## 1 Molecular Diffusion

- Notion of concentration
- Molecular diffusion, Fick's Law
- Mass balance
- Transport analogies; salt-gradient solar ponds
- Simple solutions
- Random walk analogy to diffusion
- Examples of sources and sinks

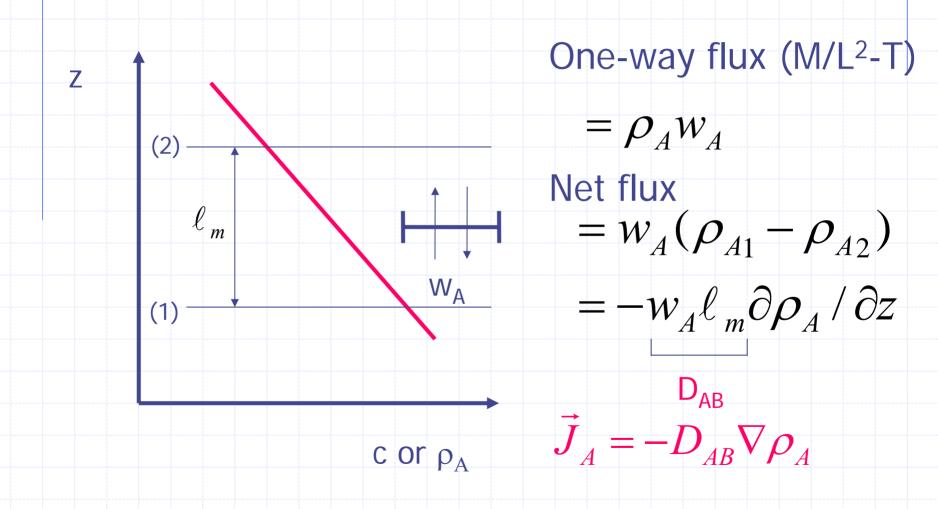
## Motivation

- Molecular diffusion is often negligible in environmental problems
- Exceptions: near interfaces, boundaries
- Responsible for removing gradients at smallest scales
- Analytical framework for turbulent and dispersive transport

#### Concentration

- Contaminant => mixture
  - Carrier fluid (B) and contaminant/tracer (A)
  - If dissolved, then solvent and solute
  - If suspended, then continuous and dispersed phase
- Concentration (c or  $\rho_A$ ) commonly based on mass/volume (e.g. mg/l); also
  - mol/vol (chemical reactions)
  - Mass fraction (salinity):  $\rho_A/\rho$
- Note: 1 mg/l ~ 1 mg/kg (water) = 1ppm

## Molecular Diffusion



## Ficks Law and Diffusivities

$$\vec{J}_A = -D_{AB} \nabla \rho_A$$

$$\nabla() = \frac{\partial}{\partial x}() \vec{i} + \frac{\partial}{\partial y}() \vec{j} + \frac{\partial}{\partial z}() \vec{k}$$

D<sub>AB</sub> is isotropic and essentially uniform (temperature dependent), but depends on A, B

Table 1.1 summarizes some values of  $D_{AB}$ 

Roughly:  $D_{air} \sim 10^{-1} \text{ cm}^2/\text{s}$ ;  $D_{water} \sim 10^{-5} \text{ cm}^2/\text{s}$ 

## Diffusivities, cont'd

Diffusivities often expressed through Schmidt no. Sc = v/D

Roughly:  $v_{air} \sim 10^{-1} \text{ cm}^2/\text{s}$ ;  $v_{water} \sim 10^{-2} \text{ cm}^2/\text{s}$ 

Sc<sub>air</sub> ~ 1

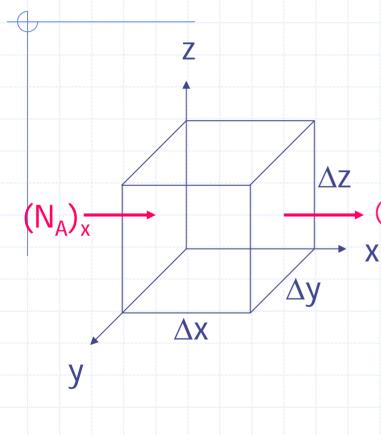
 $Sc_{water} \sim 10^3$ 

Also: Prandlt no. Pr =  $v/\kappa$  ( $\kappa$  = thermal cond.)

Add advection; total flux of A is:

$$\vec{N}_A = \rho_A \vec{q}_- - D_{AB} \nabla \rho_A$$
 macroscopic velocity vector

## Conservation of Mass



Like a bank account except expressed as rates:

(rate of) change in account =  $(N_A)_{x+\Delta x}$  (rate of) (inflow – out) +/- (rate of) prod/consumption

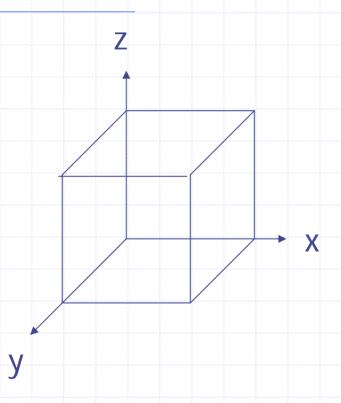
Example for x-direction

$$in = (N_A)_x \Delta y \Delta z$$

$$out = (N_A)_{x+\Delta x} \Delta y \Delta z = \left[ (N_A)_x + \left( \frac{\partial N_A}{\partial x} \right) \Delta x \right] \Delta y \Delta z$$

$$net \ in = \left[ -\left(\frac{\partial N_A}{\partial x}\right) \Delta x \right] \Delta y \Delta z$$

## Conservation of Mass, cont'd



Account balance

$$= \rho_A \Delta x \Delta y \Delta z$$

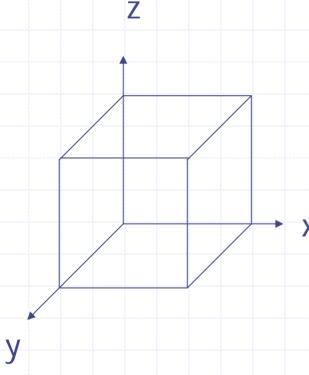
Rate of change of account balance

$$= \frac{\partial}{\partial t} (\rho_A) \Delta x \Delta y \Delta z$$

Rate of production

$$= r_A \Delta x \Delta y \Delta z$$

## Conservation of Mass, cont'd



Sum all terms (incl. advection in 3D)

$$\frac{\partial \rho_{A}}{\partial t} + \frac{\partial}{\partial x} (N_{A})_{x} + \frac{\partial}{\partial y} (N_{A})_{y} + \frac{\partial}{\partial z} (N_{A})_{z}$$

$$\frac{\partial \rho_{A}}{\partial t} + \nabla \cdot \vec{N}_{A} = r_{A}$$

Flux divergence (dot product of two vectors is scalar)

For carrier fluid B

$$\frac{\partial \rho_B}{\partial t} + \nabla \cdot \vec{N}_B = r_B$$

## Conservation of Mass, mixture

$$r_A = -r_B$$

$$\vec{N}_A + \vec{N}_B = \rho \vec{q}$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{q}) = 0$$

$$\frac{\partial \rho}{\partial t} \cong \nabla \rho \cong 0$$
$$\nabla \cdot \vec{q} = 0$$

$$\nabla \cdot \vec{q} = 0$$

Conservation of total mass

$$\rho_A + \rho_B = \rho$$

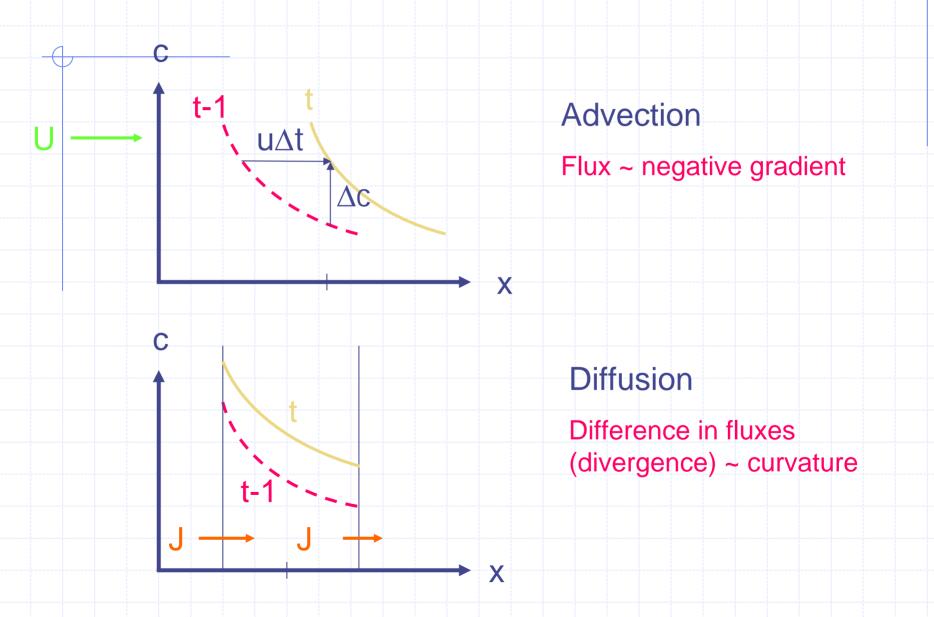
Liquids are nearly incompressible

Divergence = 0; Continuity

## Conservation of Mass, contaminant

$$\begin{split} \rho_A &= c & \text{drop subscript A} \\ \frac{\partial c}{\partial t} + \nabla \cdot (c\vec{q}) &= \nabla \cdot (D\nabla c) + r & \text{Conservative form of mass cons.} \\ \frac{\partial c}{\partial t} + c \nabla \cdot \vec{q} + \vec{q} \cdot \nabla c &= D\nabla^2 c + (\nabla D)(\nabla c) + r \\ 0 & 0 & 0 \\ \frac{\partial c}{\partial t} + \vec{q} \cdot \nabla c &= D\nabla^2 c \nabla + r & \text{N.C. form} \\ \frac{\partial c}{\partial t} &= D\nabla^2 c \nabla & \text{If} & \vec{q} = r = 0 => \text{Ficks Law of Diffusion} \end{split}$$

#### Heuristic interpretation of Advection and diffusion



## Analogs

$$\frac{\partial c}{\partial t} = D\nabla^2 c$$

$$\frac{\partial c}{\partial t} = D\nabla^2 c$$

$$\frac{\partial T}{\partial t} = \kappa \nabla^2 T$$

$$\frac{\partial \vec{q}}{\partial t} + (\vec{q} \cdot \nabla)\vec{q} = \upsilon \nabla^2 \vec{q} + \frac{\vec{r}_m}{\rho}$$
 Newton's Law (mom. Transfer)

v/DSc  $v/\kappa$  Pr

Ficks Law (mass transfer)

Fourier's Law (heat transfer)

Air Water

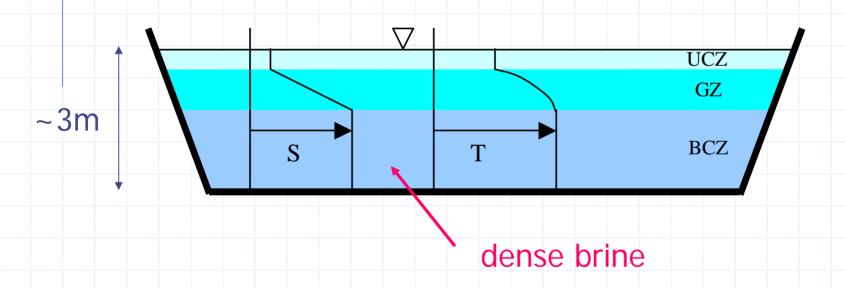
 $\sim 10^{3}$ 

~0.7

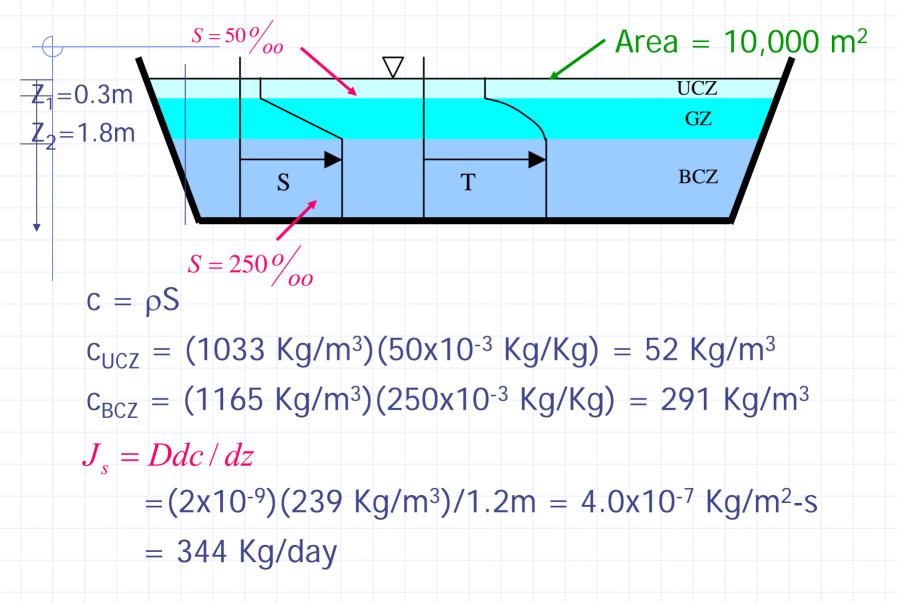
 $D < < \kappa < v$ D~K~V

## Example: Salt Gradient Solar Ponds (WE 1-1)

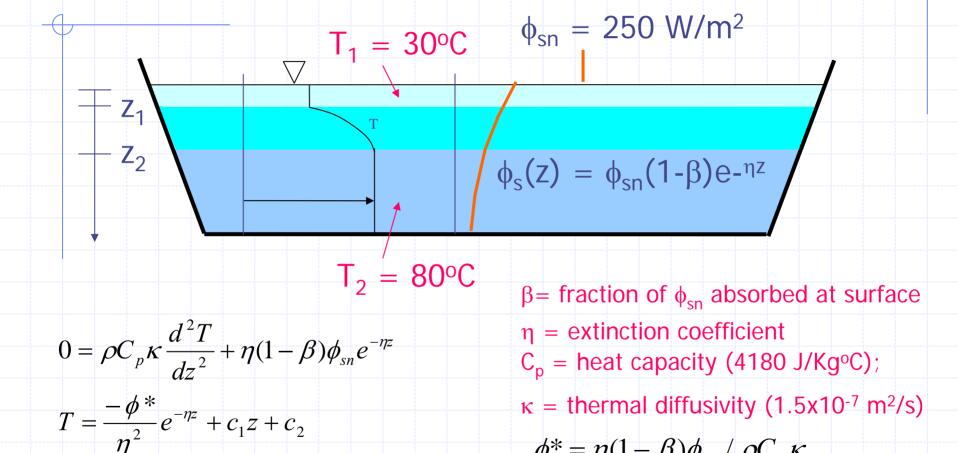
like El Paso Solar Pond



#### Solar Pond, diffusive salt flux to UCZ



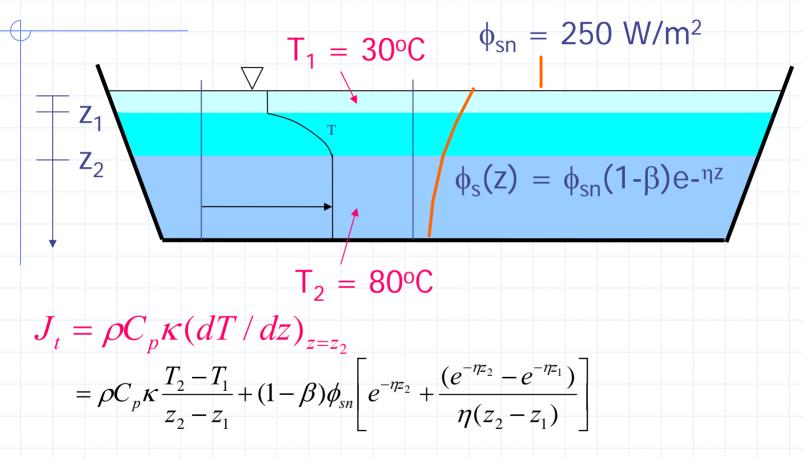
#### Solar Pond: diffusive thermal flux to UCZ



 $\phi^* = \eta(1-\beta)\phi_{sn}/\rho C_p \kappa$ 

 $c_1$ ,  $c_2$  from  $T=T_1$  at  $z_1$ ,  $T=T_2$  at  $z_2$ 

## Solar Pond: thermal flux



 $z_1 = 0.3m$ ;  $z_2 = 1.5m$ ,  $\beta = 0.5$ ,  $\eta = 0.6 = J_t = 7 \text{ W/m}^2$ 

Compare with  $(1-0.5)(250)\exp(-0.6*1.5) = 51 \text{ W/m}^2$  reaching BCZ (~13% lost)

## Rankine Cycle Heat Engine

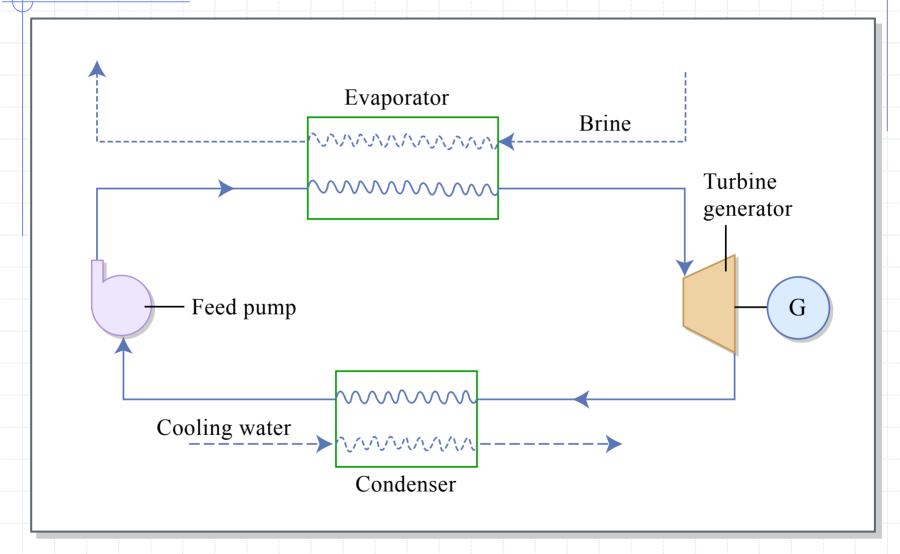
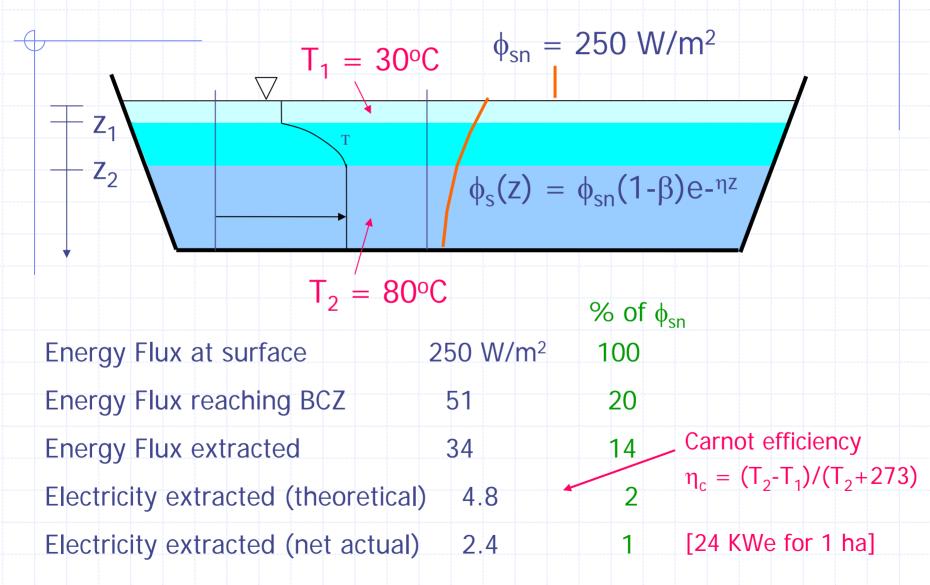


Figure by MIT OCW.

## Solar Pond: total energy extraction



## Simple Solutions

Inst. injection of mass M

$$\frac{\partial c}{\partial t} = D \frac{\partial^2 c}{\partial x^2}$$

$$bc: c = 0 \quad at \quad x = \pm \infty$$

$$ic: c = \frac{M}{A} \delta(x) \quad at \quad t = 0_+$$

$$r = \frac{M}{A} \delta(x) \delta(t) \quad with \quad c = 0 \quad at \quad t = 0$$

## Simple Solutions, cont'd

Inst. injection of mass M



$$c = \frac{B}{t^{1/2}} e^{-\frac{x^2}{4Dt}}$$

$$t$$

$$A \int_{-\infty}^{\infty} c dx = M$$

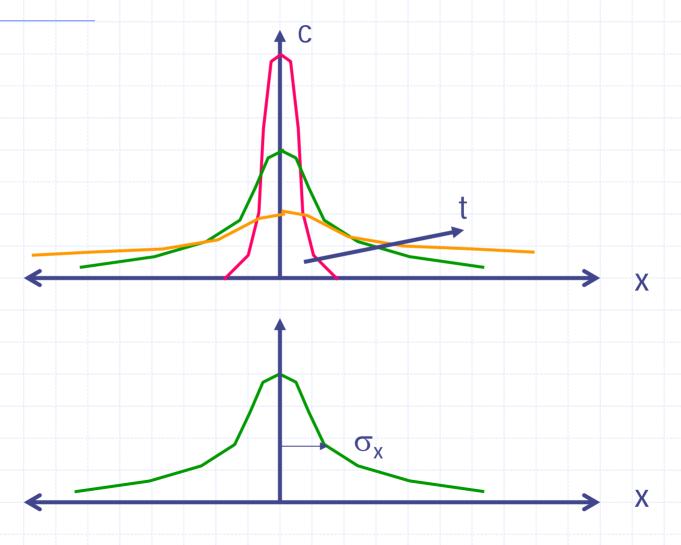
$$M = 2AB\sqrt{\pi D}$$

$$M = 2AB\sqrt{\pi D}$$

$$c(x,t) = \frac{M}{2A\sqrt{\pi Dt}}e^{-\frac{x^2}{4Dt}}$$

Add a current
$$c(x,t) = \frac{M}{2A\sqrt{\pi Dt}}e^{\frac{(x-ut)^2}{4Dt}}$$

## **Gaussian Solution**



## **Spatial Moments**

$$m_o = \int_0^\infty c(x,t)dx$$
  $M = m_o A$ 

$$m_1 = \int_{-\infty}^{\infty} cx dx$$

$$m_2 = \int_{-\infty}^{\infty} cx^2 dx$$

#### interpretation

$$M = m_o A$$

$$x_c = \frac{m_1}{m_o} = \underline{ut}$$

$$\sigma_x^2 = \frac{m_2}{m_o} - \left(\frac{m_1}{m_2}\right)^2 = 2Dt$$
 Plume variance

## Spatial Moments, cont'd

Relationship of moments to equation parameters

$$m_{2} = \int_{-\infty}^{\infty} \frac{M}{2A\sqrt{\pi Dt}} e^{\frac{-x^{2}}{4Dt}} x^{2} dt$$

$$= 2\frac{M}{A} Dt$$

$$\sigma_{x}^{2} = \frac{m_{2}}{m_{o}} = \frac{2MDt/A}{M/A} = 2Dt$$

Without current, odd moments are 0

## Spatial Moments, cont'd

#### Rewrite in terms of $\sigma$

$$c(x,t) = \frac{M}{2A\sqrt{\pi Dt}}e^{-\frac{x^2}{4Dt}}$$

$$=\frac{M}{A\sqrt{2\pi}\sigma_x}e^{-\frac{x^2}{2\sigma^2}}$$

#### or in 3-D (isotropic)

$$c(x,t) = \frac{M}{2A\sqrt{\pi Dt}}e^{-\frac{x^2}{4Dt}} \qquad c(x,t) = \frac{M}{8(\pi Dt)^{3/2}}e^{-\frac{(x^2+y^2+z^2)}{4Dt}}$$

$$=\frac{M}{(2\pi)^{3/2}\sigma^3}e^{-\frac{(x^2+y^2+z^2)}{2\sigma^2}}$$

#### Plume dilutes by spreading:

In 1-D, c ~ 
$$t^{-1/2}$$
 ~  $\sigma_x^{-1}$ 

In 3-D, c ~ 
$$t^{-3/2}$$
 ~  $\sigma^{-3}$ 

## Moment generating equation

$$\frac{\partial c}{\partial t} = D \frac{\partial^2 c}{\partial x^2}$$

$$m_i = \int_{-\infty}^{\infty} x^i c dx$$

Approach 1: moments of  $c(x,t) = > \sigma^2 = m^2/mo = 2Dt$ 

Approach 2: moments of ge => moment generation eq.

$$\int_{-\infty}^{\infty} x^{i}(each\ term)\ dx$$

## Moment generating eq., cont'd

$$\frac{\partial c}{\partial t} = D \frac{\partial^2 c}{\partial x^2} \qquad m_i = \int_{-\infty}^{\infty} x^i c dx$$

$$0^{th} \qquad \int_{-\infty}^{\infty} \frac{\partial c}{\partial t} dx = \frac{\partial}{\partial t} \int_{-\infty}^{\infty} c dx = \frac{dm_o}{dt}$$

$$moment \qquad \int_{-\infty}^{\infty} D \frac{\partial^2 c}{\partial x^2} dx = D \left( \frac{\partial c}{\partial x} \right)_{-\infty}^{\infty} = 0$$

$$2nd \qquad \int_{-\infty}^{\infty} x^2 \frac{\partial c}{\partial t} dx = \frac{\partial}{\partial t} \int_{-\infty}^{\infty} x^2 c dx = \frac{dm_2}{dt}$$

$$moment \qquad \int_{-\infty}^{\infty} Dx^2 \frac{\partial^2 c}{\partial x^2} dx = Dx^2 \left( \frac{\partial c}{\partial x} \right)_{-\infty}^{\infty} - 2 \int_{-\infty}^{\infty} x D \frac{\partial c}{\partial x} dx = -2xD(c)_{-\infty}^{\infty} + 2 \int_{-\infty}^{\infty} Dc dx = 2Dm_o$$

## Moment generating eq., cont'd

$$\frac{\partial c}{\partial t} = D \frac{\partial^2 c}{\partial x^2}$$

$$m_i = \int_{-\infty}^{\infty} x^i c dx$$

O<sup>th</sup> moment

$$\frac{dm_o}{dt} = 0$$

$$=> m_o = const = M/A$$

2nd

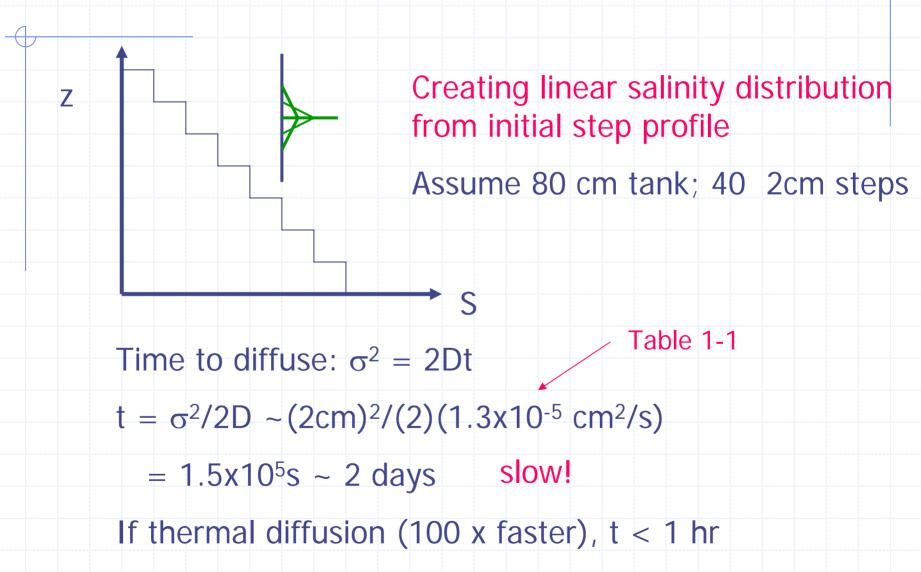
$$\frac{dm_2}{dt} = 2Dm_o$$

$$=> d\sigma^2/dt = 2D \text{ or } \sigma^2 = 2Dt$$

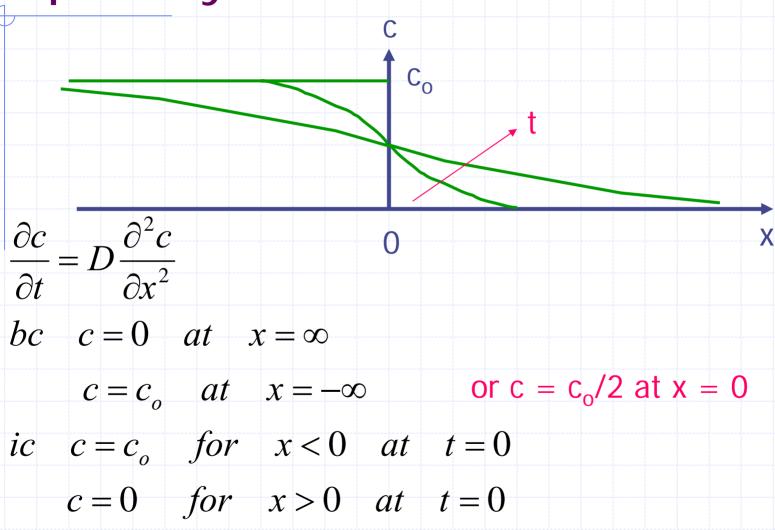
moment

$$m_o \frac{d\sigma^2}{dt} = 2Dm_o$$

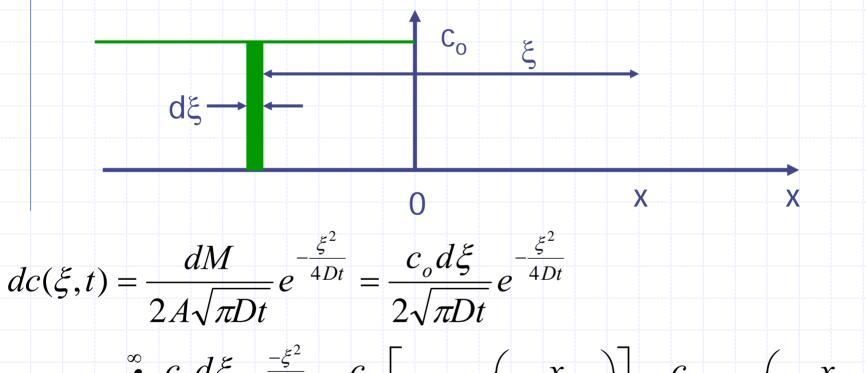
## How fast is molecular diffusion?



## Spatially distributed sources



## Spatially distributed sources



$$c(x,t) = \int_{x}^{\infty} \frac{c_o d\xi}{2\sqrt{\pi Dt}} e^{\frac{-\xi^2}{4Dt}} = \frac{c_o}{2} \left[ 1 - erf\left(\frac{x}{2\sqrt{Dt}}\right) \right] = \frac{c_o}{2} erfc\left(\frac{x}{2\sqrt{Dt}}\right)$$

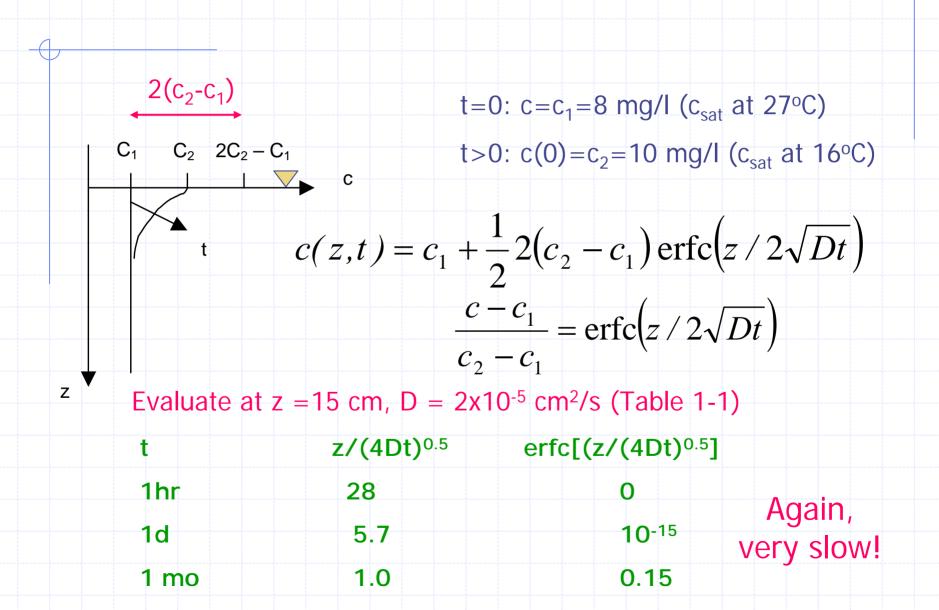
# Error Function $\frac{2}{\sqrt{\pi}}e^{-\alpha^{2}}$ erf( $\omega$ )

$$erf(\omega) = \frac{2}{\sqrt{\pi}} \int_{0}^{\omega} e^{-\alpha^{2}} d\alpha$$

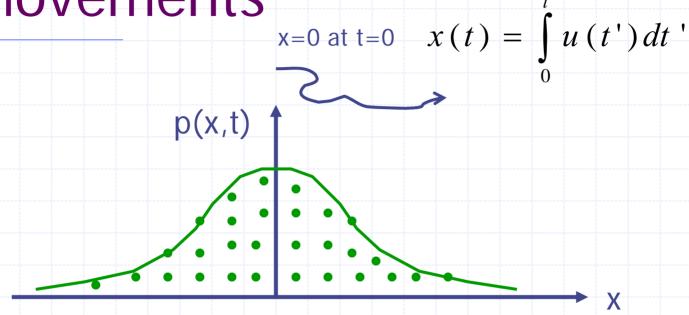
$$erfc(\omega) = \frac{2}{\sqrt{\pi}} \int_{\omega}^{\infty} e^{-\alpha^2} d\alpha$$

$$erf(0) = 0$$
  
 $erf(\infty) = 1$   
 $erfc(x) = 1 - erf(x)$ 

## Example: DO in Fish aquarium (WE 1-4)



# Diffusion as correlated movements



Analogy between p(x) and c(x); ergotic assumption

For many particles, both distributions become Normal (Gaussian) through Central Limit Theorem

## Statistics of velocity

$$\overline{u} = 0$$
 mean velocity
$$\overline{u^2} = const.$$
 variance
$$\overline{u(t)u(t-\tau)} = \overline{u(0)u(\tau)} = \overline{u(-\tau)u(0)}$$
 auto co-variance
$$\overline{u(t)u(t-\tau)} = R(\tau)$$
 auto correlation
$$1 + \overline{u(t)^2}$$

$$R(\tau)$$

## Statistics of position

$$x(t) = \int_{0}^{\infty} u(t')dt'$$

$$\overline{x} = \int_{0}^{t} u(t')dt' = \int_{0}^{t} \overline{u}(t)dt = 0$$

 $\overline{x^2(t)}$  increases with time, as follows

$$\frac{dx^2(t)}{dt} = 2x(t)\frac{dx}{dt} = 2\left[\int_0^t u(t')dt'\right]u(t) = 2\int_0^t u(t)u(t')dt'$$

$$\frac{dx^{2}(t)}{dt} = 2\overline{u^{2}(t)} \int_{0}^{t} R(t-t')dt' = 2\overline{u^{2}(t)} \int_{0}^{t} R(\tau)d\tau$$

$$D = \frac{dx^2}{2dt} = \frac{d\sigma^2}{2dt} = \overline{u^2} \int_0^t R(\tau) d\tau$$
 Taylor's Theorem (1921); classic

$$[D] = [V^2T]$$

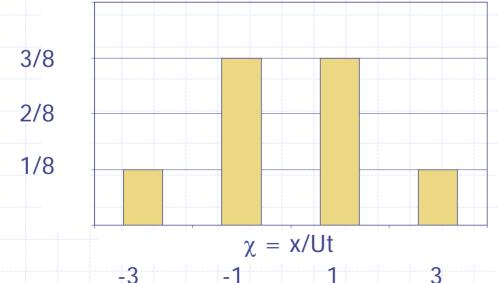
Earlier, 
$$D = w_A \ell_m$$
 [VL] or  $D = \frac{\sigma^2}{2t}$  [L<sup>2</sup>/T]

## Random Walk (WE 1-3)

Special case:  $u(t) = U \ or - U$  direction changes randomly after  $\Delta t$ 

 $x = \sum_{i=1}^{N} u \Delta t$ Walker's position at time  $t=N\Delta t$ 

Probability distribution 
$$p(\chi, N) = \frac{N!}{\left(\frac{N+\chi}{2}\right)! \left(\frac{N-\chi}{2}\right)!} \left(\frac{1}{2^N}\right)$$



Bernoulli Distribution

Approaches Gaussian for large N

Example for N=3

## Statistics of position

$$\overline{x}(t) = \sum_{i=1}^{N} \overline{u}(t) = 0$$

$$\sigma^{2} = \left(\sum_{i=1}^{N} u \Delta t\right)^{2} = \Delta t^{2} \left[ (u_{1} + u_{2} \dots u_{i} + \dots u_{N})(u_{1} + u_{2} + \dots u_{N}) \right]$$
$$= N \Delta t^{2} U^{2} = t U^{2} \Delta t = t \Delta x^{2} / \Delta t$$

$$D = \frac{\sigma^2}{2t} = \frac{U^2 \Delta t}{2} = \frac{\Delta x^2}{2\Delta t} \quad \text{[L^2/T]}$$

Alternatively, derive D from Taylor's Theorem

# Examples of Sources and Sinks (r terms)

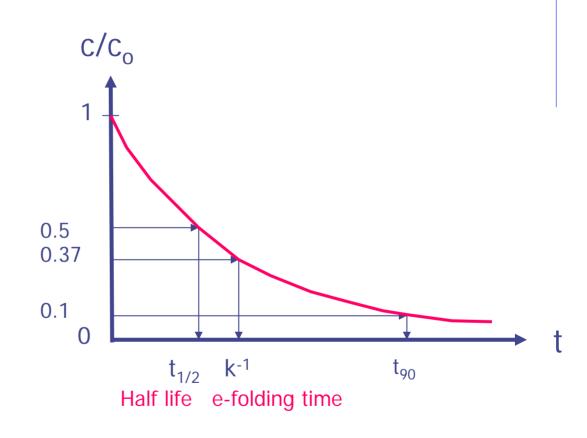
- ◆1<sup>st</sup> order
- ◆0<sup>th</sup> order
- ◆2<sup>nd</sup> order
- Coupled reactions
- Mixed order

#### 1st Order

Example: radioactive decay

$$\frac{dc}{dt} = -kc$$

$$c/c_o = e^{-kt}$$



Linearity => 1st O decay multiplies simple sol'n by e-kt; e.g.

Also very convenient in particle tracking models

$$c = \frac{M}{2A\sqrt{\pi Dt}}e^{-kt}$$

## Oth Order

Example: silica uptake by diatoms (high diatom conc)

$$\frac{dS}{dt} = -B$$
$$S = S_o - Bt$$



S=substrate (silica) concentration

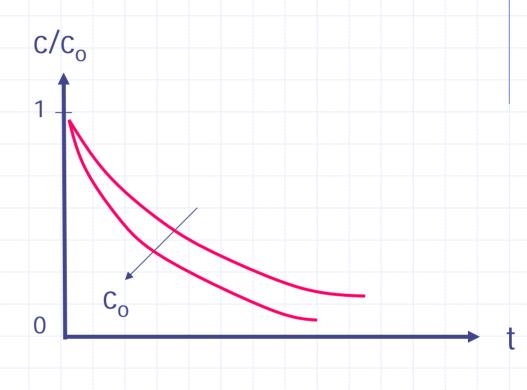
B=rate (depends on diatom population, but assume large)

## 2<sup>nd</sup> Order

Example: particleparticle collisions/reactions; flocculant settling)

$$\frac{dc}{dt} = -Bc^{2}$$

$$\frac{c}{c_{o}} = \frac{1}{1 + Btc_{o}}$$



Behavior depends on c<sub>o</sub>; slower than e<sup>-kt</sup>.

Can be confused with multiple species undergoing 1st order removal

## **Coupled Reactions**

**Example: Nitrogen oxidation** 

$$\frac{dN_1}{dt} = -K_{12}N_1 \qquad N_1 = NH_3-N$$

$$\frac{dN_2}{dt} = K_{12}N_1 - K_{23}N_2 \qquad N_2 = NO_2-N$$

$$\frac{dN_3}{dt} = K_{23}N_2 \qquad N_3 = NO_3-N$$

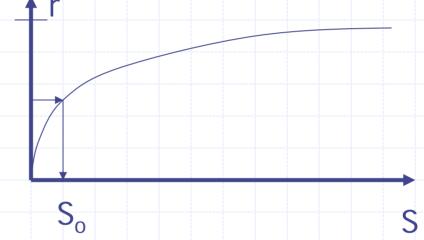
If N's are measured as molar quantities, or atomic mass, then successive K's are equal and opposite

## Mixed Order—Saturation Kinetics (Menod kinetics)

Example: algal uptake of nutrients—focus on algae

$$\frac{dc}{dt} = \frac{kS}{S + S_o}$$

 $r_{max}$   $\uparrow$  r  $R_{max}/2$ 



$$S << S_o =>$$
  $\frac{dc}{dt} \cong \frac{kS}{S_o} = k'S$  (1st Order)  
 $S >> S_o =>$   $\frac{dc}{dt} \cong k$  (0th Order)

$$S >> S_0 => \frac{dc}{dt} \cong k$$

(0th Order)

## 1 Wrap-up

Molecular diffusivities

$$D = w\ell_m$$

Molecular motion; Eulerian frame

$$D = w\ell_m$$

$$D = \frac{d\sigma^2}{2dt^2}$$

Method of moments

$$D = \overline{u^2} \int_{0}^{\infty} R(\tau) d\tau$$

 $D = \overline{u^2} \int_{0}^{\infty} R(\tau) d\tau$  Molecular motion; Lagrangian frame

- ◆ D is "small" ~ 1x10<sup>-5</sup> cm<sup>2</sup>/s for water
- Inst. point source solutions are Gaussian; other solutions built from
  - Spatial and temporal integration, coordinate translation, linear source/sink terms