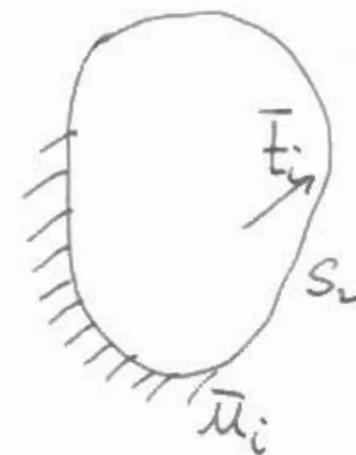


Boundary value problem

$$\begin{cases} \sigma_{ij,j} + f_i = 0 & \text{in } B \\ \sigma_{ij} n_j = \bar{t}_i & \text{on } S_2 \\ u_i = \bar{u}_i & \text{on } S_1 \\ \epsilon_{ij} = \frac{1}{2} (u_{i,j} + u_{j,i}) & \text{in } B \end{cases}$$



$\epsilon^e, \epsilon^p$  incompatible

Variational principle (minimum potential energy)

Assume:  $\exists U(\varepsilon, \gamma)$  internal energy potential  
 $\gamma = (\xi^P, \eta) /$

$$\sigma_{ij} = \frac{\partial U(\varepsilon, \gamma)}{\partial \varepsilon_{ij}}$$

Lubliner (1972):

$$U(\varepsilon, \gamma) = U^e(\varepsilon - \varepsilon^p) + U^p(\eta) \quad (\text{decoupled})$$

Example: linear elasticity  $U^e = \frac{1}{2} C_{ijkl} \varepsilon_{ij}^e \varepsilon_{kl}^e$

Equilibrium: minimize potential energy for  $B$  with respect to  $u_i$ , for fixed  $\gamma$

$$J = \int_B U(\varepsilon, \gamma) dV - \int_B f_i u_i dV - \int_{S_2} T_i u_i ds$$

$\uparrow$

$$\varepsilon_{ij} = \frac{1}{2} (u_{i,j} + u_{j,i})$$

$$\dot{y}(x, t) = f(\varepsilon(x, t), y(x, t)) \quad \text{kinetic equation}$$

Algorithms: Need to integrate constitutive relation in time at all quadrature points.

## Time-stepping algorithms for constitutive relations

Given  $\epsilon_n, \tau_n, q_n, \epsilon_n^p$  and

$\epsilon_{n+1}$  (strain driven)

Compute  $\tau_{n+1}, q_{n+1}, \epsilon_{n+1}^p$

General algorithm:

$$\tau_{n+1} = \hat{\sigma}(\epsilon_{n+1}; \epsilon_n, \tau_n, q_n, \epsilon_n^p, \Delta t)$$

combine state vector  $\Lambda = (\epsilon, \tau, q, \epsilon^p)$

$$\boxed{\tau_{n+1} = \hat{\sigma}(\epsilon_{n+1}; \Lambda_n, \Delta t)}$$

↑ strain driven

Integrate into global solution procedure

Equilibrium: Principle of virtual displacements

$$\int_B \tau_{ij} \eta_{ij} dV - \int_B f_i \eta_i dV - \int_{S_2} \bar{t}_i \eta_i ds = 0$$

$\forall \eta$  admissible

finite element discretization:

$$\sum_e \int_{\Omega^e} B^T \sigma \, dV - f^{ext} = 0$$

Enforce at time (load step)  $t = t_{n+1}$

$$\sum_e \int_{\Omega^e} B^T \sigma_{n+1} \, dV - f^{ext}_{n+1} = 0$$

Insert ... update:

$$\sum_e \int_{\Omega^e} B^T \hat{\sigma}(e_{n+1}; \Delta_n, \Delta t) \, dV - f^{ext}_{n+1} = 0$$

Compatibility:  $e_{n+1} = B u_{n+1}$

$$\boxed{\sum_e \int_{\Omega^e} B^T \hat{\sigma}(B u_{n+1}; \Delta_n, \Delta t) \, dV - f^{ext}_{n+1} = 0}$$

→ System of nonlinear algebraic equations  
for  $u_{n+1}$

Updated " $u_{n+1}$ " satisfies:

$$\left\{ \begin{array}{l} \sum_e \int_{S^e} B^T \sigma_{n+1} dV - f_{n+1}^{\text{ext}} = 0 \\ \sigma_{n+1} = \hat{\sigma}(\varepsilon_{n+1}; \Lambda_n, \Delta t) \\ \varepsilon_{n+1} = B u_{n+1} \end{array} \right.$$

Numerical quadrature:

$$\sum_e \sum_q w_p^e B^T(\xi_q) \hat{\sigma}(\xi_q) - f_{n+1}^{\text{ext}} = 0$$

↑ state variables sampled  
at quadrature points

Newton-Raphson solution procedure

$$t_n \rightarrow t_{n+1}, f_n^{\text{ext}} \rightarrow f_{n+1}^{\text{ext}}$$

$$(k)\text{-th iteration: } u_{n+1}^{(k)} \rightarrow u_{n+1}^{(k+1)}$$

$$u_{n+1}^{(k+1)} = u_{n+1}^{(k)} + \Delta u$$

$$f^{\text{int}}(u_{n+1}^{(k)} + \Delta u) = f_{n+1}^{\text{ext}}, \text{ linearize}$$

$$f^{\text{int}}(u_{n+1}^{(k)}) + \underbrace{\frac{\partial f^{\text{int}}}{\partial u}(u_{n+1}^{(k)})}_{K(u_{n+1}^{(k)})} \Delta u + \text{h.o.t} = f_{n+1}^{\text{ext}}$$

$$K(u_{n+1}^{(k)}) \Delta u = r(u_{n+1}^{(k)}) = f_{n+1}^{\text{ext}} - f(u_{n+1}^{(k)})$$

Consistent Tangent stiffness

$$K = \frac{\partial f^{\text{int}}}{\partial u} = \frac{\partial}{\partial u} \left[ \sum_e \int_{\Omega^e} B^T \hat{\Gamma}(B u; \Delta_n, \Delta t) dV \right]$$

$$= \sum_e \left[ \int_{\Omega^e} B^T \underbrace{\frac{\partial \hat{\Gamma}}{\partial \varepsilon} (B u; \Delta_n, \Delta t)}_{C(B u)} B dV \right]$$

$$C = \frac{\partial \hat{\Gamma}}{\partial \varepsilon} (\varepsilon; \Delta_n, \Delta t)$$

CONSISTENT  
TANGENT MODULI

Note  $C$  is obtained by linearization of the constitutive update algorithm

$$\Rightarrow K = \sum_e K^e = \sum_e \int_{\Omega^e} B^T C B dV$$

Newton-Raphson solution:

$$r(u_{n+1}^{(k)}) = f_{n+1}^{\text{ext}} - \sum_e \int_{\Omega^e} B^T \hat{\Gamma}(\underbrace{B u_{n+1}^{(k)}}_{\varepsilon_{n+1}^{(k)}}, \Delta_n, \Delta t) dV$$

$$K(u_{n+1}^{(k)}) = \sum_e \int_{\Omega^e} B^T C(\varepsilon_{n+1}^{(k)}; \Delta_n, \Delta t) B dV$$

Note state remains fixed at  $\Delta_n$  during the Newton Raphson iterations. The state is updated at the end of the load (time) step.

### Constitutive update algorithms

#### Backward Euler, fully implicit

$$\dot{\sigma} = C(\dot{\epsilon} - \dot{\lambda} r(\sigma, \dot{q}))$$

$$\dot{q} = \dot{\lambda} h(\sigma, \dot{q})$$

$$\dot{\lambda} = \begin{cases} \frac{\phi(\sigma, \dot{q})}{\eta} & \text{if } \phi \geq 0 \\ 0 & \text{if } \phi < 0 \end{cases}$$

$$\left\{ \begin{array}{l} \sigma_{n+1} = \sigma_n + C(\Delta \epsilon - \Delta \lambda r_{n+1}) \\ q_{n+1} = q_n + \Delta \lambda h_{n+1} \end{array} \right.$$

$$\left\{ \begin{array}{l} \Delta \lambda = \frac{\phi_{n+1}}{\eta} \end{array} \right.$$

where

$$r_{n+1} = r(\sigma_{n+1}, q_{n+1}); h_{n+1} = h(\sigma_{n+1}, q_{n+1}); \phi_{n+1} = \phi(\sigma_{n+1}, q_{n+1})$$

This defines a system of nonlinear algebraic equations in:  $\sigma_{n+1}$ ,  $q_{n+1}$ ,  $\Delta\lambda$ .

Rate independent limit:  $\eta \rightarrow 0$

$$\boxed{\phi_{n+1} = 0} \quad \begin{array}{l} \text{yield} \\ \text{criterion} \\ \text{at } n+1 \end{array}$$

### Geometrical interpretation

#### Elastic predictor

$$\left. \begin{array}{l} \sigma_{n+1}^* = \sigma_n + C \Delta\varepsilon \\ q_{n+1}^* = q_n, \quad \Delta\lambda^* = 0 \end{array} \right\} \begin{array}{l} \text{neglects} \\ \text{plasticity} \end{array}$$

Two possibilities:

- $\phi_{n+1}^* = \phi(\sigma_{n+1}^*, q_{n+1}^*) < 0$

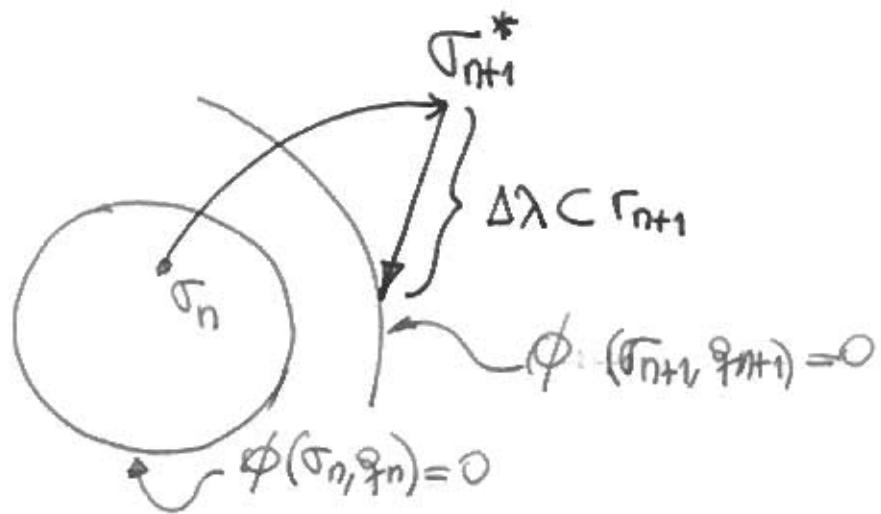
Then  $\sigma_{n+1} = \sigma_{n+1}^*$ ,  $q_{n+1} = q_{n+1}^*$ ,  $\Delta\lambda = \Delta\lambda^*$

DONE

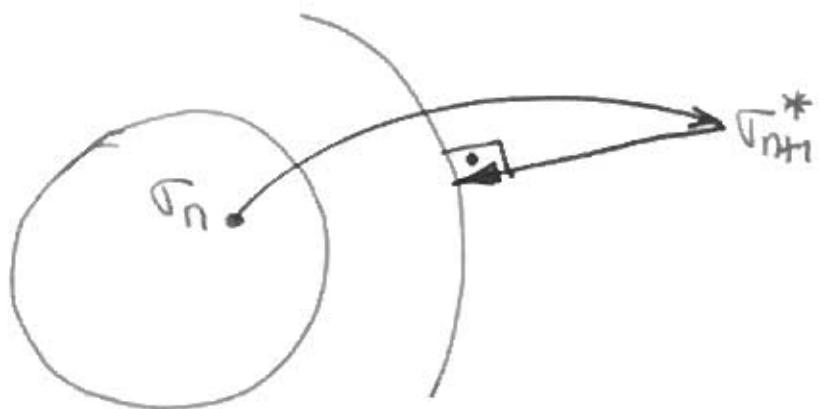
- otherwise  $\rightarrow$  plastic corrector

$$\sigma_{n+1} = \sigma_{n+1}^* - \Delta\lambda \Gamma_{n+1}$$

$$q_{n+1} = q_{n+1}^* + \Delta\lambda h_{n+1}, \quad \Delta\lambda = \phi_{n+1}/\eta$$



$$\text{if } r \propto \frac{\partial \phi}{\partial \sigma} \Rightarrow$$



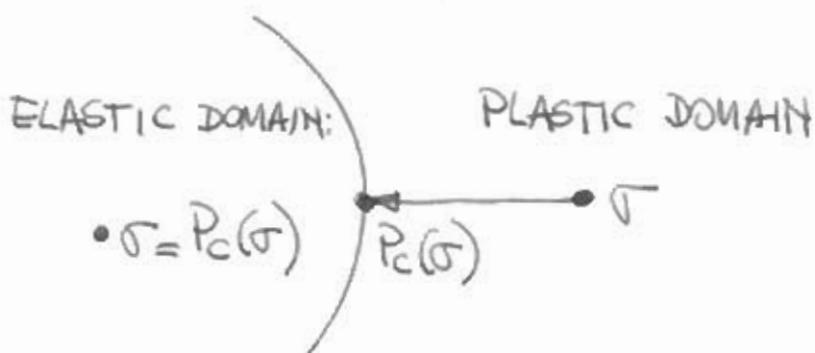
Closest point projection algorithms { always defined if elastic domain is convex

$r_{n+1}$  is closest to  $r_{n+1}^*$ :

$(r_{n+1} - r_{n+1}^*) : C^{-1} : (r_{n+1} - r_{n+1}^*)$  is minimized  
(plastic work)

$\sigma_{n+1}$  closest to  $\sigma_{n+1}^*$  in norm  $\|\sigma\| = \sigma : C^{-1} : \sigma$

$\rightarrow \|\sigma_{n+1}^* - \sigma_{n+1}\|$  minimum

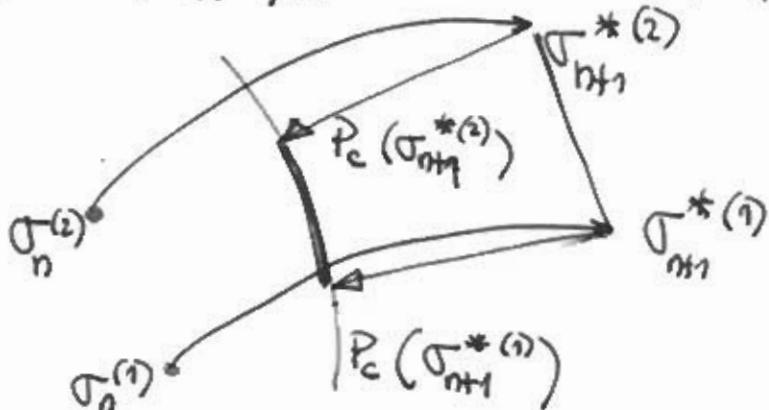


$P_C \equiv$  closest point projection onto boundary of elastic domain (yield surface)  
Well defined for any Convex elastic domain with or without corners

$$\boxed{\sigma_{n+1} = P_{C_{n+1}}(\sigma_{n+1}^*)}$$

$P_C$  is contractive if elastic domain is convex

$$\|P_C(\sigma_{n+1}^{*(1)}) - P_C(\sigma_{n+1}^{*(2)})\| \leq \|\sigma_{n+1}^{*(1)} - \sigma_{n+1}^{*(2)}\|$$



→ fully implicit algorithm is contractive  
 (errors in initial conditions are reduced  
 by algorithm)

### STRONG STATEMENT OF STABILITY

Specific model: J<sub>2</sub>-isotropic hardening -  
fully implicit

$$\dot{\epsilon}_{ij}^p = \lambda \frac{3}{2} \frac{s_{ij}}{\sigma} ; \bar{\sigma} = \left( \frac{3}{2} s_{ij} s_{ij} \right)^{1/2} ;$$

$$\tau_{ij} = \frac{3}{2} \frac{s_{ij}}{\sigma}$$

$$\tau_{ij} = C_{ijkl} \epsilon_{lk}^e$$

$$\lambda = \dot{\epsilon}_0 \left[ \left( \frac{\sigma}{\sigma_0} \right)^m - 1 \right] = \frac{\phi}{\eta} \quad \text{if } \phi \geq 0$$

$$\eta = \frac{\sigma_y}{\dot{\epsilon}_0} ; \phi = \sigma_y \left[ \left( \frac{\sigma}{\sigma_0} \right)^m - 1 \right]$$

$$\sigma_0(\bar{\epsilon}) = \sigma_y \left( 1 + \frac{\bar{\epsilon}}{\dot{\epsilon}_0} \right)^{1/m}$$

$$\bar{\epsilon} = \int_0^t \lambda dt$$

$$\sigma_{n+1} = \sigma_n + C (\Delta \varepsilon - \Delta \lambda \tau_{n+1})$$

$$\Delta \lambda = \frac{\Delta t}{\eta} \phi(\bar{\sigma}_{n+1}, \sigma_{0,n+1})$$

$$\sigma_{0,n+1} = \sigma_0 (\varepsilon_n + \Delta \lambda)$$

Isotropic elasticity

$$\boxed{p_{n+1} = p_n + K \Delta \varepsilon_{kk}}$$

pressure update  
is elastic

$$s_{n+1} = s_n + 2\mu \left( \Delta \varepsilon - \Delta \lambda \frac{3}{2} \frac{s_{n+1}}{\bar{\sigma}_{n+1}} \right)$$

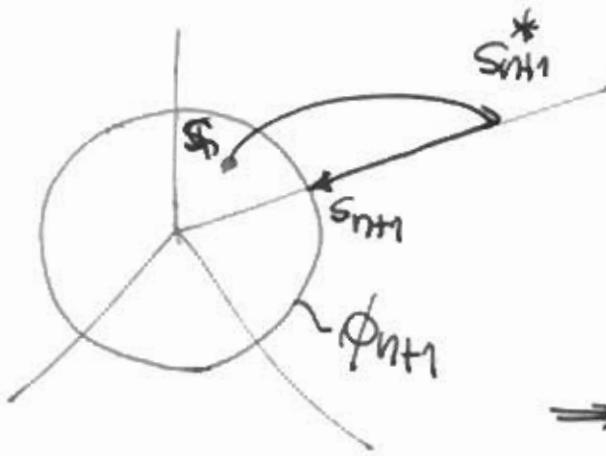
$$\Delta \varepsilon = \frac{\Delta \varepsilon_{kk}}{3} I + \Delta \varepsilon$$

$$s_{n+1}^* = s_n + 2\mu \Delta \varepsilon$$

$$\bar{\sigma}_{n+1}^* = \left( \frac{3}{2} s_{n+1}^* : s_{n+1}^* \right)^{1/2}$$

$$s_{n+1} = s_{n+1}^* - \left( \Delta \lambda \frac{3}{2} \frac{s_{n+1}}{\bar{\sigma}_{n+1}} \right) 2\mu$$

$$= s_{n+1}^* - 3\mu \Delta \lambda \frac{s_{n+1}}{\bar{\sigma}_{n+1}}$$



$$s_{n+1} = C s_{n+1}^*$$

$$\bar{s}_{n+1} = C \bar{s}_{n+1}^*$$

$$\Rightarrow \frac{s_{n+1}}{\bar{s}_{n+1}} = \frac{s_{n+1}^*}{\bar{s}_{n+1}^*}$$

$$C s_{n+1}^* = s_{n+1}^* - 3\mu \Delta \lambda \frac{s_{n+1}^*}{\bar{s}_{n+1}^*}$$

$$\left( C - 1 + 3\mu \frac{\Delta \lambda}{\bar{s}_{n+1}^*} \right) s_{n+1}^* = 0$$

$$\Rightarrow C = 1 - 3\mu \frac{\Delta \lambda}{\bar{s}_{n+1}^*}$$

$$\begin{cases} \bar{s}_{n+1} = \bar{s}_{n+1}^* - 3\mu \Delta \lambda \\ \Delta \lambda = \frac{\Delta t}{\eta} \phi(\bar{s}_{n+1}, s_{n+1}) \\ s_{n+1} = s_0 (\bar{s}_n + \Delta \lambda) \end{cases}$$

3 equations, 3 unknowns. Can be transformed into 1 equation with 1 unknown ( $\Delta \lambda$ )

$$\Delta\lambda = \frac{\Delta t}{\eta} \phi(\bar{\tau}_{n+1}^* - 3\mu \Delta\lambda, \sigma(\epsilon_n + \Delta\lambda))$$

One scalar equation with one unknown " $\Delta\lambda$ ",  
 solve by local Newton Raphson iteration.  
 For power-law viscosity:

$$f(\Delta\lambda) = \left( \frac{\Delta\lambda}{\dot{\varepsilon}_0 \Delta t} + 1 \right)^{1/m} - \frac{\bar{\tau}_{n+1}^* - 3\mu \Delta\lambda}{\sigma_0 (\epsilon_n + \Delta\lambda)} = 0$$