## Lecture 4

### PN Junction and MOS Electrostatics(I)

# Semiconductor Electrostatics in Thermal Equilibrium

### Outline

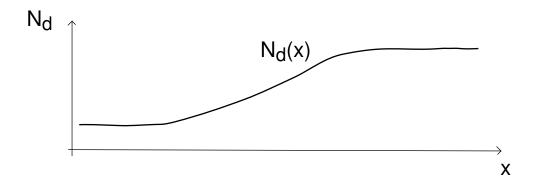
- Non-uniformly doped semiconductor in thermal equilibrium
- Relationships between potential,  $\phi(x)$  and equilibrium carrier concentrations,  $p_0(x)$ ,  $n_0(x)$ 
  - -Boltzmann relations & "60 mV Rule"
- Quasi-neutral situation

### **Reading Assignment:**

Howe and Sodini; Chapter 3, Sections 3.1-3.2

# 1. Non-uniformly doped semiconductor in thermal equilibrium

Consider a piece of n-type Si in thermal equilibrium with non-uniform dopant distribution:

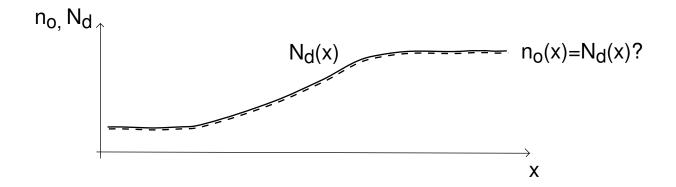


What is the resulting electron concentration in thermal equilibrium?

n-type ⇒ lots of electrons, few holes ⇒ focus on electrons

**OPTION 1:** electron concentration follows doping concentration EXACTLY ⇒

$$n_o(x) = N_d(x)$$

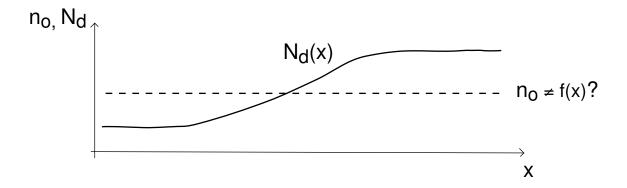


Gradient of electron concentration

- ⇒ net electron diffusion
- $\Rightarrow$  not in thermal equilibrium!

**OPTION 2:** electron concentration uniform in space

$$n_0(x) = n_{ave} \neq f(x)$$



Think about space charge density:

$$\rho(x) \approx q[N_d(x) - n_o(x)]$$

If 
$$N_d(x) \neq n_o(x)$$
  
 $\Rightarrow \rho(x) \neq 0$   
 $\Rightarrow$  electric field  
 $\Rightarrow$  net electron drift

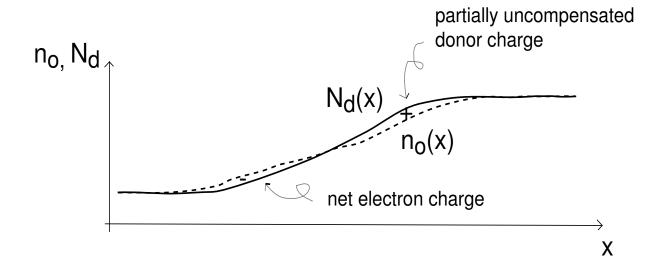
⇒ not in thermal equilibrium!

**OPTION 3:** Demand that  $J_n = 0$  in thermal equilibrium at every x ( $J_p = 0$  too)

Diffusion precisely balances Drift

$$J_n(x) = J_n^{drift}(x) + J_n^{diff}(x) = 0$$

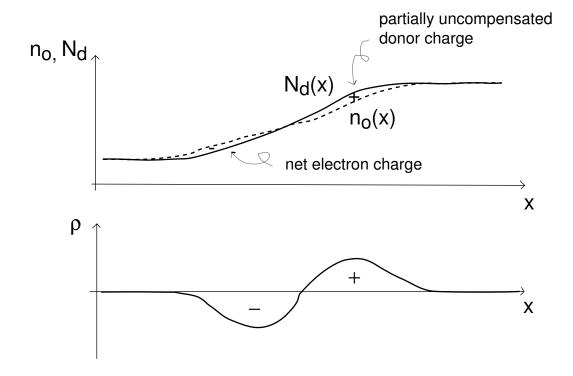
What is  $n_o(x)$  that satisfies this condition?



Let us examine the electrostatics implications of  $n_o(x) \neq N_d(x)$ 

# **Space charge density**

$$\rho(x) = q[N_d(x) - n_o(x)]$$



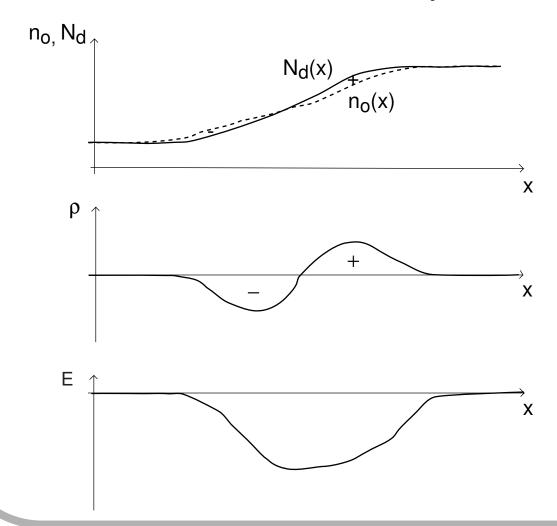
### **Electric Field**

Poisson's equation:

$$\frac{d\mathbf{E}}{d\mathbf{x}} = \frac{\rho(\mathbf{x})}{\varepsilon_{\mathbf{s}}}$$

Integrate from x = 0:

$$E(x) - E(0) = \frac{1}{\varepsilon_s} \int_{0}^{x} \rho(x') dx'$$



### **Electrostatic Potential**

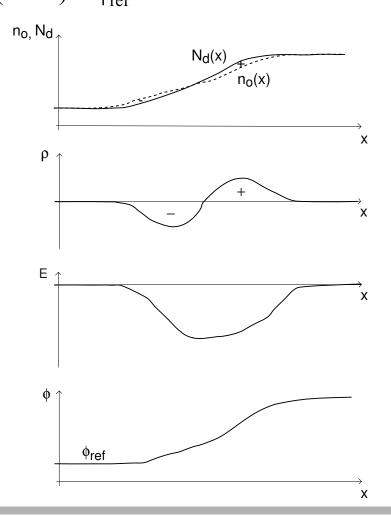
$$\frac{\mathrm{d}\phi}{\mathrm{d}x} = -\mathbf{E}(\mathbf{x})$$

Integrate from x=0:

$$\phi(x) - \phi(0) = -\int_{0}^{x} E(x')dx'$$

Need to select reference

(physics is in the potential difference, not in absolute value!); Select  $\varphi(x=0)=\varphi_{ref}$ 



# 2. Relationships between potential, $\phi(x)$ and equilibrium carrier concentrations, $p_0(x)$ , $n_0(x)$ (Boltzmann relations)

$$J_{n} = 0 = qn_{o}\mu_{n}E + qD_{n}\frac{dn_{o}}{dx}$$
$$\frac{\mu_{n}}{D_{n}} \cdot \frac{d\phi}{dx} = \frac{1}{n_{o}} \cdot \frac{dn_{o}}{dx}$$

Using Einstein relation:

$$\frac{q}{kT} \bullet \frac{d\phi}{dx} = \frac{d(\ln n_o)}{dx}$$

Integrate:

$$\frac{q}{kT}\left(\phi - \phi_{ref}\right) = \ln n_o - \ln n_{o,ref} = \ln \frac{n_o}{n_{o,ref}}$$

$$n_o = n_{o,ref} \exp \left| \frac{q(\phi - \phi_{ref})}{kT} \right|$$

Any reference is good

In 6.012,  $\phi_{ref} = 0$  at  $n_{o,ref} = n_i$ 

Then:

$$n_o = n_i e^{q\phi/kT}$$

If we do same with holes (starting with  $J_p = 0$  in thermal equilibrium, or simply using  $n_o p_o = n_i^2$ );

$$p_o = n_i e^{-q\phi/kT}$$

We can re-write as:

$$\phi = \frac{kT}{q} \bullet \ln \frac{n_o}{n_i}$$

and

$$\phi = -\frac{kT}{q} \bullet \ln \frac{p_o}{n_i}$$

## "60 mV" Rule

At room temperature for Si:

$$\phi = (25 \,\mathrm{m}) \cdot \ln \frac{\mathrm{n_0}}{\mathrm{n_i}} = (25 \,\mathrm{mV}) \cdot \ln(10) \cdot \log \frac{\mathrm{n_0}}{\mathrm{n_i}}$$

Or

$$\phi \approx (60 \,\mathrm{m}) \cdot \log \frac{\mathrm{n_0}}{\mathrm{n_i}}$$

### **EXAMPLE 1**:

$$n_0 = 10^{18} \text{cm}^{-3} \implies \phi = (60 \text{m}) \times 8 = 480 \text{mV}$$

## "60 mV" Rule: contd.

With holes:

$$\phi = -(25 \,\mathrm{m}) \cdot \ln \frac{p_0}{n_i} = -(25 \,\mathrm{m}) \cdot \ln (10) \cdot \log \frac{p_0}{n_i}$$

Or

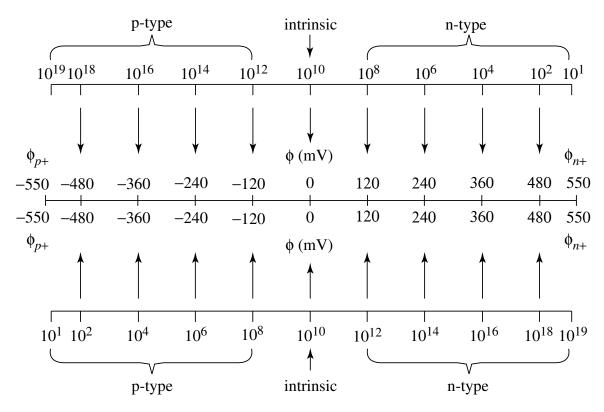
$$\phi \approx -(60 \,\mathrm{m}) \cdot \log \frac{\mathrm{p_0}}{\mathrm{n_i}}$$

#### **EXAMPLE 2**:

$$n_o = 10^{18} \text{ cm}^{-3} \Rightarrow p_o = 10^2 \text{ cm}^{-3}$$
  
  $\Rightarrow \phi = -(60 \text{ m}) \times -8 = 480 \text{ mV}$ 

# Relationship between $\phi$ , $n_o$ and $p_o$ :

 $p_o$  equilibrium hole concentration (cm<sup>-3</sup>)



 $n_o$  equilibrium electron concentration (cm<sup>-3</sup>)

Note: φ cannot exceed 550 mV or be smaller than -550 mV. (Beyond this point different physics come into play.)

**Example 3:** Compute potential difference in thermal equilibrium between region where  $n_o = 10^{17}$  cm<sup>-3</sup> and  $n_o = 10^{15}$  cm<sup>-3</sup>.

$$\phi(n_0 = 10^{17} \text{ cm}^{-3}) = 60 \times 7 = 420 \text{ mV}$$

$$\phi(n_o = 10^{15} cm^{-3}) = 60 \times 5 = 300 mV$$

$$\phi(n_o = 10^{17} \ cm^{-3}) - \phi(n_o = 10^{15} \ cm^{-3}) = 120 \ mV$$

**Example 4:** Compute potential difference in thermal equilibrium between region where  $p_o = 10^{20}$  cm<sup>-3</sup> and  $p_o = 10^{16}$  cm<sup>-3</sup>.

$$\phi(p_o = 10^{20} cm^{-3}) = \phi_{\text{max}} = -550 mV$$

$$\phi(p_0 = 10^{16} cm^{-3}) = -60 \times 6 = -360 mV$$

$$\phi(p_o = 10^{20} \text{ cm}^{-3}) - \phi(p_o = 10^{16} \text{ cm}^{-3}) = -190 \text{ mV}$$

# 3. Quasi-neutral situation

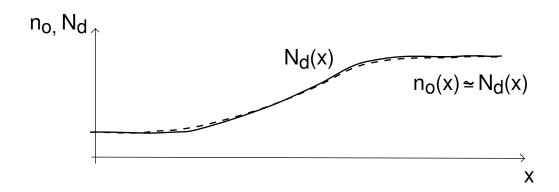
If  $N_d(x)$  changes slowly with  $x \Rightarrow n_o(x)$  also changes slowly with x. WHY?

Small dn<sub>o</sub>/dx implies a small diffusion current. We do not need a large drift current to balance it.

Small drift current implies a small electric field and therefore a small space charge

Then: 
$$n_o(x) \approx N_d(x)$$

 $n_o(x)$  tracks  $N_d(x)$  well  $\Rightarrow$  minimum space charge  $\Rightarrow$  semiconductor is quasi-neutral



# What did we learn today?

### **Summary of Key Concepts**

- It is possible to have an electric field inside a semiconductor in thermal equilibrium
  - $\Rightarrow Non-uniform doping distribution.$
- In thermal equilibrium, there is a fundamental relationship between the  $\phi(x)$  and the equilibrium carrier concentrations  $n_o(x)$  &  $p_o(x)$ 
  - Boltzmann relations (or "60 mV Rule").
- In a slowly varying doping profile, majority carrier concentration tracks well the doping concentration.

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6.012 Microelectronic Devices and Circuits Spring 2009

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