

Linear elasticity:

Dirichlet form: $a(u, v) = \int_B C_{ijkl} u_{ij} v_{kl} dv$

Energy norm: $\|u\|_E = [a(u, u)]^{1/2}$

If $C_{ijkl}(x) \in L^\infty(B)$, coercive, convex \Rightarrow

$$C \|u\|_h \leq \|u\|_E \leq C' \|u\|_h$$

Again: $\|u - u_h\| \leq \inf_{v_h \in V_h} \|u - v_h\|_E$ (best approx.)
 (of FE solution)

$$\Rightarrow \|u - u_h\|_E \leq \|u - u_I\|_E \leq C \|u - u_I\|_h$$

$$\Rightarrow \|u - u_h\|_E \leq \sum_{e=1}^E C \sigma_h^e (h^e)^k |u^e|_{k+1}$$

($m=1$)

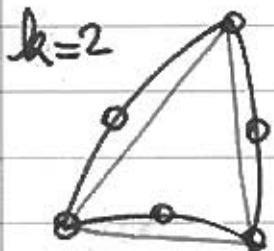
(order of derivative in Dirichlet form)

The previous estimate provides the rate of convergence as the mesh is refined ($h \rightarrow 0$). However a knowledge of the bound requires a knowledge of the unknown exact solution "u" (a priori error estimates). Want a posteriori

Need $|u^e|_{k+1}$, $e=1, \dots, E$

Cannot use $|\mu_h^e|_{k+1}$ since = 0
 $(\mu_h^e \in P_k(\Omega_h^e))$

- Assume $k \geq 2$



$$\mu_h \in R_k(\Omega_h^e)$$

$v_h \in P_{k-1}(\Omega_h^e)$: P_{k-1} interpolant to μ_h

v_h also defines a converging approximation

$$\|u - v_h\|_E \leq \|u - \mu_h + \mu_h - v_h\|_E \leq \|u - \mu_h\|_E + \|\mu_h - v_h\|_E$$

$$\leq \sum_{e=1}^E C \sigma_h^e (h^e)^k |\mu^e|_{k+1} + \sum_{e=1}^E C \sigma_h^e (h^e)^{k-1} |\mu_h^e|_k$$

$$\Rightarrow \bullet \|u - v_h\|_E \rightarrow 0 \text{ as } h^e \rightarrow 0$$

the first term converges faster, i.e.;

as $h^e \rightarrow 0$, $(h^e)^k$ cancels compared to $(h^e)^{k-1}$

$$\rightarrow \|u - v_h\|_E \leq \sum_{e=1}^E C \tau_h^e (h^e)^{k-1} \|u_h^e\|_k$$

"A posteriori" error estimate, "local" can be computed element by element.

Numerical integration errors

Fully integrated case: $a(u_h, v_h) = \langle f, v_h \rangle \quad \forall v_h \in V_h$

$$\text{where } a(u, v) = \int_B c_{ijkl} u_{i,j} v_{k,l} dV$$

When we introduce numerical quadrature we obtain:

$$\tilde{a}(u, v) = \sum_{e=1}^E \sum_{q=1}^Q w_q (c_{ijke} u_{i,j} v_{k,l}) (\xi_q^e)$$

$$\tilde{a}(\tilde{u}_h, v_h) = \langle f, v_h \rangle \quad \forall v_h \in V_h$$

In general $\tilde{u}_h \neq u_h$

We know that:

$$\|u - u_h\|_E \sim O(h^{k+1-m}) \text{ as } h \rightarrow 0$$

Under what conditions:

$$\|u - \tilde{u}_h\|_E \rightarrow 0 \text{ as } h \rightarrow 0?$$

At what rate?

Strong & Fix, p. 181

Proposition: Assume $\exists c > 0 / \tilde{a}(u_h, u_h) \geq c \|u_h\|_{V_h}$
 $\forall u_h \in V_h$
 (definiteness)

• Assume: (exact quadrature for special cases)

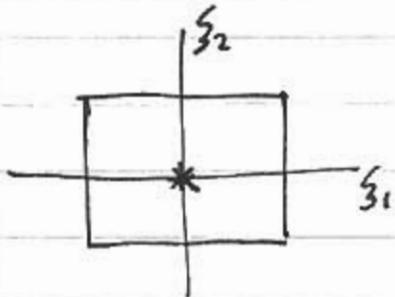
$$\sum_{q=1}^Q w_q [C_{ijkl} \epsilon_{kl} v_{ij}] (s_q) = \int_{\Omega} C_{ijkl} \epsilon_{kl} v_{ij} d\Omega$$

$$+\varepsilon_{kl} \in P_{l-m}(\Omega), +v_h \in V_h$$

$$\Rightarrow \|u - \tilde{u}_h\|_m \sim O(h^{l+1-m}) \text{ as } h \rightarrow 0$$

• Corollary: $\ell = k$ for same rate of convergence as fully integrated solution

Examples:

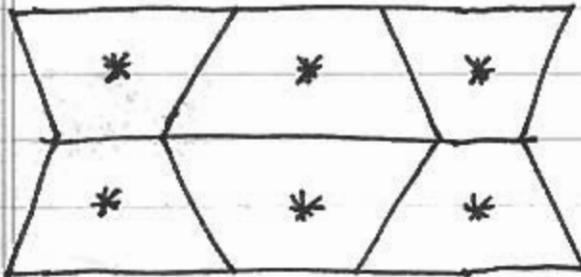


Numerical quadrature must be exact for:

$$\varepsilon_{ij} \in P_0(\Omega)$$

(constant strains)

It seems it would be enough to use "one" quadrature point. However:



(hourglass mode)

$$\varepsilon_{ij}(\xi_3) = 0 \Rightarrow$$

$$\tilde{a}(u_h, u_h) = 0$$

"spurious zero energy mode"

\tilde{a} is not positive definite

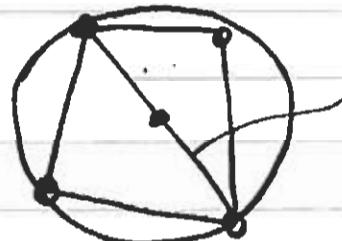
Check always that " K^e " has no zero energy

modes other than rigid body modes.

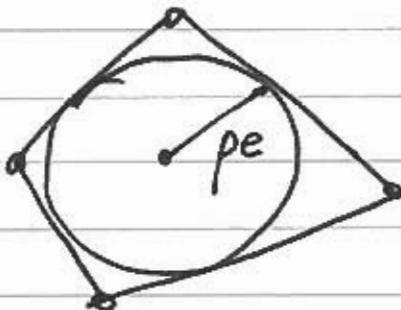
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Notation:

- "m" orders of derivatives in Dirichlet form. ($m=1$ for linear elasticity)
($m=2$ for plate theory)
- "k" highest order of polynomials
fully included in interpolation



h^e : radius of smallest circumscribed circle
↳ element size



ρ^e = biggest circle inscribed in the element.

Basic error estimates

$$\|u - u_h\| \leq C(u) \frac{h^{k+1}}{\rho^m}$$

$h = h^e$ } for " e " such that $\frac{(h^e)^{k+1}}{(\rho^e)^m}$ is maximum
 $\rho = \rho^e$ }

Define $\sigma = \frac{h^e}{\rho^e}$ element aspect ratio

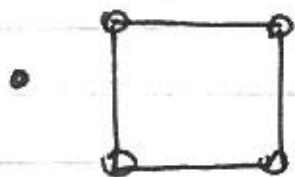
$$\|u - u_h\| \leq C(u) h^{k+1} \frac{(\sigma^e)^m}{(h^e)^m} \leq C(u) h^{k+1-m} (\sigma)^m$$

Where $\sigma_h \leq \sigma$, assume regular refinements

$$\|u - u_h\| \leq C'(u) h^{k+1-m}$$

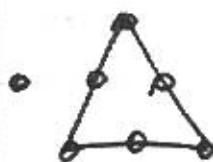
→ $k+1-m$ = rate of convergence

Examples

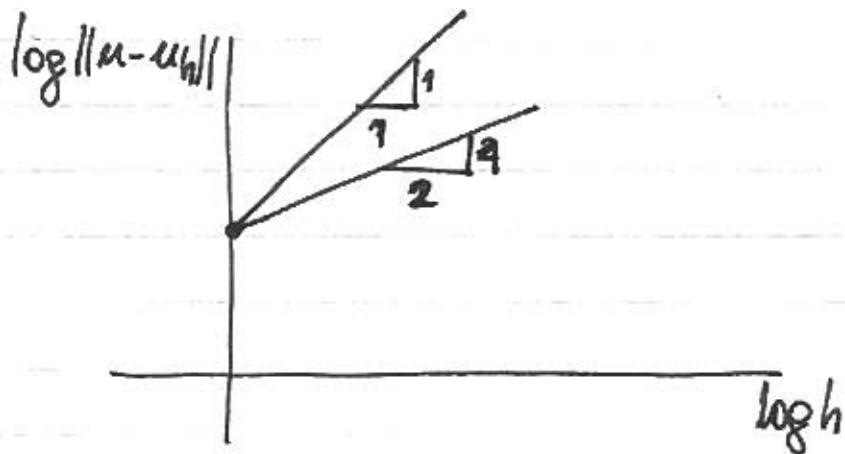


linear elasticity $k=1$
 $m=1$

$$\|u - u_h\| = C(u) h$$



$k=2$ $m=1$ $\|u - u_h\| = C(u) h^2$



Conditions for convergence

① $\|u_h\| < \infty$, otherwise $\|u - u_h\| \rightarrow \infty$

Finite element interpolation must give u_h with finite energy.

Sufficient conditions $N_a \in C^m(\Omega^e)$

$N_a \in C^{m-1}(\partial\Omega_{int}^e)$

Linear elasticity: $N_a \in C^1(\Omega^e)$
 $N_a \in C^0(\partial\Omega^e)$

Beam theory: $N_a \in C^2(\Omega^e)$
 $N_a \in C^1(\partial\Omega^e)$ interior

Elasticity: $m=1 \Rightarrow k > 0$, "k" at least 1

Patch Test

$$\textcircled{2} \quad \|u - u_h\| \leq C(u) h^{k+1-m}$$

For convergence

$$k+1-m > 0$$

for fixed "m"

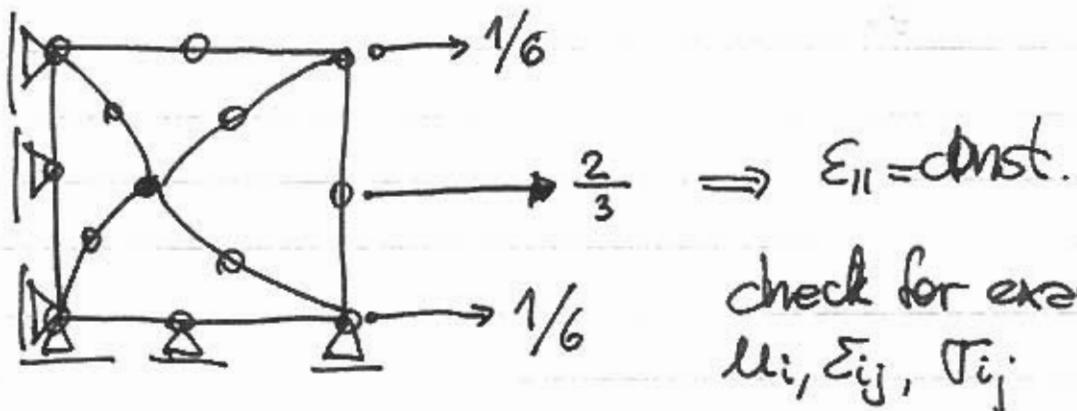
$$k > m-1$$

COMPLETENESS

Elasticity: $m=1 \Rightarrow k > 0$, "k" at least 1.

Patch test: Completeness \Rightarrow

constant strain state must be included
 exactly in the interpolation (up to machine precision)



- $\|u - u_h\| \leq C(u) \sigma^m h^{k+1-m}$

$$C(u) = C \|u\|_{m+1}$$

$$\|u\|_{m+1} = \left[\int_B |\nabla^{m+1} u|^2 dV \right]^{1/2}$$

strain gradients

\Rightarrow high strain-gradients in exact solution slow down convergence.