

MASSACHUSETTS INSTITUTE OF TECHNOLOGY

Molecular, Cellular and Tissue Biomechanics  
BEH.410 / 2.978J / 6.524J / 10.537J

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**Problem Set #1**

**Issued: 2/13/03**

**Due: 2/19/03 (in class)**

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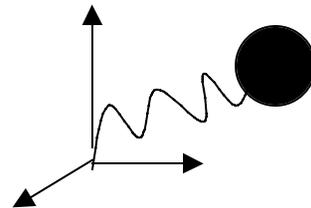
Please staple each problem separately as we will have different graders grading each problem. Also please put your name on each problem!

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**Problem #1: Boltzmann's Relation and Cell Micromechanics**

Recently D. Discher studied the mechanical behavior of the red blood cell cytoskeleton (we will learn soon that this is a polymer network attached to the cell plasma membrane). Discher attached a bead of 40 nm to this network (see figure A below) and tracked the center of mass of the bead over time (figure B). The network acts as a *spring* constraining the motion of the bead.

Image removed due to copyright considerations.



a) Consider a simple system consisting of a bead attached to a spring (with spring constant  $k$ ) at constant  $N, V, T$ . Calculate the mean squared displacement relative to the average position  $\langle (x - \langle x \rangle)^2 \rangle$ . (Note: you should start by first proving the general statistical relationship:  $\langle (x - \langle x \rangle)^2 \rangle = \langle x^2 \rangle - \langle x \rangle^2$ )

b) From the data above, estimate the cytoskeleton's spring constant.

c) Your analysis in part a) can be generalized to prove the Principle of Equipartition of Energy: the mean value of each independent quadratic term in the energy of a system is equal to  $1/2 kT$ . Prove that in general if the energy of a molecule depends on the square of a parameter (such as position), then the mean energy  $\langle U \rangle$  associated with the parameter is equal to  $1/2 kT$ .

## Problem #2: Microrheology and the Langevin Equation

In problem #1 we used equilibrium Statistical Mechanics to relate the mean squared displacement of the bead to the spring constant of the cytoskeleton. Here we will employ a simplified 1-dimensional form of the Langevin equation:

$$\zeta \frac{dx}{dt} = -\kappa x + f(t)$$

$$\langle f(t) \rangle = 0$$

$$\langle f(t)f(t') \rangle = 2\zeta kT \delta(t - t')$$

You will notice that we have neglected inertia and the Brownian force is delta correlated. (Note: the prefactor of  $2kT$  is obtained using the inertial form of the Langevin equation without spring forces to solve for the velocity correlation function

$$\langle v(t + \tau)v(t) \rangle = \frac{F}{2m\zeta} \exp(-\tau\zeta / m) \text{ then using equipartition of energy at } \tau = 0.)$$

a) Show that the position autocorrelation function for the 1-D model above is

$$\langle x(t + \tau)x(t) \rangle = \frac{kT}{\kappa} \exp(-\tau\kappa / \zeta). \text{ Note: you can start with the fact that}$$

$$x(t) = \int_{-\infty}^t \exp\left(\frac{-\kappa(t-t')}{\zeta}\right) dt'.$$

b) In a particle tracking experiment one usually measures  $\langle (x(t + \tau) - x(t))^2 \rangle$ .

Use your result from part a) and equipartition of energy to arrive at a simple expression for  $\langle (x(t + \tau) - x(t))^2 \rangle$ .

c) Calculate the limits of  $\langle (x(t + \tau) - x(t))^2 \rangle$  at short  $\tau$  and long  $\tau$ .

d) Sketch your answer from part c) and describe how you would calculate the viscosity and spring constant of a complex medium (such as a cell) from a plot of  $\langle (x(t + \tau) - x(t))^2 \rangle$  versus  $\tau$ .

## Problem #3: Collapse of a Macromolecule

Problem #11 in chapter 10 of Dill and Bromberg