} \end{aligned}$$

Examples

Like in Figure 1



Then, let us work out $E(x)$, by splitting calculations into 4 regions:

1. $E(x \leq -2 \mu\text{m}) = E(x \geq 2 \mu\text{m})$ (\because no charge outside these regions)

2. For $-2 \mu\text{m} < x < -1 \mu\text{m}$:

$$E(x) - E(-2 \mu\text{m}) = \frac{1}{\epsilon_s} \int_{-2}^x \rho(x) dx = \frac{1}{\epsilon_s} \rho_o \cdot (x + 2)$$

$$\rho_o = 1.6 \times 10^{-5} \text{ C/cm}^3 = 1.6 \times 10^{-17} \text{ C/cm}^3$$

$$\text{Particularly, } E(-1 \mu\text{m}) = \frac{1}{\epsilon_s} \rho_o \times 1 \mu\text{m}$$

$$\text{Since } \epsilon_s = 1 \times 10^{-12} \text{ F/cm,}$$

$$E(-1 \mu\text{m}) = \frac{1.6 \times 10^{-5} \times 10^{-4} \text{ C/cm}^3 \times \text{cm}}{1 \times 10^{-12} \text{ F/cm}} = 1.6 \times 10^3 \text{ V/cm}$$

3. For $-1 < x < 1 \mu\text{m}$:

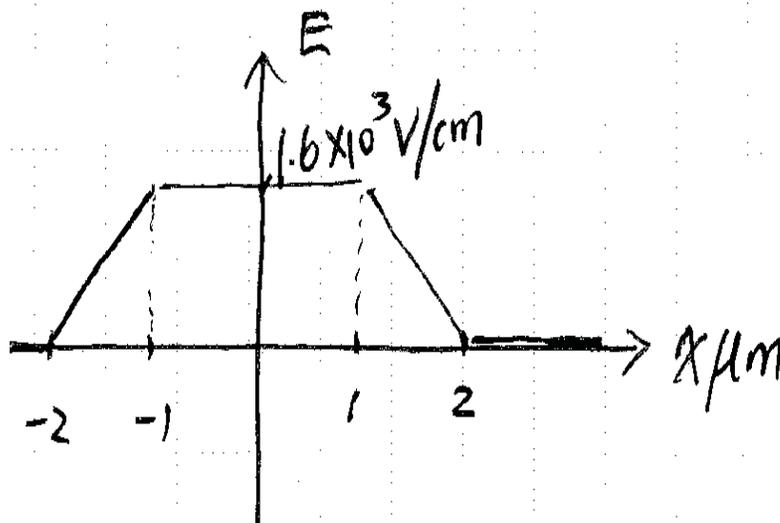
$$\rho(x) = 0, \text{ same dielectric material} \implies E \text{ constant}$$

$$E(x) = E(-1 \mu\text{m}) = 1.6 \times 10^3 \text{ V/cm}$$

4. For $1 < x < 2 \mu\text{m}$:

$$E(x) - E(1 \mu\text{m}) = \frac{1}{\epsilon_s} \int_1^x \rho(x) dx = -\frac{1}{\epsilon_s} \rho_o (x - 1)$$

$$\text{particularly, } E(2 \mu\text{m}) = E(1 \mu\text{m}) - \frac{1}{\epsilon_s} \rho_o (x - 1) \Big|_{x=2 \mu\text{m}} = 0$$



Now for the electrostatic potential:

As we do the integration, the results will be relative. Then we can use the doping of Si to find the actual value. For this thought-example, let us make $\phi(0) = 0$.

For $0 < x < 1 \mu\text{m}$, $E(x) = \text{constant} = 1.6 \times 10^3 \text{ V/cm}$.

$$\begin{aligned}\phi(x) - \phi(0) &= - \int_0^x E(x') dx' = -1.6 \times 10^3 \cdot x \text{ (V/cm} \cdot \mu\text{m)} \\ \text{particularly } \phi(1 \mu\text{m}) &= -0.16\text{V}\end{aligned}$$

For $1 < x < 2 \mu\text{m}$,

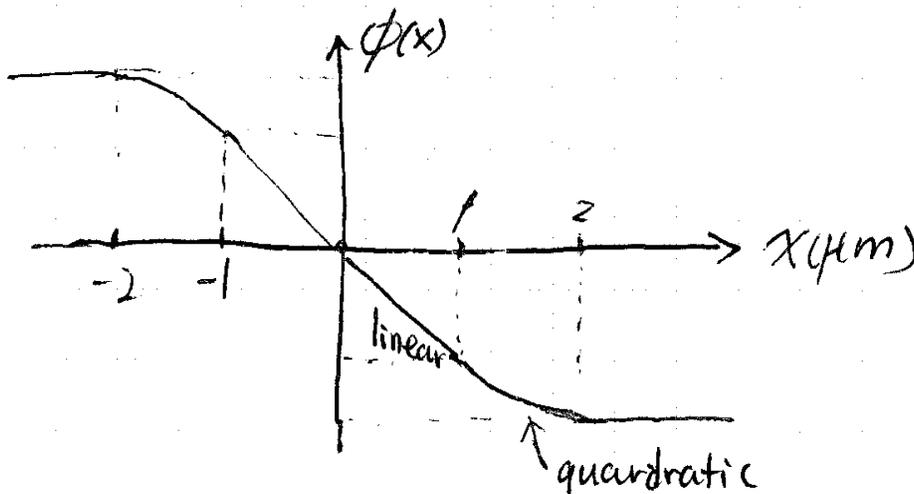
$$\begin{aligned}\phi(x) - \phi(1 \mu\text{m}) &= - \int_1^x E(x') dx' \\ &= - \int_1^x (1.6 \times 10^3 \text{ V/cm} - \frac{1}{\epsilon_s} \rho_o (x' - 1)) dx' \\ &= -1.6 \times 10^3 \text{ V/cm} (x - 1) + \frac{1}{2\epsilon_s} \rho_o (x - 1)^2 \text{ quadratic}\end{aligned}$$

$$\begin{aligned}\text{particularly, } \phi(2 \mu\text{m}) &= \phi(1 \mu\text{m}) - 1.6 \times 10^3 \text{ V/cm} \cdot (1 \times 10^{-4}) \text{ cm} + \frac{1}{2} \frac{1.6 \times 10^{-5} \text{ C/cm}^3}{1 \times 10^{-12} \text{ F/cm}} \cdot (1 \times 10^{-8} \text{ cm}^2) \\ &= -0.16\text{V} - 0.08\text{V} = -0.24\text{V}\end{aligned}$$

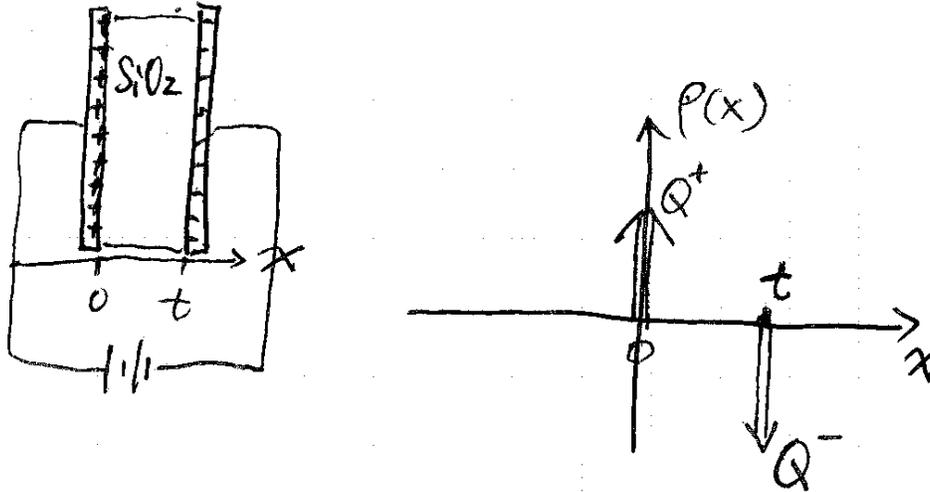
For $x > 2 \mu\text{m}$:

$\therefore E(x) = 0$ electrostatic potential will be constant

Similarly, we can work out the other side:



Parallel Plate Capacitor



Two sheets of charge at the interface of the capacitor Q^+/Q^-
 In 1D, this can be modeled as δ function

Now let us consider the electric field:

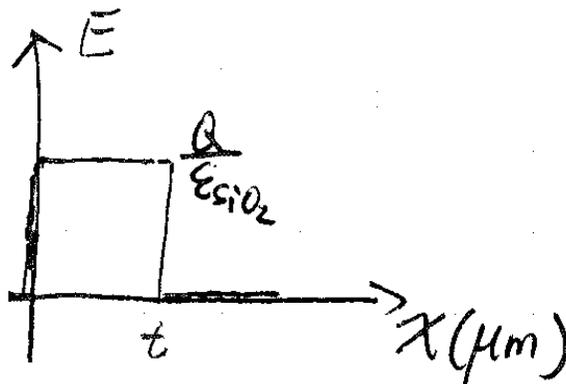
Inside the metal, no charge, no field, $\implies E(x < 0) = E(x > t_d) = 0$:

$$\epsilon_{\text{SiO}_2} - 0 = \int_{0^-}^{0^+} \rho(x') dx' = \int_{0^-}^{0^+} \delta(x) dx = Q$$

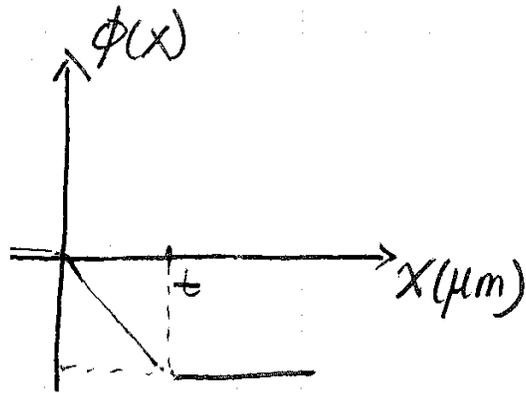
$$\therefore E(0^+) = \frac{Q}{\epsilon_{\text{SiO}_2}}$$

$$\epsilon_{\text{SiO}_2} = 3.45 \times 10^{-13} \text{ F/cm}$$

As there is no charge inside SiO_2 region, E is constant $E(x) = \frac{Q}{\epsilon_{\text{SiO}_2}}$ for $0 < x < t$



If we use $\phi(0) = 0$ again,



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