

Supplemental Notes for Chapter 17
General Treatment of Phase and Chemical Equations

1. Phase rule constrained parameter variability

- $F = n + 2 - \pi - r$
- invariant ($F = 0$)
- monovariant ($F = 1$)
- divariant ($F = 2$)
- other constraints (stoichiometric ratio, “indifferent states”)

2. Matrix/determinant formalism applied to generalized Gibbs-Duhem and reaction equilibrium expressions

- phase and chemical equilibrium criteria combined
- Gibbs-Duhem combined with reaction equilibrium criteria to produce a generalized expression for π phases, r reactions and n components

$$\sum_{s=1}^{\pi} \left[-\underline{S}^{(s)} dT + \underline{V}^{(s)} dP - \sum_{i=1}^n N_i^{(s)} d\mu_i \right] \quad \text{and} \quad \sum_{i=1}^n \sum_{k=1}^r v_i^{(k)} d\mu_i$$

$$-\begin{bmatrix} \underline{S}^{(1)} \\ \underline{S}^{(\pi)} \\ 0^{(1)} \\ 0^{(r)} \end{bmatrix} \begin{bmatrix} dT \\ dP \end{bmatrix} + \begin{bmatrix} \underline{V}^{(1)} \\ \vdots \\ \underline{V}^{(\pi)} \\ 0^{(1)} \\ \vdots \\ 0^{(r)} \end{bmatrix} - \begin{bmatrix} N_1^{(1)} & \dots & N_j^{(1)} & \dots & N_n^{(1)} \\ \vdots & & \vdots & & \vdots \\ N_1^{(\pi)} & \dots & N_j^{(\pi)} & \dots & N_n^{(\pi)} \\ v_1^{(1)} & \dots & v_j^{(1)} & \dots & v_n^{(1)} \\ \vdots & & \vdots & & \vdots \\ v_1^{(r)} & \dots & v_j^{(r)} & \dots & v_n^{(r)} \end{bmatrix} \begin{bmatrix} d\mu_1 \\ d\mu_j \\ d\mu_n \end{bmatrix} = 0$$

or in shorthand vector notation

$$-\underline{S}^{(s)} dT + \underline{V}^{(s)} dP - \underline{\underline{N}}_j^{(s)} d\mu_j = 0$$

3. Invariant systems ($F=0$)

- $d\mu_i = 0$
- mass and mole constraints
- $|\tilde{N}^{(s)}| \neq 0$ required or the rank of the $N^{(s)}$ matrix must be reduced to avoid an over-constrained system
- indifferent states result if $|\tilde{N}^{(s)}| = 0$

4. Monovariant systems ($F=1$) pressure-temperature variations

- generalized Clapeyron equation
- use Gibbs-Duhem with chemical reactions included
- use a chemical potential or fugacity approach to get

$$\left. \frac{dP}{dT} \right|_{F=1} = \frac{|\Delta H|}{T |\Delta V|} \quad \text{or equivalently} \quad \frac{|\Delta S|}{|\Delta V|}$$

$$|\Delta H| \equiv \begin{vmatrix} H^{(1)} & x_1^{(1)} & x_2^{(1)} & \dots & x_n^{(1)} \\ \vdots & \vdots & & & \vdots \\ H^{(\pi)} & x_1^{(\pi)} & x_2^{(\pi)} & \dots & x_n^{(\pi)} \\ 0 & v_1^{(1)} & v_2^{(1)} & \dots & v_n^{(1)} \\ \vdots & \vdots & & & \vdots \\ 0 & v_1^{(r)} & v_2^{(r)} & \dots & v_n^{(r)} \end{vmatrix} \approx |\Delta V| \equiv \begin{vmatrix} V^{(1)} & x_1^{(1)} & x_2^{(1)} & \dots & x_n^{(1)} \\ \vdots & \vdots & & & \vdots \\ V^{(\pi)} & x_1^{(\pi)} & x_2^{(\pi)} & \dots & x_n^{(\pi)} \\ 0 & v_1^{(1)} & v_2^{(1)} & \dots & v_n^{(1)} \\ \vdots & \vdots & & & \vdots \\ 0 & v_1^{(r)} & v_2^{(r)} & \dots & v_n^{(r)} \end{vmatrix}$$

5. Isobaric monovariant equilibria, temperature-composition variations

-apply $dP = 0$ constraint to the set of generalized Gibbs-Duhem equations and
use the Gibbs-Helmholtz relationship to simplify

$$\left[\frac{\partial T}{\partial x_1^\beta} \right]_{P,[\alpha,\dots,\pi]} = \frac{-T \left| \tilde{x}_i^{(s)} \right| \left(\partial \mu_1^\beta / \partial x_1^\beta \right)_{T,P}}{\left[|\Delta H| - \overline{H}_1^\beta \left| \tilde{x}_i^{(s)} \right| \right]}$$

or with the fugacity introduced

$$\left[\frac{\partial T}{\partial x_1^\beta} \right]_{P,[\alpha,\dots,\pi]} = \frac{-RT^2 \left| \tilde{x}_1^{(s)} \right| \left[\frac{\partial \ln \hat{f}_1^\beta}{\partial x_1^\beta} \right]_{T,P}}{\left[|\Delta \tilde{H}| - \bar{H}_1^\beta \left| \tilde{x}_i^{(s)} \right| \right]}$$

6. Indifferent states and azeotropic behavior ($F \geq 2$)

$$\left| \tilde{N}_j^{(s)} \right| \text{ or } \left| \tilde{x}_j^{(s)} \right| = 0$$

- rank of the mole or composition matrix must be reduced which is equivalent to reducing the number of components
- Gibbs-Konovalows' first and second theorems

$$-\mathcal{S}^{(s)}dT + \mathcal{V}^{(s)}dP - \tilde{N}_j^{(s)}d\mu_j = 0$$

if $\left| \tilde{N}_j^{(s)} \right| = 0$ and if $dP = 0$ then $dT = 0$ (extremum in T)

or

if $\left| \tilde{N}_j^{(s)} \right| = 0$ and if $dT = 0$ then $dP = 0$ (extremum in P)

$$\tilde{N}_j^{(s)} = \begin{bmatrix} N_1^{(1)} & \dots & N_j^{(1)} & \dots & N_n^{(1)} \\ \vdots & & \vdots & & \vdots \\ N_1^{(\pi)} & \dots & N_j^{(\pi)} & \dots & N_n^{(\pi)} \\ v_1^{(1)} & \dots & v_j^{(1)} & \dots & v_n^{(1)} \\ \vdots & & \vdots & & \vdots \\ v_1^{(r)} & \dots & v_j^{(r)} & \dots & v_n^{(r)} \end{bmatrix}$$

and

$$\tilde{x}_j^{(s)} = \begin{bmatrix} x_1^{(1)} \dots x_j^{(1)} \dots x_n^{(1)} \\ \vdots & \vdots & \vdots \\ x_1^{(\pi)} & x_j^{(\pi)} \dots x_n^{(\pi)} \\ v_1^{(1)} & v_j^{(1)} \dots v_n^{(1)} \\ \vdots & \vdots & \vdots \\ v_1^{(r)} \dots v_j^{(r)} \dots v_n^{(\pi)} \end{bmatrix}$$