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16.346 Astrodynamics
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Locus of Mean Points

#6.4

Definition of **mean point**: At the point \mathbf{r}_0 , the velocity \mathbf{v}_0 is parallel to the chord. Using the eccentricity vector at this point, we have $\mathbf{v}_0 \times \mathbf{h} \cdot (\mathbf{r}_2 - \mathbf{r}_1) = 0$ Therefore:

$$\mu \mathbf{e} \cdot (\mathbf{r}_2 - \mathbf{r}_1) = \left(\mathbf{v}_0 \times \mathbf{h} - \frac{\mu}{r_0} \mathbf{r}_0 \right) \cdot (\mathbf{r}_2 - \mathbf{r}_1) = 0 - \frac{\mu}{r_0} \mathbf{r}_0 \cdot (\mathbf{r}_2 - \mathbf{r}_1)$$

Hence:
$$\underbrace{\mathbf{e} \cdot (\mathbf{r}_2 - \mathbf{r}_1)}_{= r_1 - r_2} = -\frac{1}{r_0} \mathbf{r}_0 \cdot (\mathbf{r}_2 - \mathbf{r}_1) \implies \boxed{\mathbf{r}_0 \cdot (\mathbf{r}_2 - \mathbf{r}_1) = r_0(r_2 - r_1)}$$

The loci of all mean points are the lines through the focus F and the extremities of the minor axis of the fundamental ellipse.

The Line Segment FS:

Page 266

The line FS is the distance along mean point locus from the focus F to intersection with chord. The flight direction angle at P_0 is γ_0 and δ is the angle opposite the line segment SP_1 . Use the law of sines for the triangles:

$$\triangle FP_1S : \frac{FS}{\sin(\gamma_0 + \delta)} = \frac{r_1}{\sin \gamma_0} \quad \triangle FP_1P_2 : \frac{c}{\sin \theta} = \frac{r_2}{\sin(\gamma_0 + \delta)}$$

and use the calculation on the previous page for $1 - e_F^2$:

$$\triangle FP_0C : \sin \gamma_0 = \frac{b_F}{a_F} = \sqrt{1 - e_F^2} = \frac{2}{c} \sqrt{r_1 r_2} \sin \frac{1}{2} \theta$$

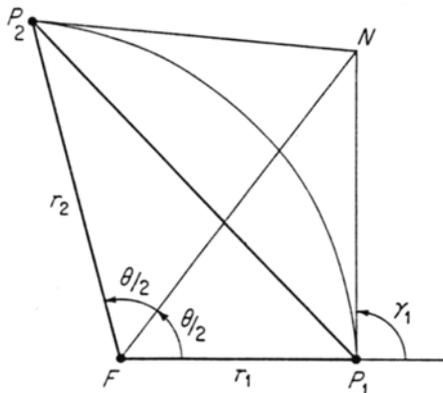
Therefore:

$$\boxed{FS = \sqrt{r_1 r_2} \cos \frac{1}{2} \theta}$$

Recall the proposition:

Lecture 8, Page 2

The line connecting the focus and the point of intersection of the orbital tangents at the terminals bisects the transfer angle.



$$\sqrt{r_1 r_2} = \begin{cases} FN \cos \frac{1}{2} (E_2 - E_1) & \text{ellipse} \\ FN & \text{parabola} \\ FN \cosh \frac{1}{2} (H_2 - H_1) & \text{hyperbola} \end{cases}$$

Fig. 6.7 from *An Introduction to the Mathematics and Methods of Astrodynamics*. Courtesy of AIAA. Used with permission.

For the ellipse $FN_1 \cos \frac{1}{2}(E_0 - E_1) = \sqrt{r_1 r_0}$ $FN_2 \cos \frac{1}{2}(E_2 - E_0) = \sqrt{r_0 r_2}$

and for the parabola $FN_{1p} = \sqrt{r_1 r_{0p}}$ $FN_{2p} = \sqrt{r_{0p} r_2}$

Triangle $\Delta FN_{1p}N_{2p}$ is similar to triangle ΔFN_1N_2 . Therefore

$$\frac{FN_{1p}}{FN_{2p}} = \frac{FN_1}{FN_2} \implies \cos \frac{1}{2}(E_0 - E_1) = \cos \frac{1}{2}(E_2 - E_0)$$

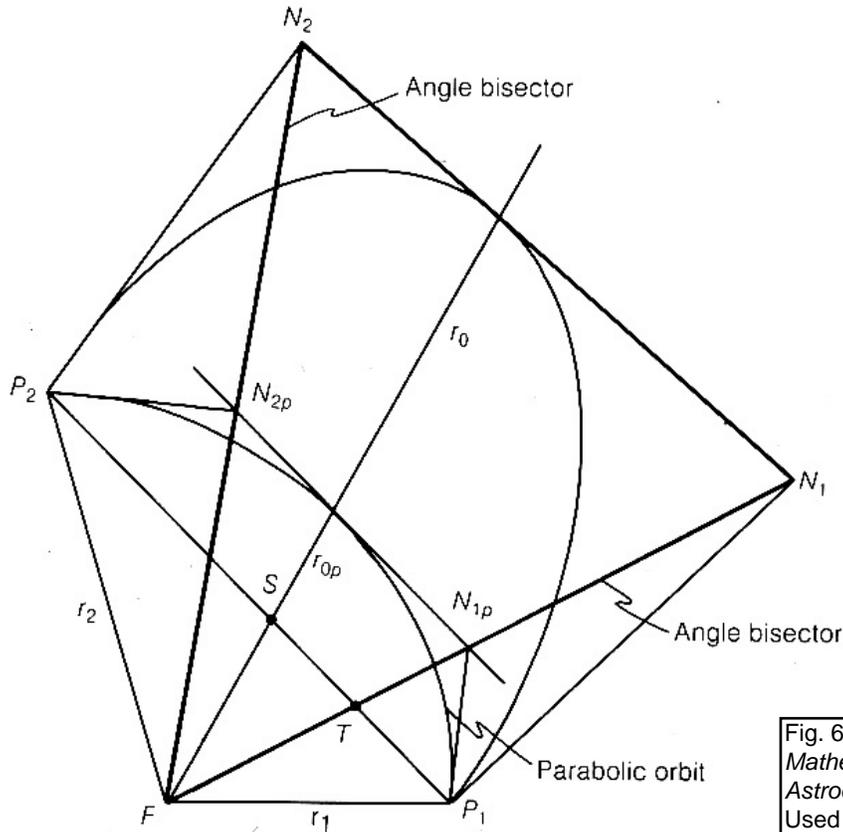


Fig. 6.16 from *An Introduction to the Mathematics and Methods of Astrodynamics*. Courtesy of AIAA. Used with permission.

Hence

$$E_0 = \frac{1}{2}(E_1 + E_2)$$

6.73

The eccentric anomaly of the mean point of an orbit connecting two termini is the arithmetic mean between the eccentric anomalies of those termini.

The parameter of the parabola is obtained from

$$\left(\frac{p_p}{p_m}\right)^2 - 2D\frac{p_p}{p_m} + 1 = 0 \quad \text{where} \quad D = \frac{r_1 + r_2}{c} - \frac{s(s-c)}{ac} = \frac{r_1 + r_2}{c}$$

so that

$$\frac{p_p}{p_m} = \frac{1}{c}(r_1 + r_2 + 2\sqrt{r_1 r_2} \cos \frac{1}{2}\theta) = \frac{1}{c}(\sqrt{s} + \sqrt{s-c})^2$$

The mean point radius of the parabola is

$$r_{0p} = \frac{p_p}{1 + \cos 2\phi_F} = \frac{p_p}{2 \cos^2 \phi_F} = \frac{p_p}{2(1 - e_F^2)} = \frac{p_p c}{4p_m}$$

Hence

$$r_{0p} = \frac{1}{4}(r_1 + r_2 + 2\sqrt{r_1 r_2} \cos \frac{1}{2}\theta) = \frac{1}{2}(a_F + FS)$$

The mean point radius of the parabola extends to the midpoint between the chord and the extremity of the minor axis of the fundamental ellipse.

From the derivation of the eccentric anomaly of the mean point, we have

$$FN_2 \cos \frac{1}{2}(E_2 - E_0) = FN_2 \cos \frac{1}{2}[E_2 - \frac{1}{2}(E_1 + E_2)] = FN_2 \cos \frac{1}{4}(E_2 - E_1) = \sqrt{r_0 r_2}$$

$$FN_{2p} = \sqrt{r_{0p} r_2}$$

so that

$$\frac{FN_2}{FN_{2p}} \cos \frac{1}{4}(E_2 - E_1) = \sqrt{\frac{r_0}{r_{0p}}}$$

But, from similar triangles,

$$\frac{FN_2}{FN_{2p}} = \frac{r_0}{r_{0p}}$$

Therefore, we have the truly elegant expression

$$r_0 = r_{0p} \sec^2 \frac{1}{4}(E_2 - E_1) = r_{0p} \sec^2 \frac{1}{2}\psi = r_{0p}(1 + \tan^2 \frac{1}{2}\psi)$$

and, as we might expect,

$$r_0 = r_{0p} \operatorname{sech}^2 \frac{1}{4}(H_2 - H_1)$$

obtains also for hyperbolic orbits.

$$r_0 = \begin{cases} a[1 - e \cos \frac{1}{2}(E_1 + E_2)] = a(1 - \cos \phi) \\ a[1 - e \cosh \frac{1}{2}(H_1 + H_2)] \end{cases} = \begin{cases} r_{0p}(1 + \tan^2 \frac{1}{2}\psi) \\ r_{0p}(1 - \tanh^2 \frac{1}{4}(H_2 - H_1)) \end{cases}$$