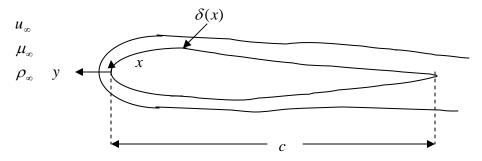
Laminar Boundary Layer Order of Magnitude Analysis



Assumptions (in addition to incompressible, steady & 2-D)

* Changes in x – direction occur over a distance c

$$\Rightarrow \frac{\partial}{\partial x} \sim \frac{1}{C} \quad \text{or, we write} \quad \frac{\partial}{\partial x} = O(\frac{1}{C})$$

* Changes in y – direction occur over a distance δ

$$\Rightarrow \frac{\partial}{\partial y} \sim \frac{1}{\delta} \quad \text{or,} \qquad \qquad \frac{\partial}{\partial y} = O(\frac{1}{\delta})$$

- * $\delta \ll C$ (boundary layer is thin)
- * Re = $\frac{\rho_{\infty}u_{\infty}c}{\mu_{\infty}}$ >> 1 (Reynolds number is large)
- * Changes in x velocity are proportional to u_{∞} $\Delta u = O(u_{\infty})$
- * Changes in pressure are proportional to $\rho_{\infty}u_{\infty}^2$ $\Delta p = O(\rho_{\infty}u_{\infty}^2)$

Now, let's do the order of magnitude analysis:

Conservation of Mass

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

Now, we introduce our assumed orders. For example:

$$\frac{\partial u}{\partial x} = O(\frac{u_{\infty}}{c})$$

Since we have not assumed an order for v variations, just leave this as simply Δv and say,

$$\frac{\partial v}{\partial y} = O(\frac{\Delta v}{\delta})$$

Since $\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$, they obviously have equal (though opposite) magnitudes and therefore orders. Thus, the changes in v scale as:

$$\Delta v = \mathcal{O}(u_{\infty} \frac{\delta}{C})$$

Since v = 0 at a wall, $\Delta v = O(v)$ itself and we can say

$$\frac{v}{u_{\infty}} = O(\frac{\delta}{c})$$

x – momentum

$$\rho u \frac{\partial u}{\partial x} + \rho v \frac{\partial u}{\partial y} = -\frac{\partial p}{\partial x} + \mu \frac{\partial^2 u}{\partial x^2} + \mu \frac{\partial^2 u}{\partial y^2}$$

$$O(\rho_{\infty} \frac{u_{\infty}^2}{c}) O(\rho_{\infty} \frac{u_{\infty}^2}{c}) O(\rho_{\infty} \frac{u_{\infty}^2}{c}) O(\frac{\mu_{\infty} u_{\infty}}{c^2}) O(\frac{\mu_{\infty} u_{\infty}}{\delta^2})$$

If we compare the last two terms, clearly

$$\frac{\mu_{\infty}u_{\infty}}{c^2} << \frac{\mu_{\infty}u_{\infty}}{\delta^2}$$

Since we have assumed $\delta << C$. Thus, we can assume that $\mu \frac{\partial^2 u}{\partial y^2}$ is small.

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In order for any viscous terms to remain, this requires that $\mu \frac{\partial^2 u}{\partial x^2}$ be equal in order to the remaining inviscid terms. That is, e.g

$$O(\rho u \frac{\partial u}{\partial x}) = O(\mu \frac{\partial^2 u}{\partial y^2})$$

$$\Rightarrow O(\frac{\rho_{\infty} u_{\infty}^2}{c}) = O(\frac{\mu_{\infty} u_{\infty}}{\delta^2})$$

$$\Rightarrow (\frac{\delta}{c})^2 = O(\frac{\mu_{\infty}}{\rho_{\infty} u_{\infty}})$$

In other words,

$$\frac{\delta}{C} = O(\frac{1}{\sqrt{\text{Re}}})$$
 \tag{\text{Major result!}}

y – momentum

$$\rho u \frac{\partial v}{\partial x} + \rho v \frac{\partial v}{\partial y} = -\frac{\partial p}{\partial y} + \mu \frac{\partial^2 v}{\partial x^2} + \mu \frac{\partial^2 v}{\partial y^2}$$

$$O(\rho_{\infty} u_{\infty}^2 \frac{\delta}{c^2}) O(\rho_{\infty} u_{\infty}^2 \frac{\delta}{c^2}) O(\frac{\rho_{\infty} u_{\infty}^2}{\delta}) O(\mu_{\infty} u_{\infty} \frac{\delta}{c^3}) O(\mu_{\infty} u_{\infty} \frac{1}{\delta c})$$

Normalizing all of these orders by $\frac{1}{\rho_{\infty}u_{\infty}^2/C}$ and substituting in for $\mathrm{Re}=\frac{\rho_{\infty}u_{\infty}C}{\mu_{\infty}}$ and $\frac{\delta}{C}=\mathrm{O}(\frac{1}{\sqrt{\mathrm{Re}}})$ gives:

$$O(\frac{1}{\sqrt{Re}}) \ O(\frac{1}{\sqrt{Re}}) \ O(\sqrt{Re}) \ O(\frac{1}{Re^{\frac{3}{2}}}) \ O(\frac{1}{\sqrt{Re}})$$

Clearly, except for $\frac{\partial p}{\partial y}$, all other terms are small as Reincreases. Thus, we must conclude that:

$$\Rightarrow \boxed{0 = \frac{\partial p}{\partial y} \text{ is small}}$$

$$\Rightarrow \boxed{0 = \frac{\partial p}{\partial y}}$$
 is the y – momentum equation for a boundary layer

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