

# Lecture 3

## Semiconductor Physics (II)

### Carrier Transport

#### Outline

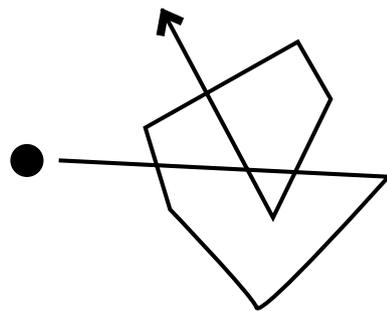
- Thermal Motion
- Carrier Drift
- Carrier Diffusion

**Reading Assignment:**  
Howe and Sodini; Chapter 2, Sect. 2.4-2.6

# 1. Thermal Motion

In thermal equilibrium, carriers are not sitting still:

- Undergo collisions with vibrating Si atoms (*Brownian motion*)
- Electrostatically interact with each other and with ionized (charged) dopants



Characteristic time constant of thermal motion:

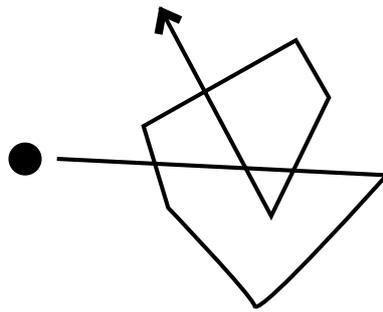
⇒ **mean free time between collisions**

$$\tau_c \equiv \textit{collision time} [s]$$

In between collisions, carriers acquire high velocity:

$$\mathbf{v}_{th} \equiv \textit{thermal velocity} [cms^{-1}]$$

.... but get nowhere!



Characteristic length of thermal motion:

$$\lambda \equiv \textit{mean free path} \text{ [cm]}$$

$$\lambda = v_{th} \tau_c$$

Put numbers for Si at room temperature:

$$\tau_c \approx 10^{-13} \text{ s}$$

$$v_{th} \approx 10^7 \text{ cm s}^{-1}$$

$$\Rightarrow \lambda \approx 0.01 \text{ } \mu\text{m}$$

For reference, state-of-the-art production MOSFET:

$$L_g \approx 0.1 \text{ } \mu\text{m}$$

$\Rightarrow$  **Carriers undergo many collisions as they travel through devices**

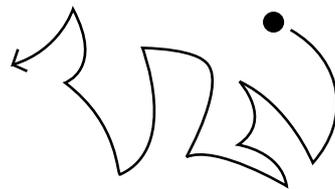
## 2. Carrier Drift

Apply electric field to semiconductor:

**E**  $\equiv$  electric field [ $\text{V cm}^{-1}$ ]

$\Rightarrow$  net force on carrier

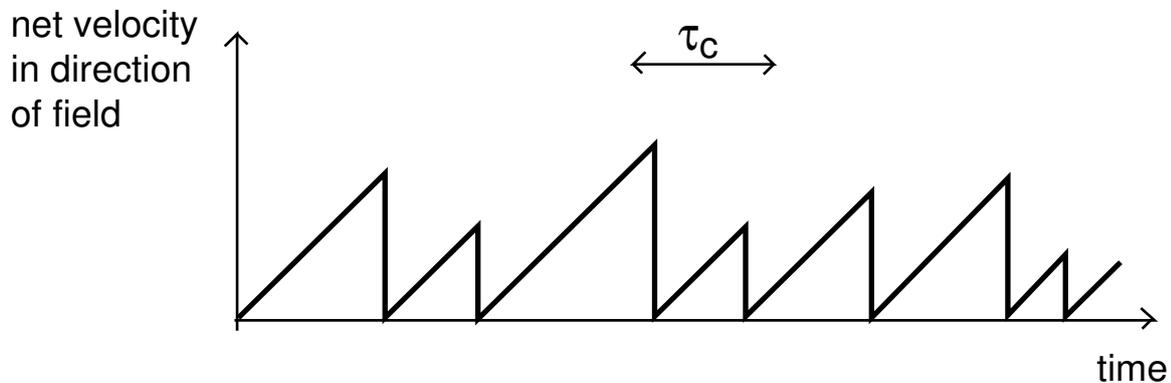
$$F = \pm qE$$



Between collisions, carriers accelerate in the direction of the electrostatic field:

$$\mathbf{v}(\mathbf{t}) = \mathbf{a} \bullet \mathbf{t} = \pm \frac{q\mathbf{E}}{m_{n,p}} \mathbf{t}$$

But there is (on the average) a collision every  $\tau_c$  and the velocity is randomized:



The average net velocity in direction of the field:

$$\bar{v} = v_d = \pm \frac{qE}{2m_{n,p}} \tau_c = \pm \frac{q \tau_c}{2m_{n,p}} E$$

This is called **drift velocity** [ $\text{cm s}^{-1}$ ]

Define:

$$\mu_{n,p} = \frac{q \tau_c}{2m_{n,p}} \equiv \text{mobility} [\text{cm}^2 \text{V}^{-1} \text{s}^{-1}]$$

**Then, for electrons:**

$$v_{dn} = -\mu_n E$$

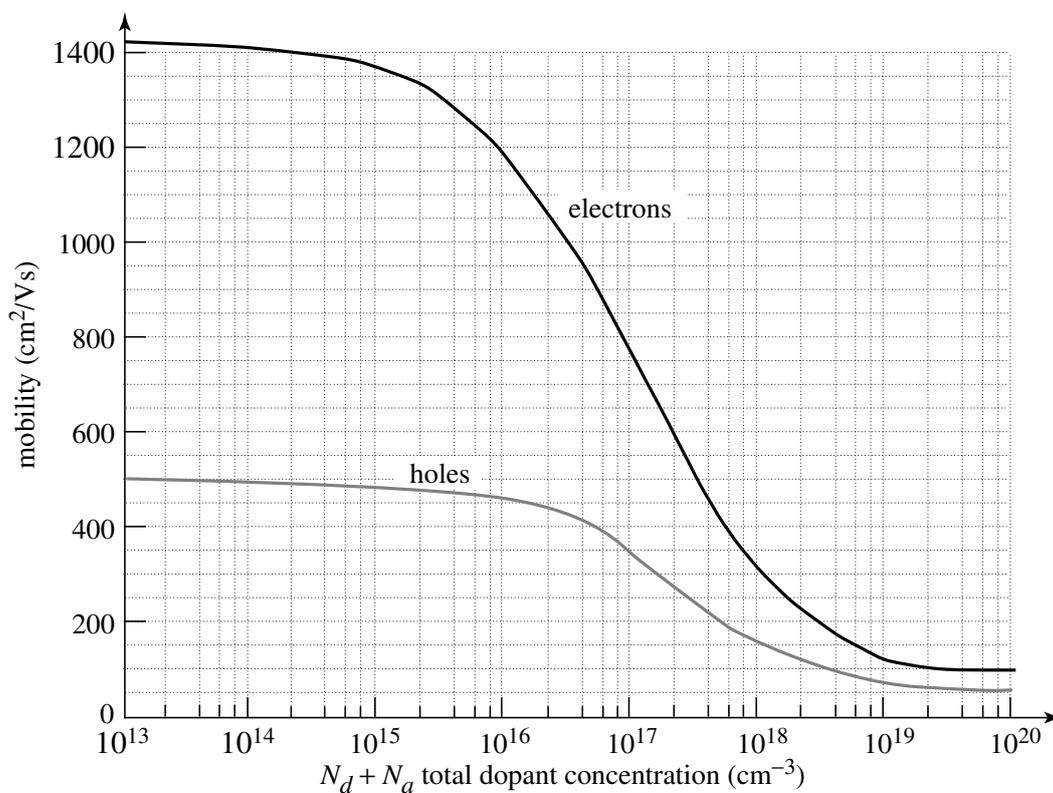
**and for holes:**

$$v_{dp} = \mu_p E$$

## Mobility - is a measure of ease of carrier drift

- If  $\tau_c \uparrow$ , longer time between collisions  $\Rightarrow \mu \uparrow$
- If  $m \downarrow$ , “lighter” particle  $\Rightarrow \mu \uparrow$

At room temperature, mobility in Si depends on doping:



- For low doping level,  $\mu$  is limited by collisions with lattice. As Temp  $\rightarrow$  **INCREASES**;  $\mu \rightarrow$  **DECREASES**
- For medium doping and high doping level,  $\mu$  limited by collisions with ionized impurities
- Holes “heavier” than electrons
  - For same doping level,  $\mu_n > \mu_p$

# Drift Current

Net velocity of charged particles  $\Rightarrow$  electric current:

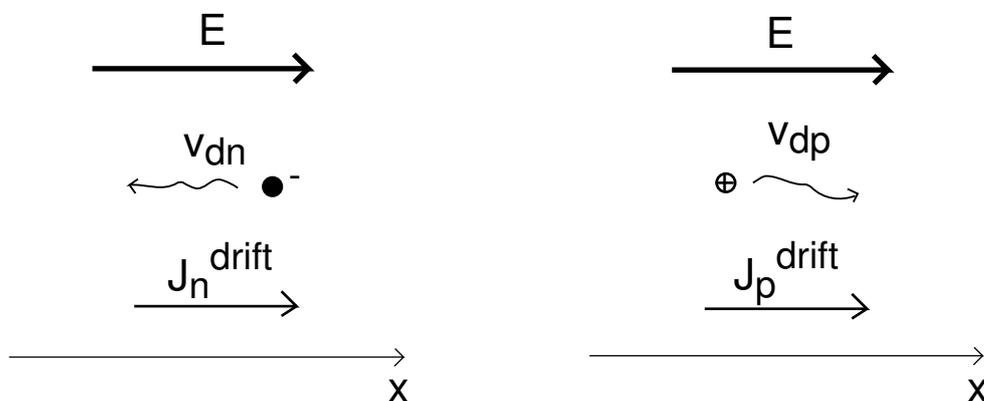
*Drift current density*  $\propto$  carrier drift velocity  
 $\propto$  carrier concentration  
 $\propto$  carrier charge

Drift current densities:

$$J_n^{drift} = -qnv_{dn} = qn\mu_n E$$

$$J_p^{drift} = qp v_{dp} = qp\mu_p E$$

Check signs:



## Total Drift Current Density :

$$\mathbf{J}^{\text{drift}} = \mathbf{J}_n^{\text{drift}} + \mathbf{J}_p^{\text{drift}} = q(n\mu_n + p\mu_p)\mathbf{E}$$

Has the form of *Ohm's Law*

$$\mathbf{J} = \sigma\mathbf{E} = \frac{\mathbf{E}}{\rho}$$

Where:

$\sigma \equiv$  conductivity [ $\Omega^{-1} \cdot \text{cm}^{-1}$ ]

$\rho \equiv$  resistivity [ $\Omega \cdot \text{cm}$ ]

Then:

$$\sigma = \frac{1}{\rho} = q(n\mu_n + p\mu_p)$$

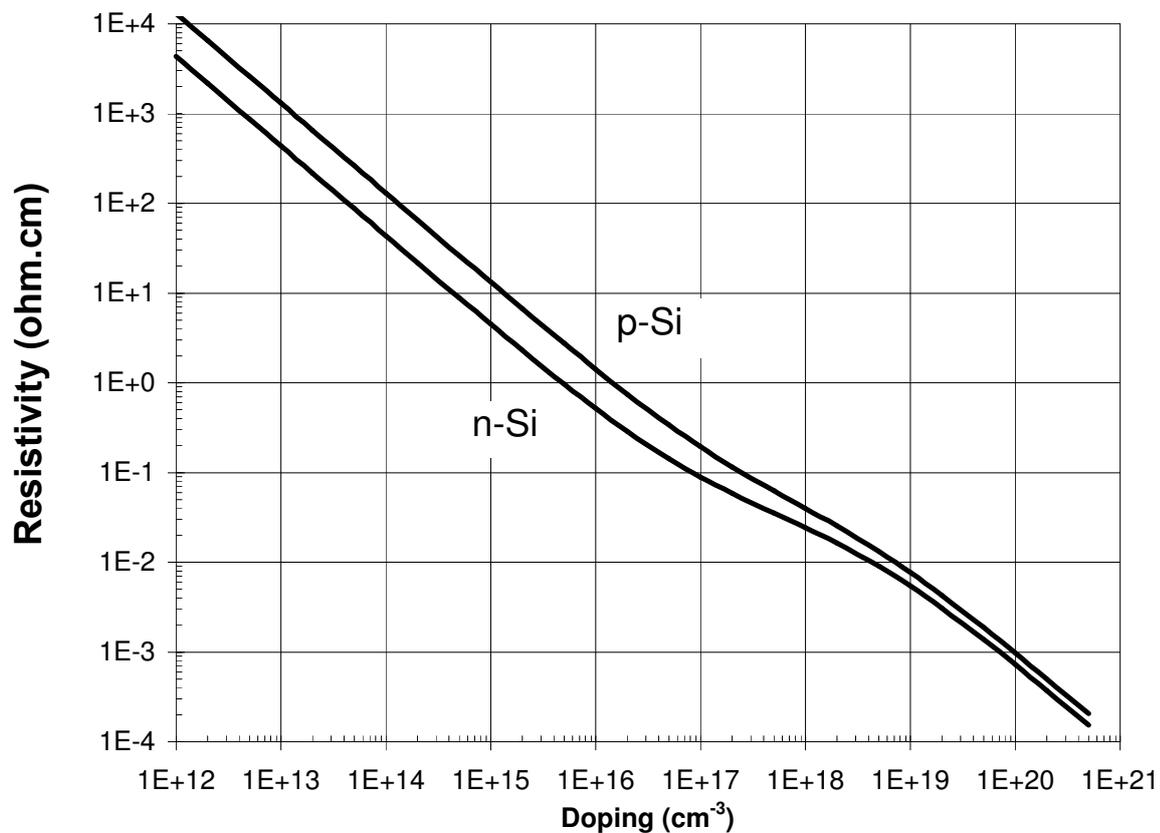
**Resistivity** is commonly used to specify the doping level

- In n-type semiconductor:

$$\rho_n \approx \frac{1}{qN_d\mu_n}$$

- In p-type semiconductor:

$$\rho_p \approx \frac{1}{qN_a\mu_p}$$



## Numerical Example:

Si with  $N_d = 3 \times 10^{16} \text{ cm}^{-3}$  at room temperature

$$\mu_n \approx 1000 \text{ cm}^2 / \text{V} \cdot \text{s}$$

$$\rho_n \approx 0.21 \Omega \cdot \text{cm}$$

$$n \approx 3 \times 10^{16} \text{ cm}^{-3}$$

Apply  $E = 1 \text{ kV/cm}$

$$v_{dn} \approx -10^6 \text{ cm/s} \ll v_{th}$$

$$J_n^{drift} \approx qn v_{dn} = qn \mu_n E = \sigma E = \frac{E}{\rho}$$

$$J_n^{drift} \approx 4.8 \times 10^3 \text{ A/cm}^2$$

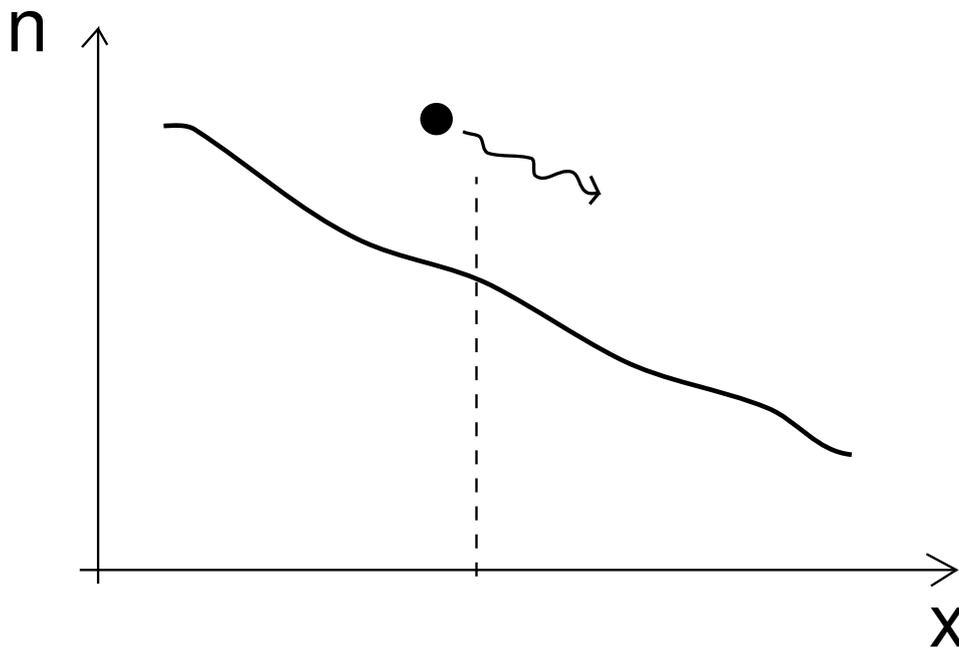
Time to drift through  $L = 0.1 \mu\text{m}$

$$t_d = \frac{L}{v_{dn}} = 10 \text{ ps}$$

**fast!**

### 3. Carrier Diffusion

**Diffusion** = particle movement (flux) in response to concentration gradient



#### Elements of diffusion:

- A medium (*Si Crystal*)
- A gradient of particles (*electrons and holes*) inside the medium
- Collisions between particles and medium send particles off in random directions
  - Overall result is to erase gradient

# Fick's first law-

## Key diffusion relationship

*Diffusion flux  $\propto$  - concentration gradient*

**Flux**  $\equiv$  number of particles crossing a unit area per unit time [ $\text{cm}^{-2} \cdot \text{s}^{-1}$ ]

**For Electrons:**

$$F_n = -D_n \frac{dn}{dx}$$

**For Holes:**

$$F_p = -D_p \frac{dp}{dx}$$

$D_n \equiv$  electron diffusion coefficient [ $\text{cm}^2 \text{s}^{-1}$ ]

$D_p \equiv$  hole diffusion coefficient [ $\text{cm}^2 \text{s}^{-1}$ ]

D measures the *ease* of carrier diffusion in response to a concentration gradient:  $D \uparrow \Rightarrow F^{\text{diff}} \uparrow$

D limited by vibration of lattice atoms and ionized dopants.

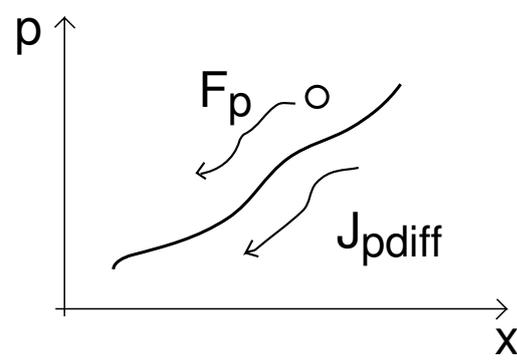
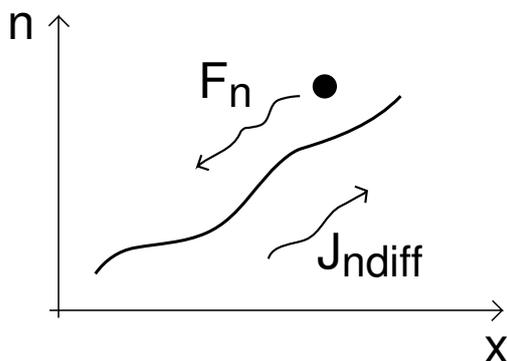
# Diffusion Current

*Diffusion current density = charge  $\times$  carrier flux*

$$J_n^{diff} = qD_n \frac{dn}{dx}$$

$$J_p^{diff} = -qD_p \frac{dp}{dx}$$

Check signs:



# Einstein relation

At the core of drift and diffusion is same physics:  
collisions among particles and medium atoms  
⇒ there should be a relationship between  $D$  and  $\mu$

**Einstein relation** [will not derive in 6.012]

$$\frac{D}{\mu} = \frac{kT}{q}$$

In semiconductors:

$$\frac{D_n}{\mu_n} = \frac{kT}{q} = \frac{D_p}{\mu_p}$$

$kT/q \equiv$  thermal voltage

At room temperature:

$$\frac{kT}{q} \approx 25 \text{ mV}$$

For example: for  $N_d = 3 \times 10^{16} \text{ cm}^{-3}$

$$\mu_n \approx 1000 \text{ cm}^2 / V \cdot s \Rightarrow D_n \approx 25 \text{ cm}^2 / s$$

$$\mu_p \approx 400 \text{ cm}^2 / V \cdot s \Rightarrow D_p \approx 10 \text{ cm}^2 / s$$

# Total Current Density

In general, total current can flow by drift and diffusion separately. **Total current density:**

$$\mathbf{J}_n = \mathbf{J}_n^{\text{drift}} + \mathbf{J}_n^{\text{diff}} = qn\mu_n\mathbf{E} + qD_n \frac{dn}{dx}$$

$$\mathbf{J}_p = \mathbf{J}_p^{\text{drift}} + \mathbf{J}_p^{\text{diff}} = qp\mu_p\mathbf{E} - qD_p \frac{dp}{dx}$$

$$\mathbf{J}_{\text{total}} = \mathbf{J}_n + \mathbf{J}_p$$

# What did we learn today?

## Summary of Key Concepts

- Electrons and holes in semiconductors are mobile and charged

–  $\Rightarrow$  **Carriers of electrical current!**

- **Drift current**: produced by electric field

$$\mathbf{J}^{\text{drift}} \propto \mathbf{E} \quad \mathbf{J}^{\text{drift}} \propto \frac{d\phi}{dx}$$

- **Diffusion current**: produced by concentration gradient

$$\mathbf{J}^{\text{diffusion}} \propto \frac{dn}{dx}, \frac{dp}{dx}$$

- Diffusion and drift currents are sizeable in modern devices
- Carriers move fast in response to fields and gradients

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