

for $0 < x < L/5$:

$$\theta^2(x) = \frac{0.45\omega}{(5u_0 x/L)^6} \int_0^x (5u_0 \xi/L)^5 d\xi \\ = 0.015 \frac{\omega L}{u_0}$$

$$\therefore \theta(x) = 3.87 \times 10^{-5} L$$

$$\lambda = \frac{\theta^2}{\nu} \frac{du}{dx} = 0.075 \rightarrow H = 2.36$$

for $L/5 < x < L$

$$\theta^2(x) = \theta^2(L/5) + \frac{0.45\omega}{u_0^6} \int_{L/5}^x u_0^5 dx \\ = 0.015 \left(\frac{\omega}{u_0 L}\right) L^2 + 0.45 \left(\frac{\omega}{u_0 L}\right) \cdot L^2 \left(\frac{x}{L} - \frac{1}{5}\right)$$

$$\therefore \theta(x) = 2.12 \times 10^{-4} \left\{ x/L - 1/5 \right\}^{1/2} L$$

$$\lambda = 0 \rightarrow H = 2.61$$

16) Michel's Criterion

$$\textcircled{1} \quad Re_\infty = Re_{crit} = 2.9 Re_x^{0.4}$$

for $0 < x < L/5$:

$$x/L = 10.9 \text{ so no transition}$$

$$L/5 < x < L : \quad x/L = 0.77 \text{ (transition)}$$

$$\textcircled{2} \quad Re_\infty = Re_{crit} = 1.174 \left[1 + \left(\frac{22,400}{Re_x, \bar{u}} \right) \right] Re_{x, \bar{u}}^{0.46}$$

No transition point

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Granville's Method

$$Re_0(x_{\bar{w}}) \approx Re_0(x_i) + 450 + 400 e^{60 \lambda_m}$$

$$\lambda_m = \frac{1}{x_{\bar{w}} - x_i} \int_{x_i}^{x_{\bar{w}}} \lambda(x) dx$$

for $x < L/5$, $\lambda(x) = \lambda_m = 0.075$. At initial instability $Re_0(x_0) \approx 2400$
 for $H = 2.36$, $Re_0(x=L/5) = 387 \Rightarrow$ clearly no transition point
 for $x < L/5$.

for $x > L/5$, $\lambda = 0$, $H = 2.61$, $Re_0(x_i) = 201$ which is $< Re_0(x=L/5)$
 \therefore initial instability will occur at $x = L/5$.

$$Re_0(x_{\bar{w}}) = 387 + 450 + 400 = 1237$$

$$\begin{aligned} Re_0(x) &= \left(\frac{\partial}{L}\right) \cdot Re_L \\ &= 2.12 \times 10^{-4} \left\{ x/L - \frac{1}{6} \right\}^{1/2} \cdot 1 \times 10^7 \end{aligned}$$

$$\Rightarrow \frac{x_{\bar{w}}}{L} \approx 0.5$$

eⁿ Method (Envelop Method)

$$q = \int_{Re_0}^{Re_0} \frac{dn}{dRe_0} \cdot dRe_0$$

$$\text{For } 0 < x < L/5 : H = 2.36, Re_{00} = 2462, \frac{dn}{dRe_0} = 0.005$$

At $x = L/5$, $Re_0 = 387 \Rightarrow$ no instability, no transition

$$\text{For } L/5 < x < L : H = 2.61, Re_{00} = 206, \frac{dn}{dRe_0} = 0.0117$$

$$q = (Re_{00} - 387) 0.0117 \Rightarrow Re_{00} = 1156 \rightarrow \frac{x_{\bar{w}}}{L} = 0.467$$

(3)

Rapid initial acceleration makes Michel's criterion dubious, since Re_0 for a given Re_x is lower than it would be for Blasius flow. The use of Re_x is inappropriate.

Both Granger's method and e^n method are more suitable since they track the suitability from $x=2/5$ where it really begins.

$$2\triangleright \text{Time averaging products: } \bar{uv} = \frac{1}{2T} \int_{-T}^T uv dt$$

$$\overline{(\bar{u} + \tilde{u} + u')(\bar{v} + \tilde{v} + v')} = \bar{u}\bar{v} + \overline{\tilde{u}\tilde{v}} + \overline{u'v'}$$

Terms like $\overline{\tilde{u}v'}$, $\overline{\tilde{u}u'}$, etc. are zero, since (\cdot) and $(\cdot)'$ have no frequencies in common.

Time Avg. equations:

$$\frac{\partial \bar{u}}{\partial x} + \frac{\partial \bar{v}}{\partial y} = 0 \quad \begin{matrix} \bar{u}(x,y) \\ \bar{v}(x,y) \end{matrix} - \text{dependant variables}$$

$$\bar{u} \frac{\partial \bar{u}}{\partial x} + \bar{v} \frac{\partial \bar{u}}{\partial y} = \bar{u} \frac{d \bar{u}}{dx} + v \frac{\partial^2 \bar{u}}{\partial y^2} - \frac{\partial}{\partial x} \left[\underline{\bar{u}^2} + \underline{\bar{u}'^2} \right]$$

terms which need to be modeled

$$2b\triangleright \text{Phased-locked ensemble averaging: } \langle uv \rangle = \frac{1}{N} \sum u_i v_i$$

$$\langle (\bar{u} + \tilde{u} + u')(\bar{v} + \tilde{v} + v') \rangle = (\bar{u} + \tilde{u})(\bar{v} + \tilde{v}) + \langle u'v' \rangle$$

Ensemble averaging has no effect on (\cdot) , $(\cdot)'$, or their products, since those quantities are always the same for all occurrences in the averaging summation.

Ensemble averaged TSL equations:

$$\frac{\partial \bar{\tilde{u}}}{\partial y} + \frac{\partial \bar{\tilde{v}}}{\partial y} = 0 \quad \begin{matrix} \bar{\tilde{u}}(x,y,t) = \bar{u} + \tilde{u} \\ \bar{\tilde{v}}(x,y,t) = \bar{v} + \tilde{v} \end{matrix} - \text{dep. variables}$$

$$\frac{\partial \bar{u}}{\partial t} + \bar{u} \frac{\partial \bar{u}}{\partial x} + \bar{v} \frac{\partial \bar{u}}{\partial y} = \frac{\partial \bar{u}_c}{\partial t} + \bar{u}_c \frac{\partial \bar{u}_c}{\partial x} \\ + 2 \frac{\partial^2 \bar{u}}{\partial y^2} - \frac{\partial}{\partial x} (\underline{\underline{\langle u'^2 \rangle}}) - \frac{\partial}{\partial y} (\underline{\underline{\langle u'v' \rangle}})$$

The equations are unsteady, "unsteady" terms need modeling

- 2c) The ensemble-averaged equations are more suitable for periodic-unsteady flows, since "mean" flow (dep. variables) ($\bar{u} + \bar{u}'$) retain the inherent unsteadiness. Turbulence model is needed for $\langle u'^2 \rangle$ and $\langle u'v' \rangle$. The time averaged equations lump the unsteadiness into the additional stresses. This would require an "unsteadiness model".